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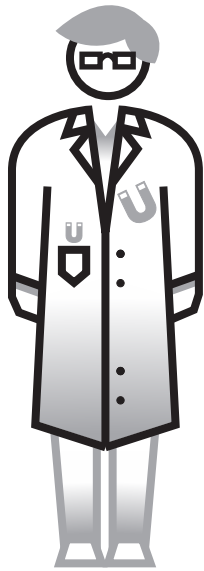
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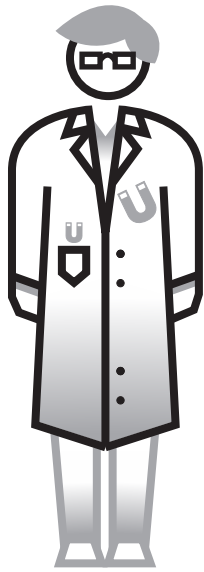
**Albert**



**Albert Fert**



# **Albert Fert Nobel Price**



**Albert Fert  
Nobel Prize  
2007**

**Albert Fert**  
**Nobel Price**  
**2007**





**Albert Fert**

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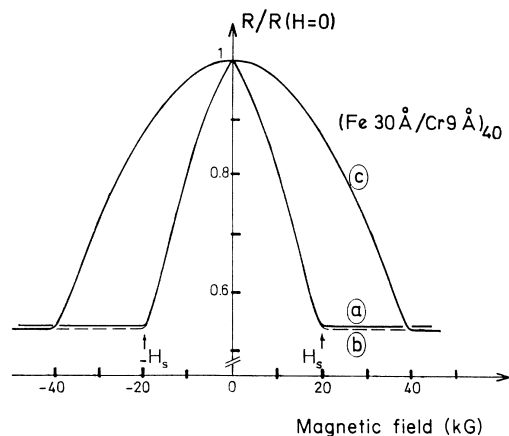


FIG. 2. Magnetoresistance of a  $[(\text{Fe } 30 \text{ \AA})/(\text{Cr } 9 \text{ \AA})]_{40}$  superlattice of 4.2 K. The current is along  $[110]$  and the field is in the layer plane along the current direction (curve *a*), in the layer plane perpendicular to the current (curve *b*), or perpendicular to the layer plane (curve *c*). The resistivity at zero field is  $54 \mu\Omega \text{ cm}$ . There is a small difference between the curves in increasing and decreasing field (hysteresis) that we have not represented in the figure. The superlattice is covered by a  $100\text{-\AA}$  Ag protection layer. This means that the magnetoresistance of the superlattice alone should be slightly higher.

of Grünberg *et al.*<sup>4</sup> and by the spin-polarized low-energy electron-diffraction experiments of Carbone and Alvarado.<sup>10</sup> The AF coupling between the Fe layers has been ascribed to indirect exchange interactions through the Cr layers, but a theoretical model of these interactions is still lacking.<sup>4,9</sup>

The magnetoresistance of the Fe/Cr superlattices has been studied by a classical ac technique on small rectangular samples. Examples of magnetoresistance curves at 4.2 K are shown in Figs. 2 and 3. The resistance decreases during the magnetization process and becomes practically constant when the magnetization is saturated. The curves *a* and *b* in Fig. 2 are obtained for applied fields in the plane of layers in the longitudinal and transverse directions, respectively. The field  $H_S$  is the field needed to overcome the AF couplings and to saturate the magnetization (compare with Fig. 1). In contrast, fields applied perpendicularly to the layers (curve *c*) have to overcome not only the AF coupling but also the magnetic anisotropy, so that the magnetoresistance is saturated at a field higher than  $H_S$ .

The most remarkable result exhibited in Figs. 2 and 3 is the huge value of the magnetoresistance. For  $t_{\text{Cr}}=9 \text{ \AA}$  and  $T=4.2 \text{ K}$ , see Fig. 2, there is almost a factor of 2 between the resistivities at zero field and in the saturated state, respectively (in absolute value, the resistivity change is about  $23 \mu\Omega \text{ cm}$ ). By comparison of the results for three different samples in Fig. 3, it can be seen

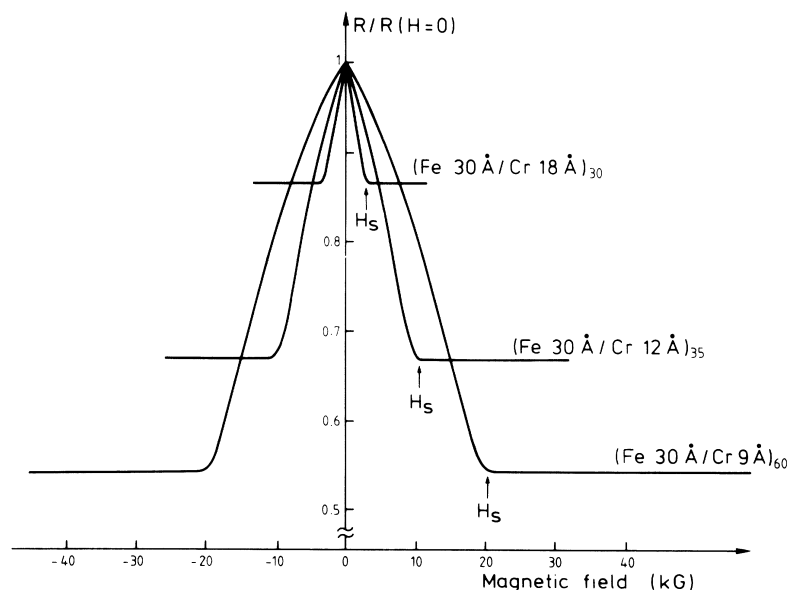


FIG. 3. Magnetoresistance of three Fe/Cr superlattices at 4.2 K. The current and the applied field are along the same  $[110]$  axis in the plane of the layers.

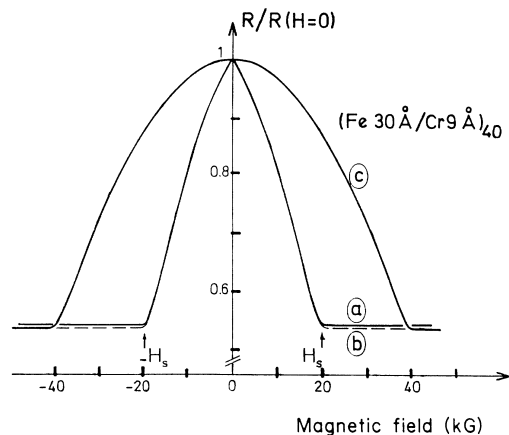


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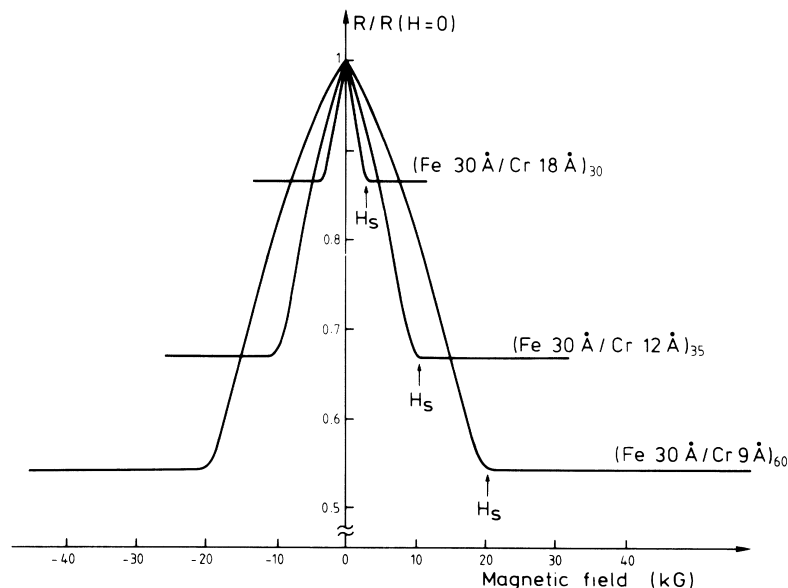


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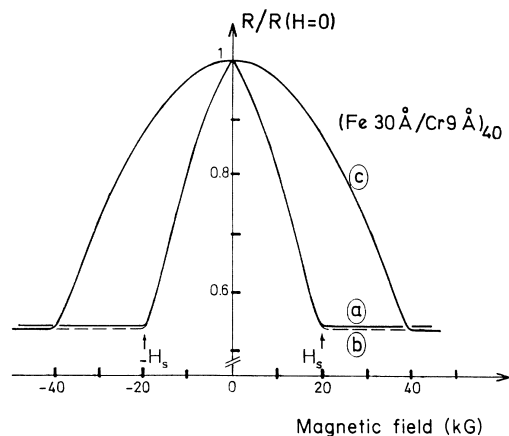


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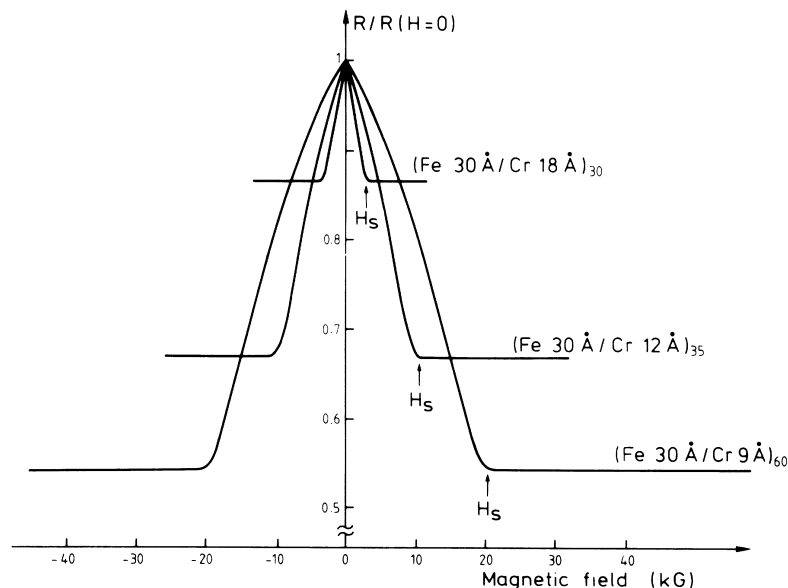
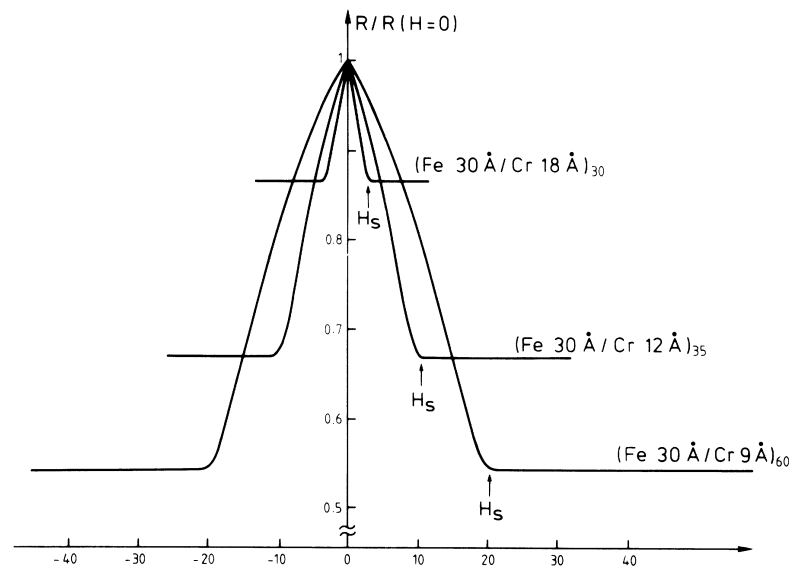
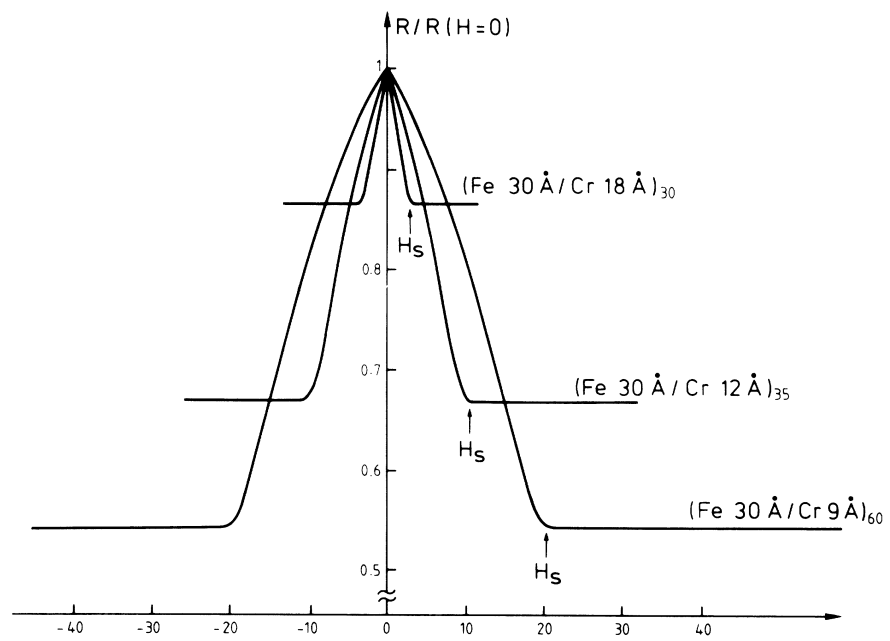
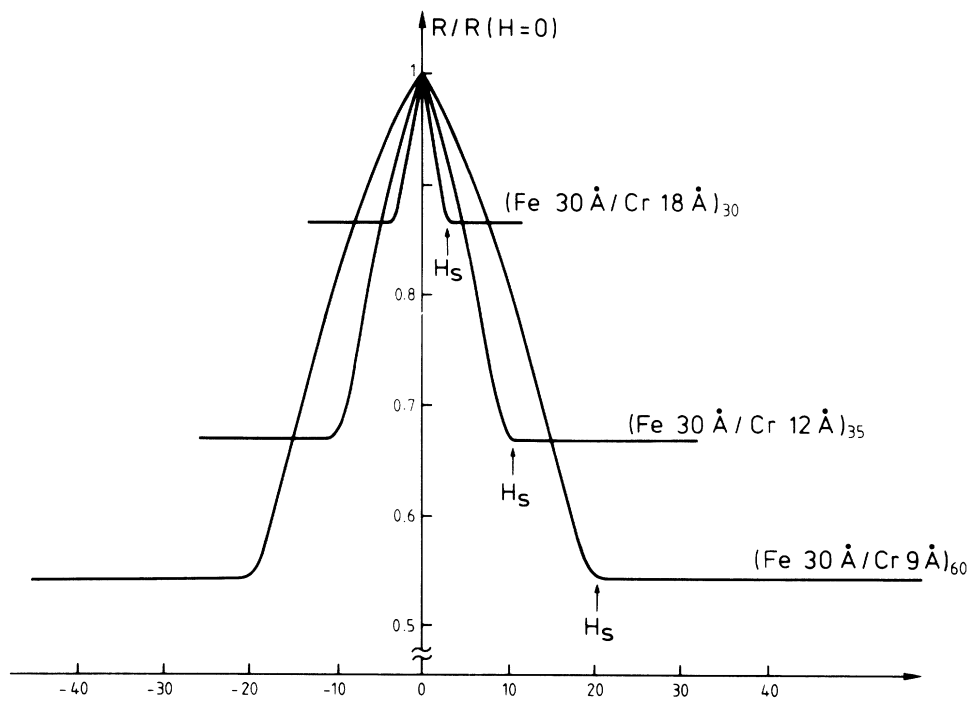


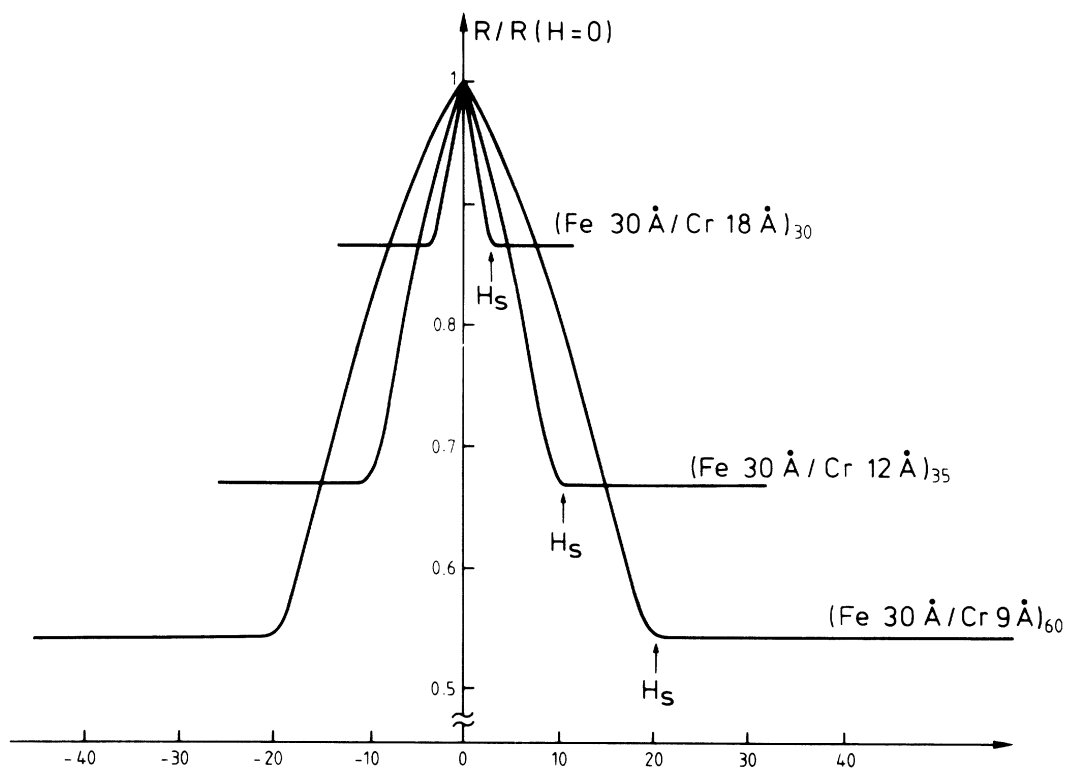
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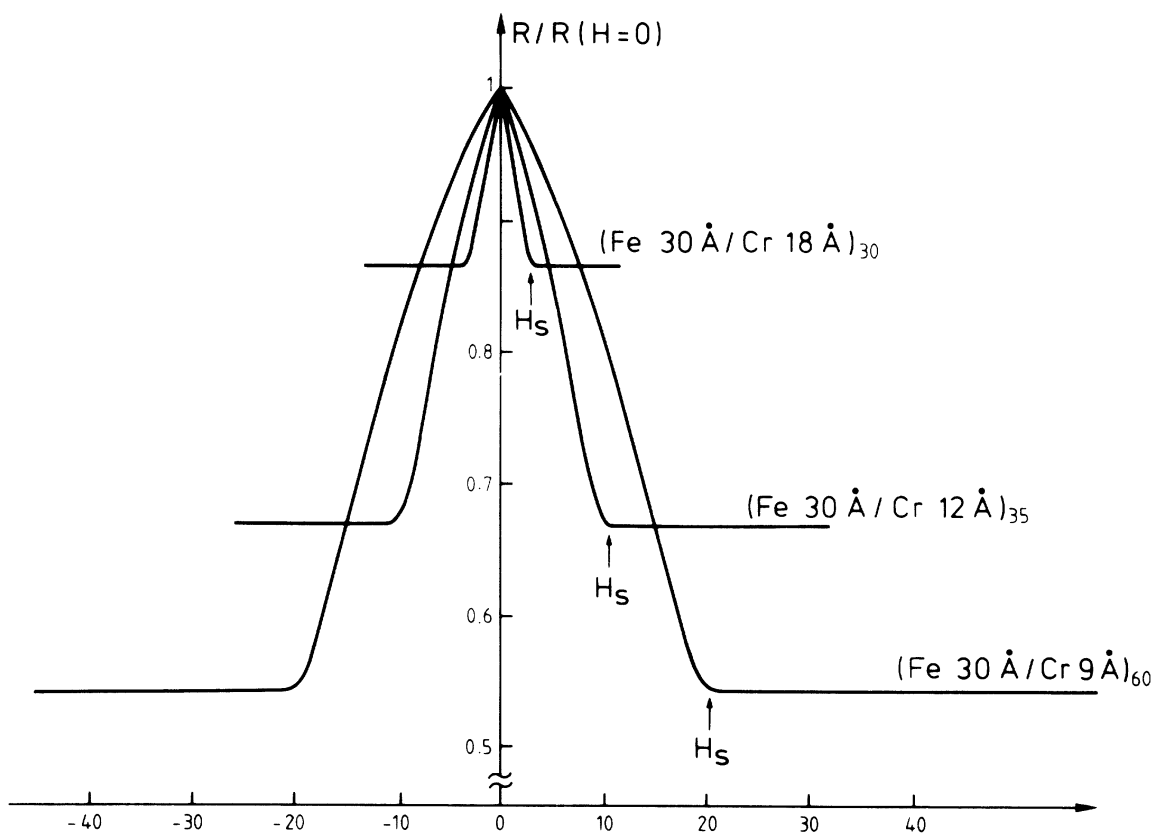


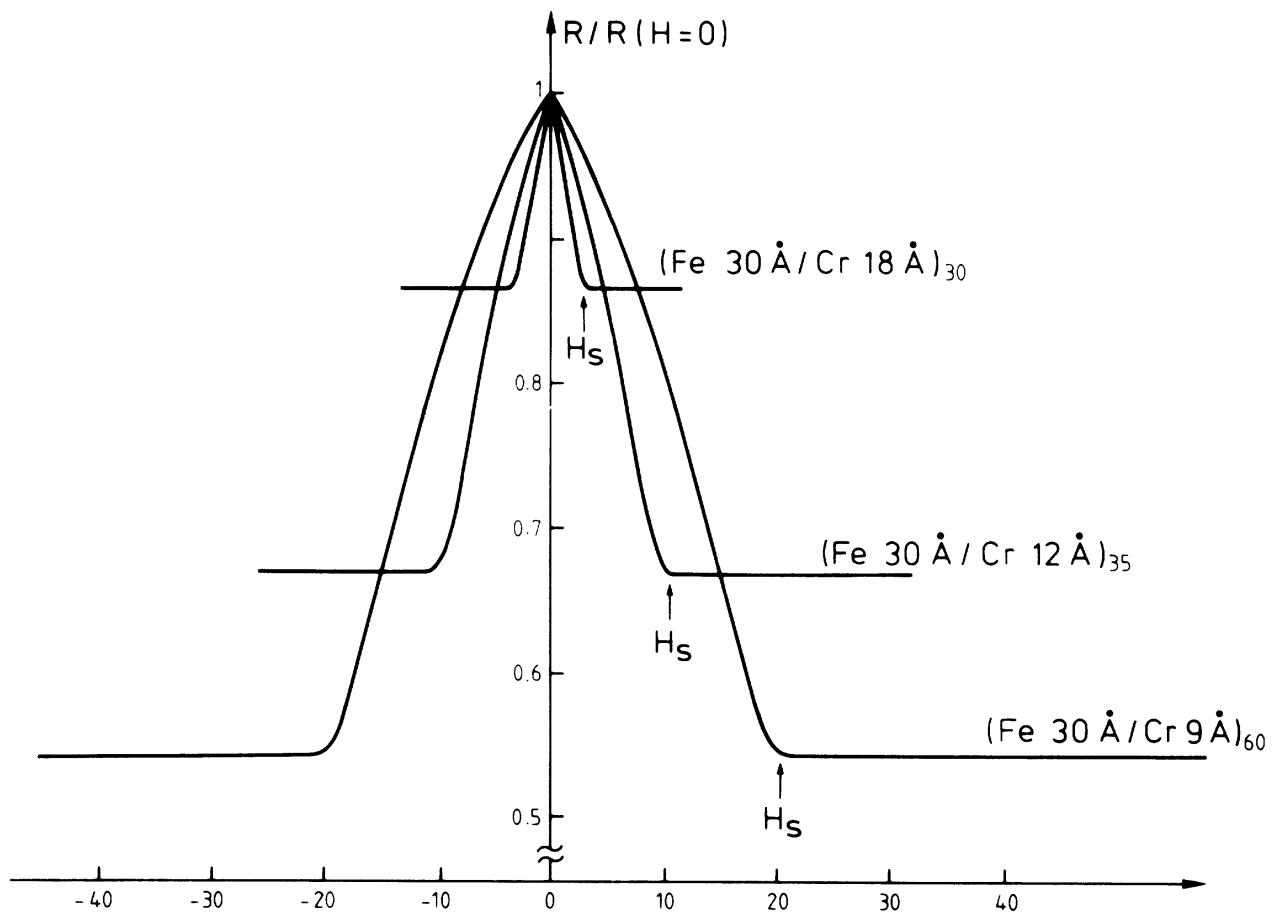


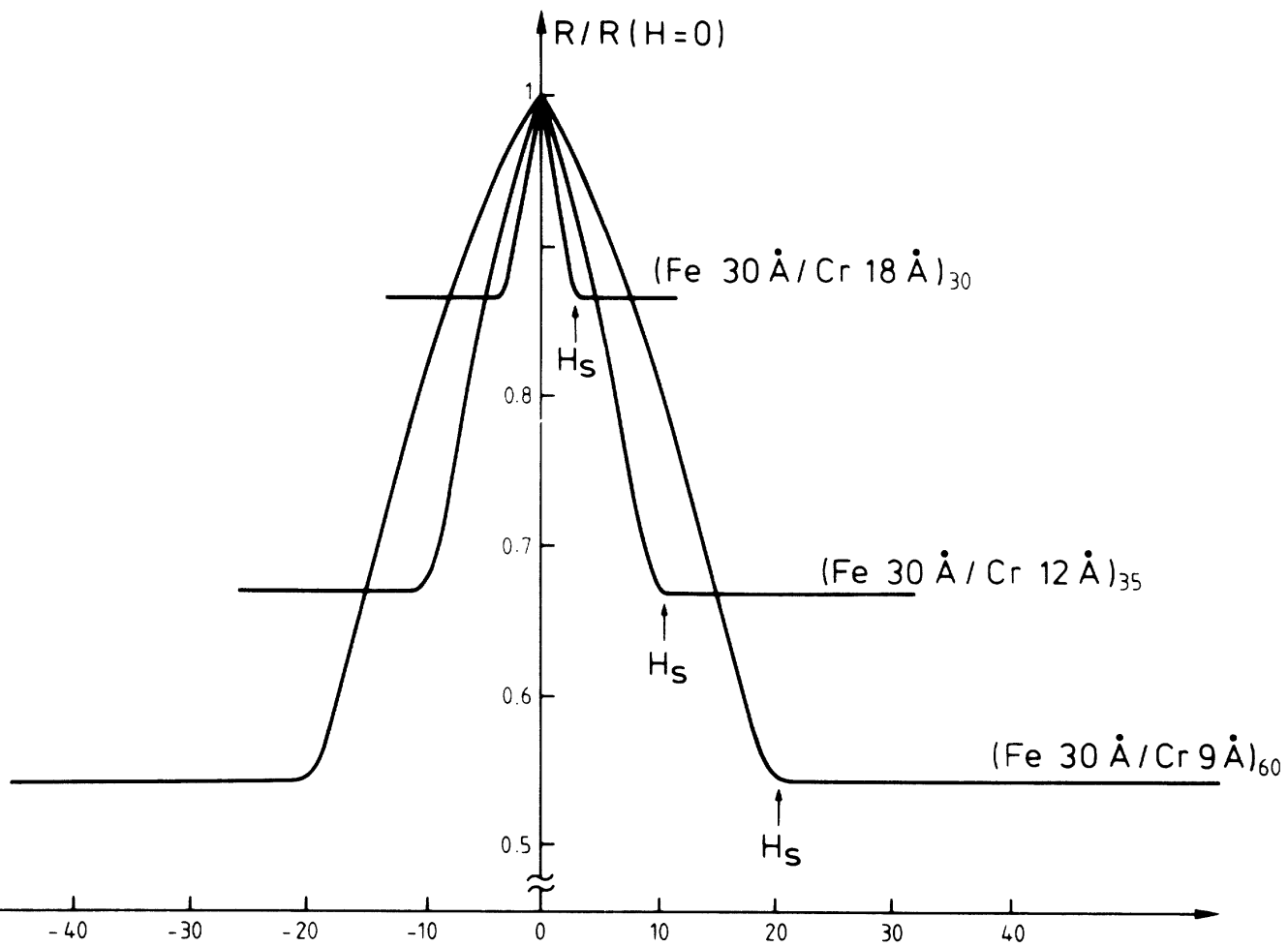


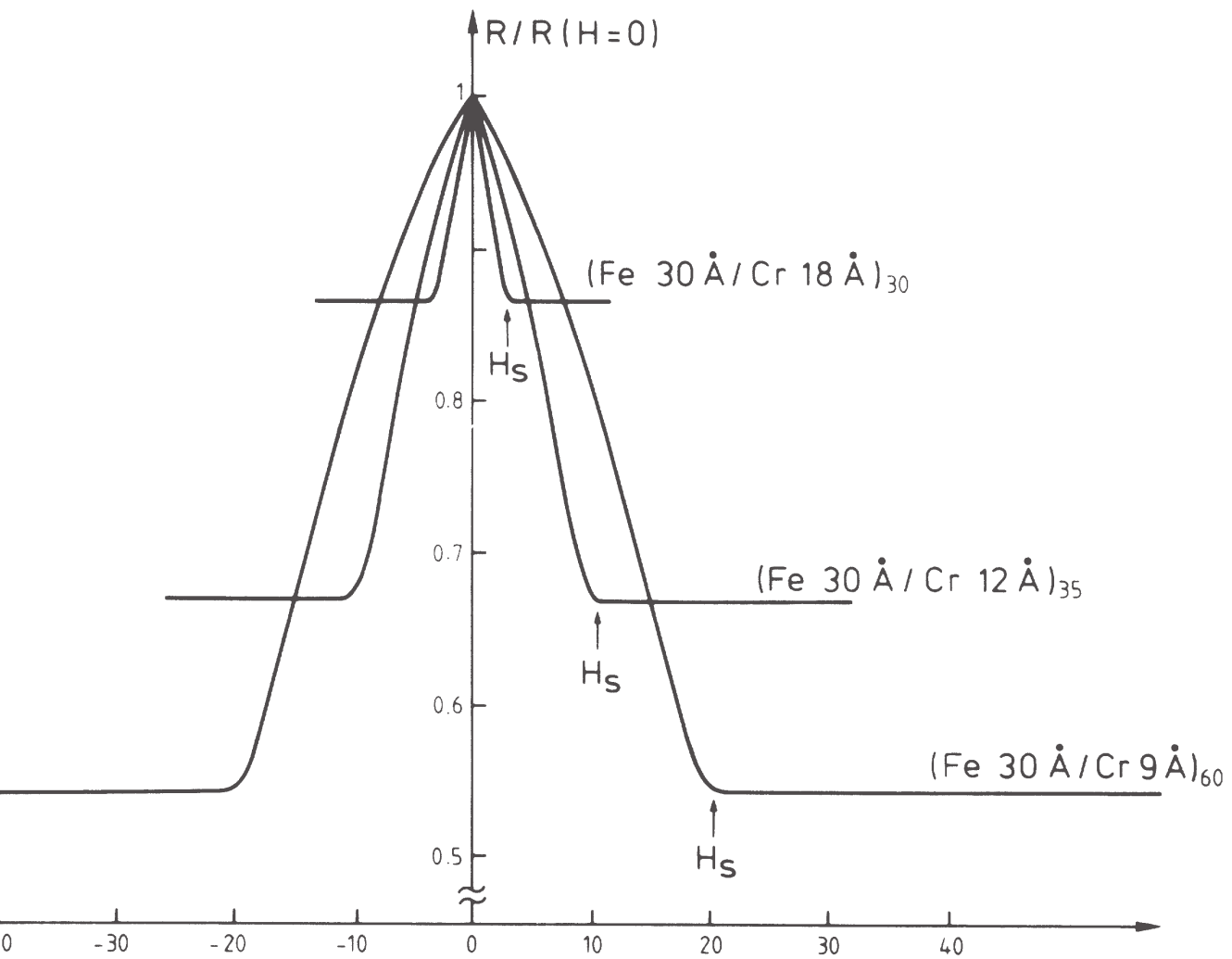






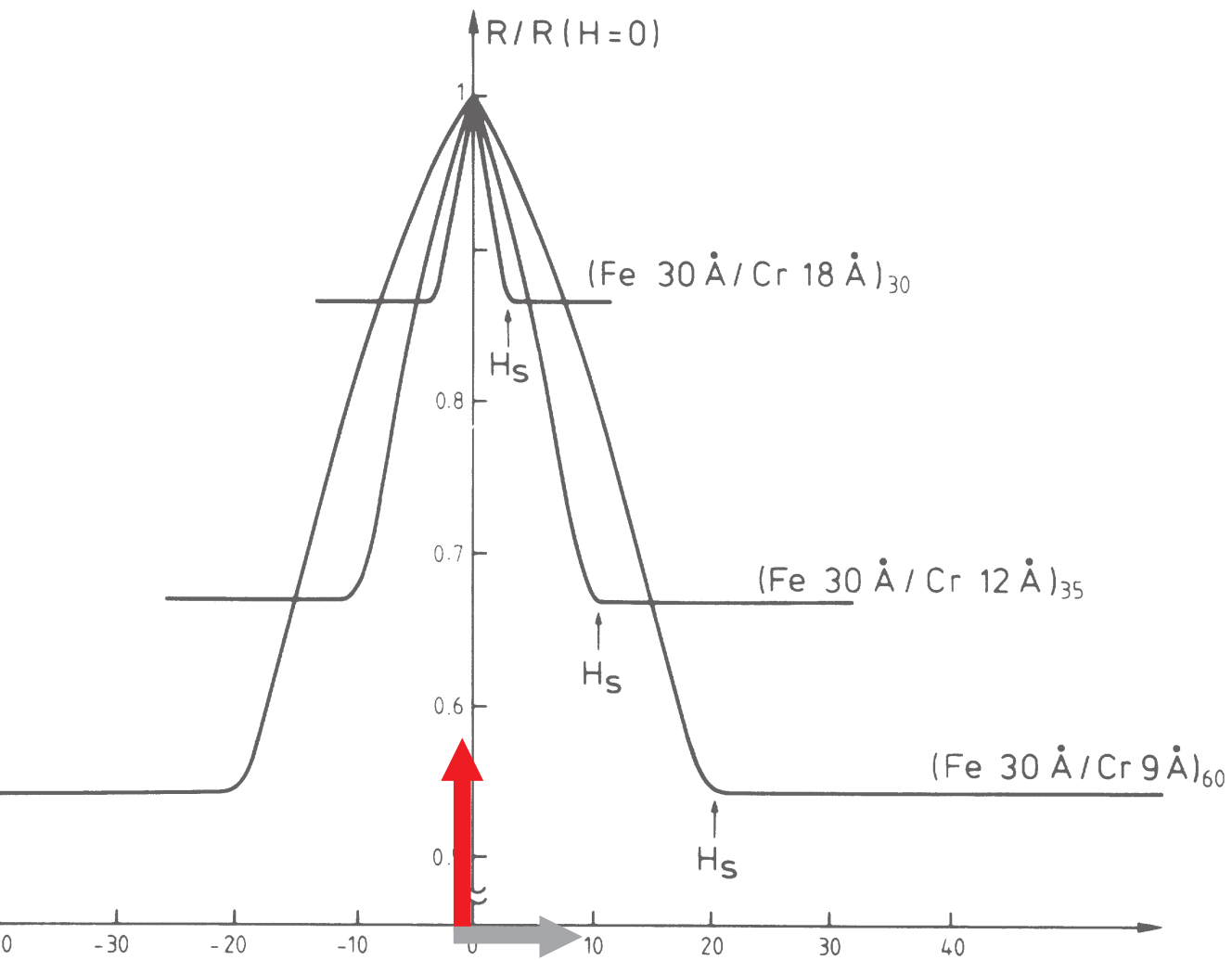






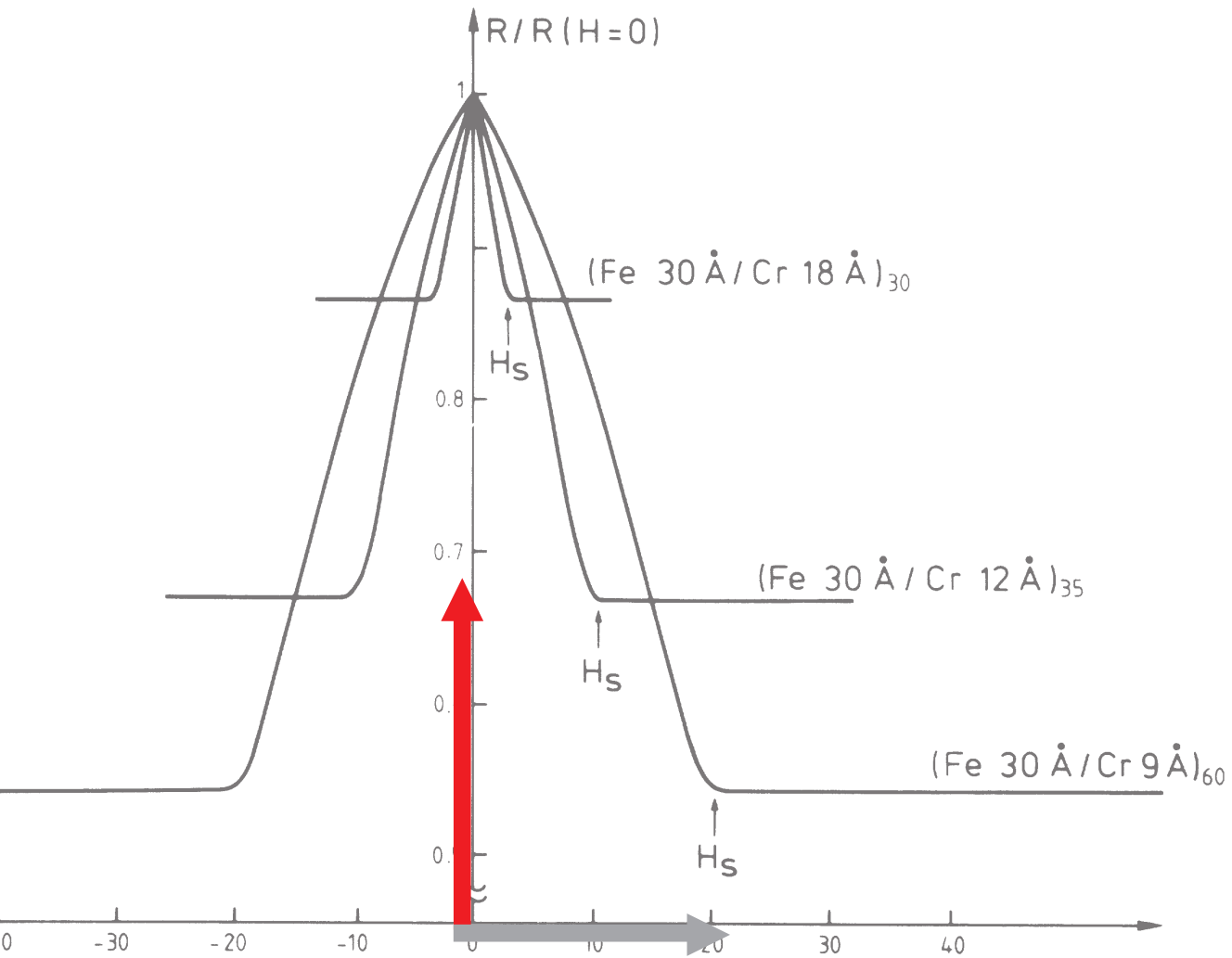


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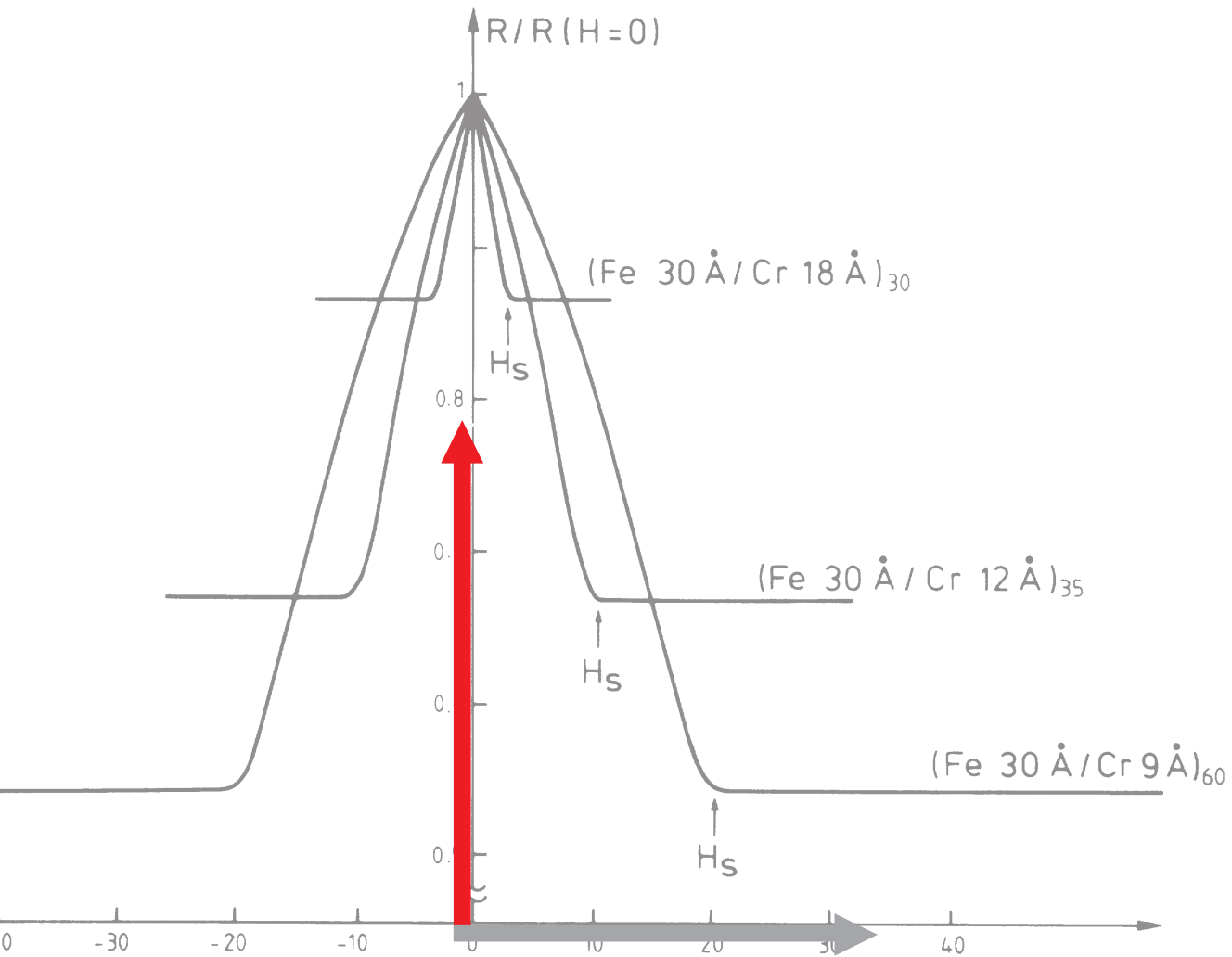
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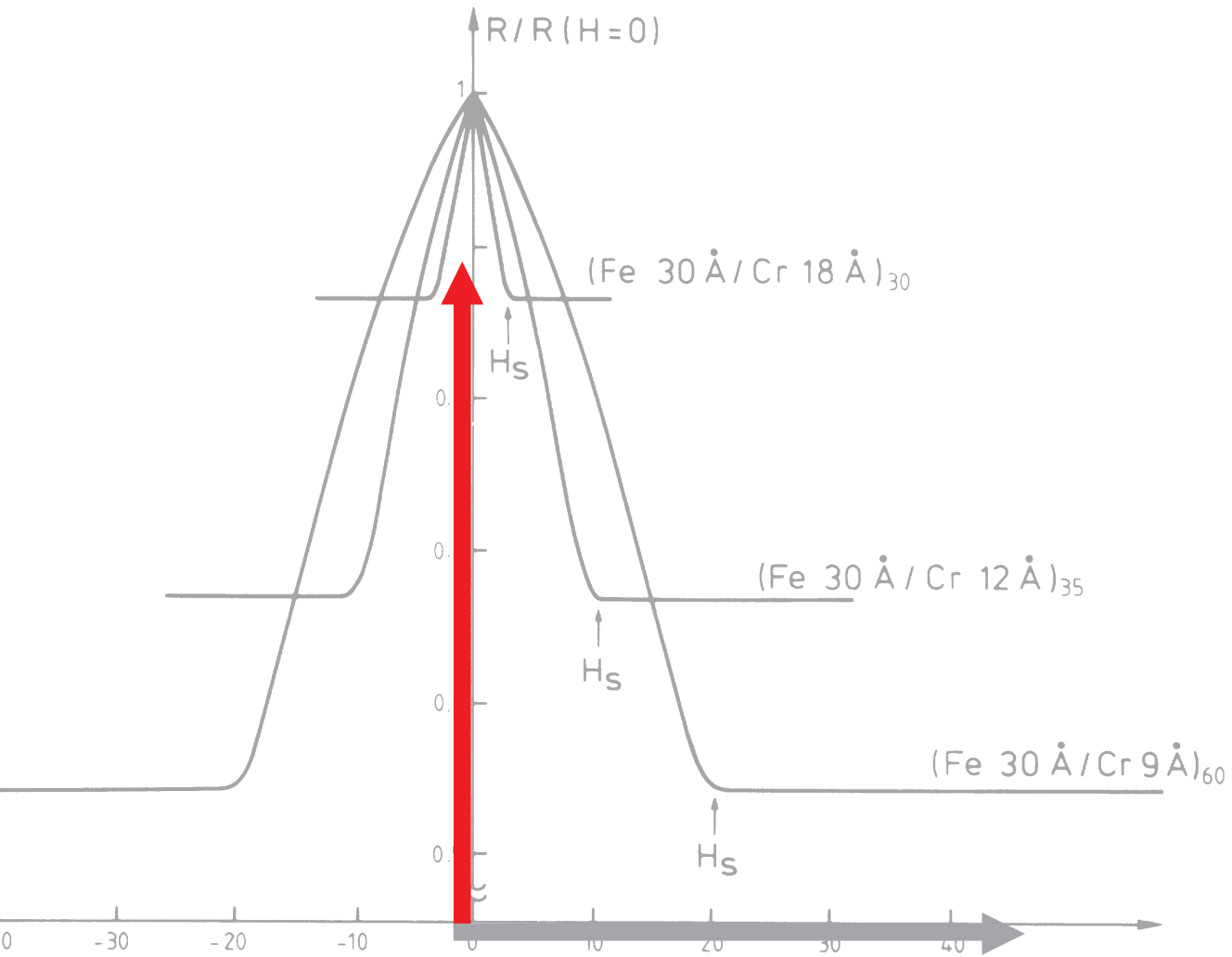
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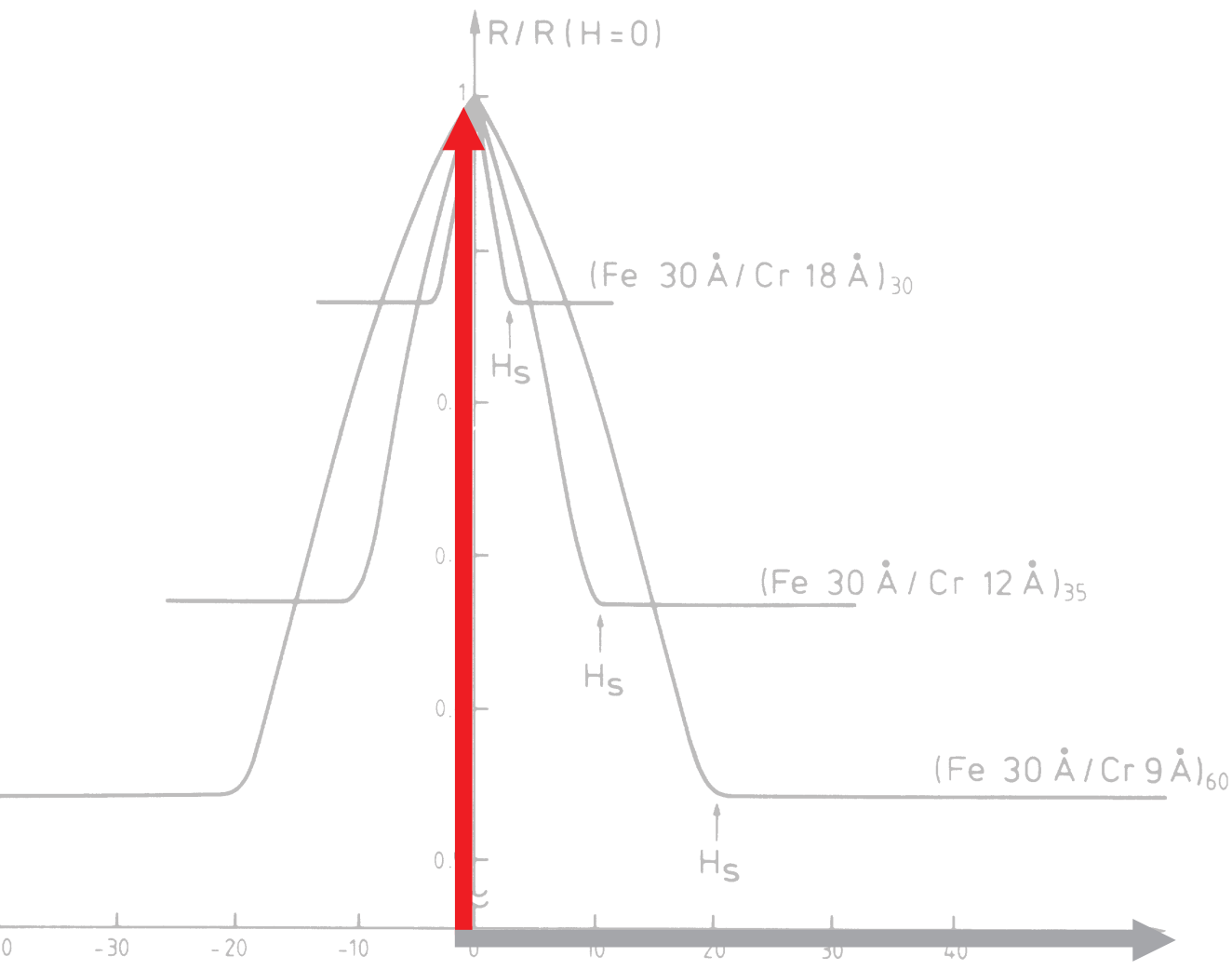
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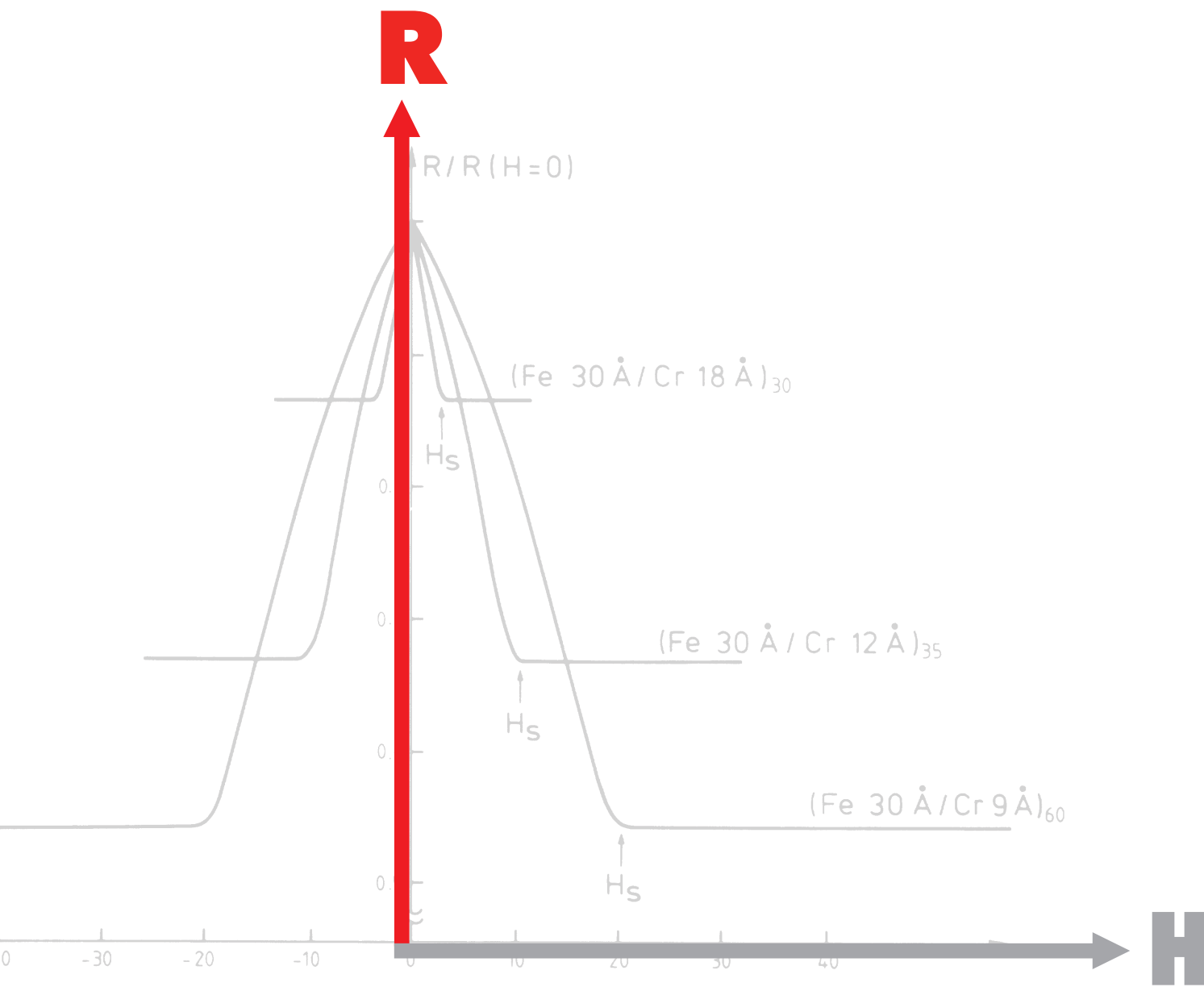


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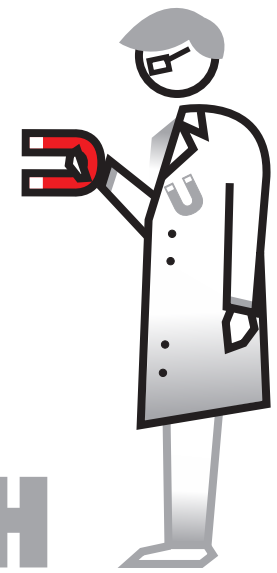
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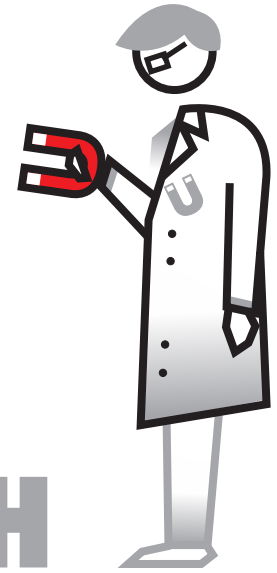
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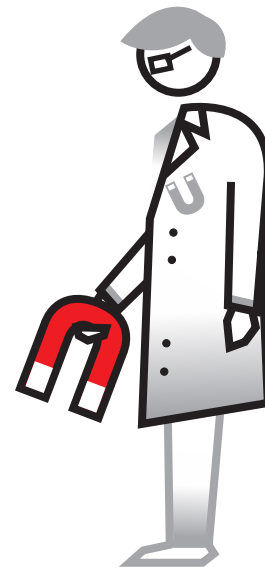
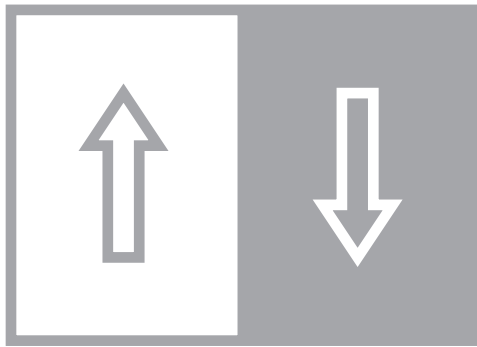


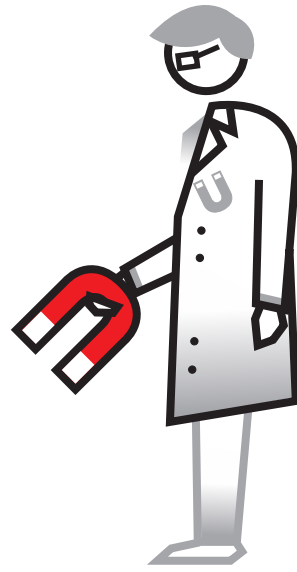
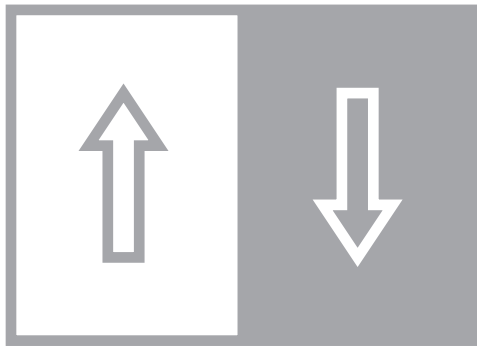
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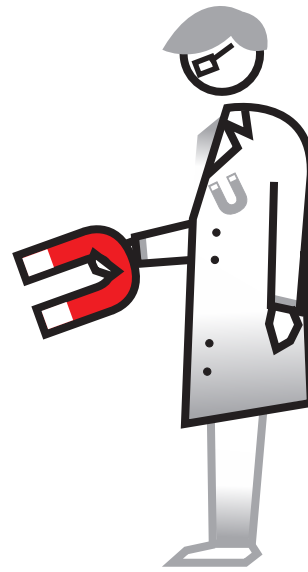
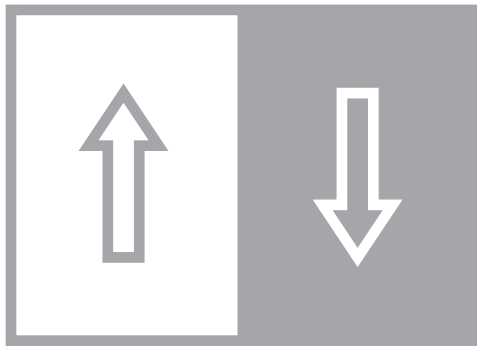
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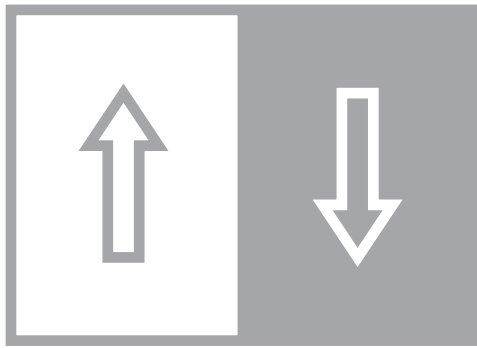


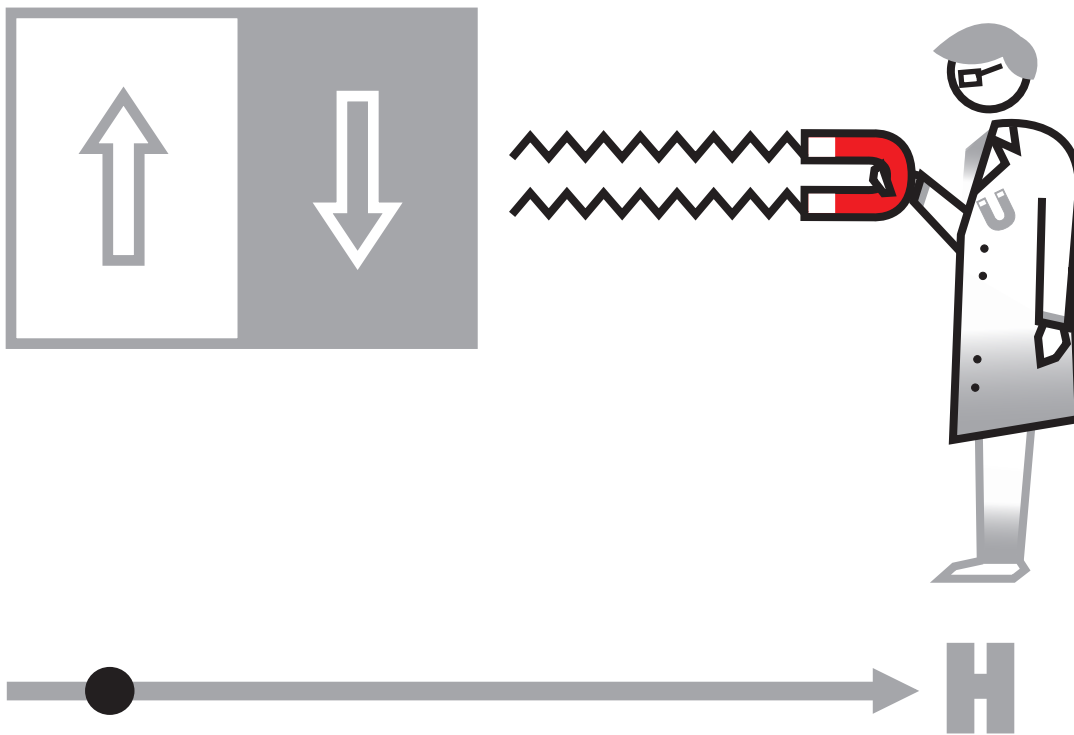


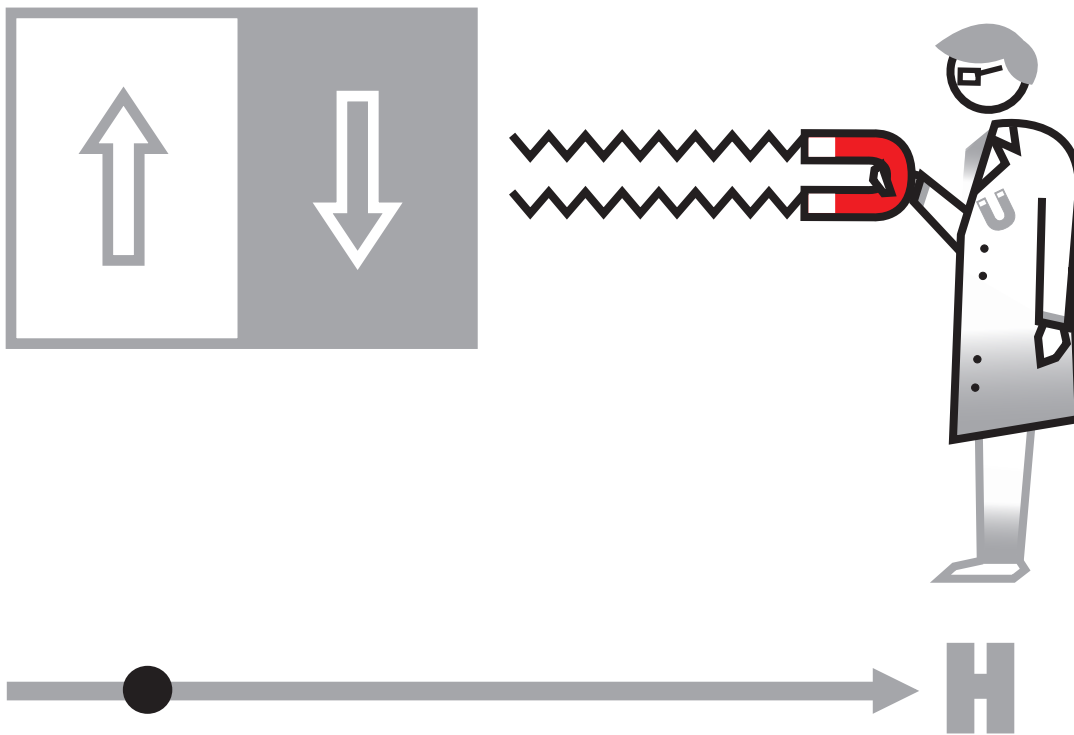


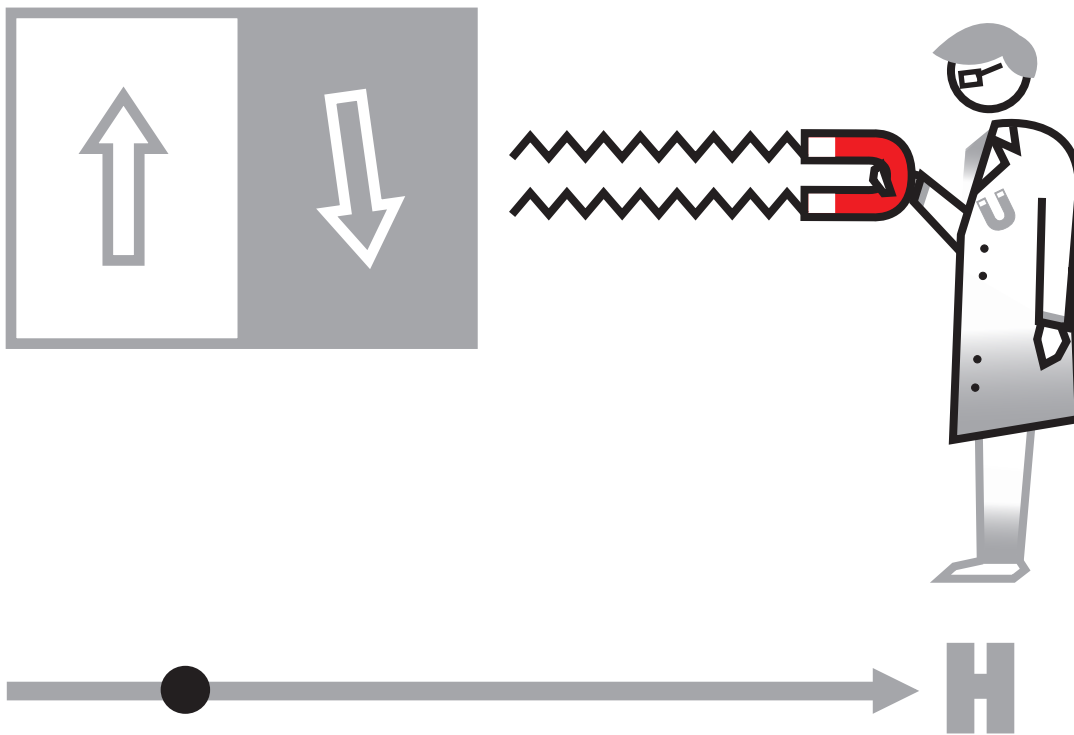


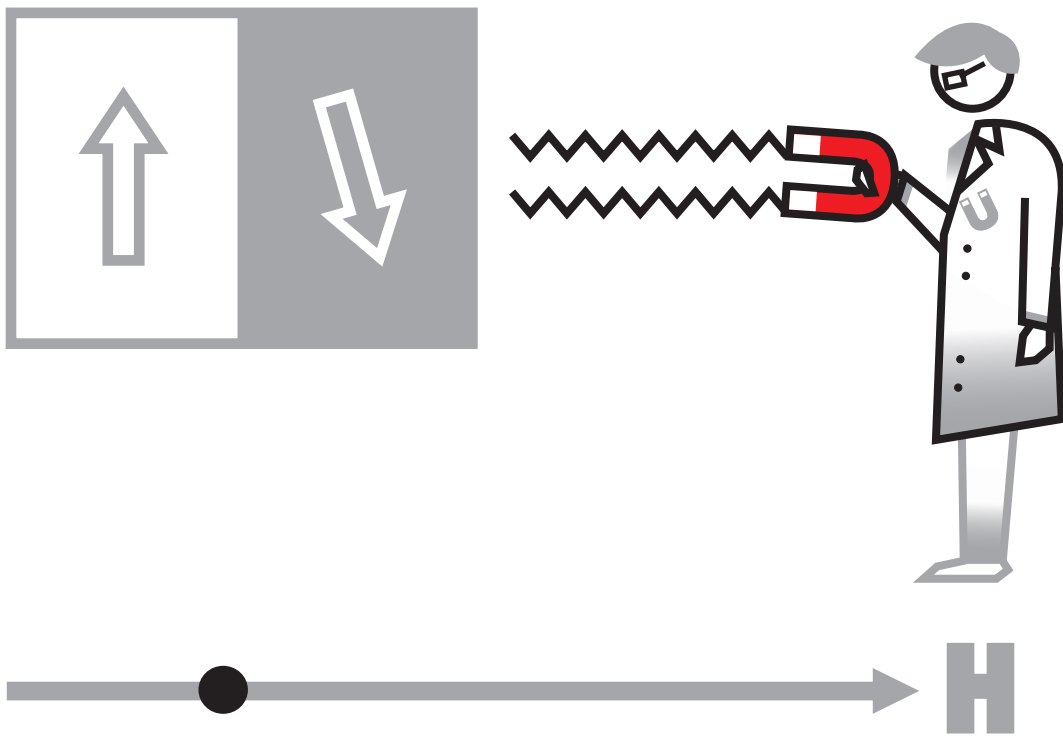


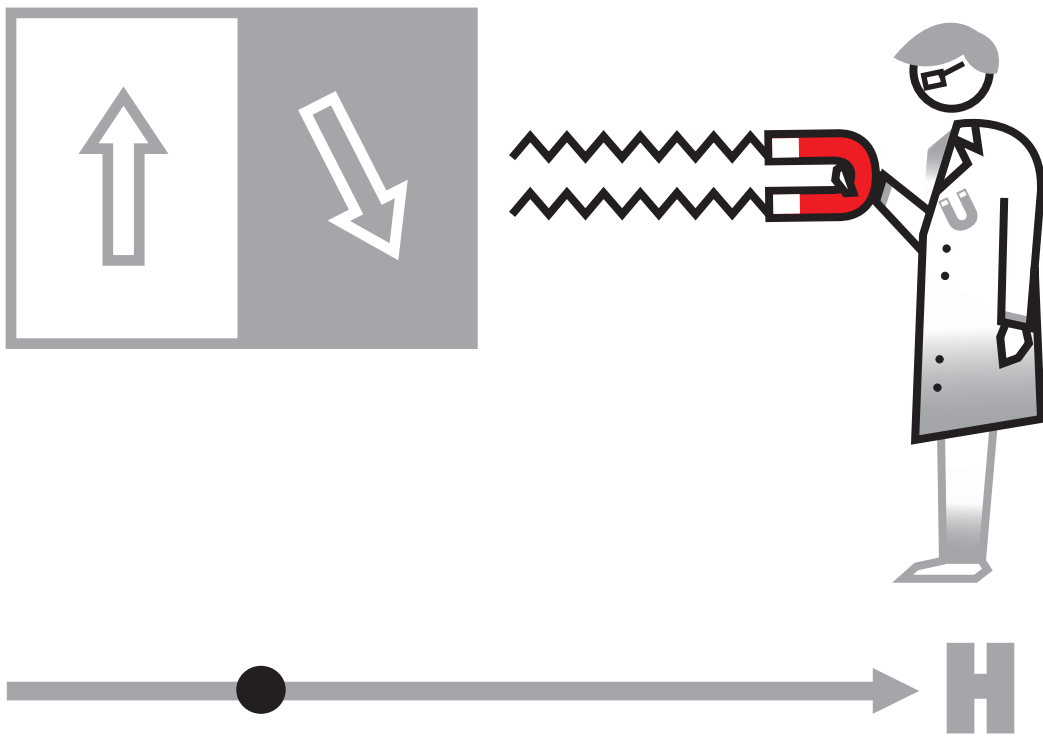


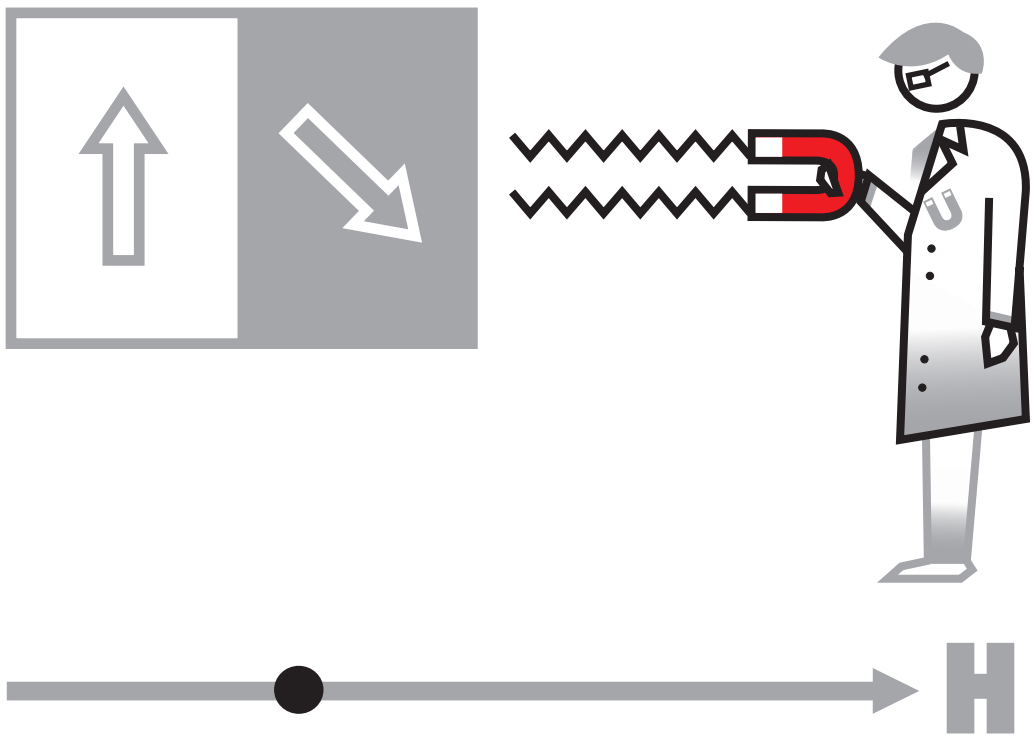




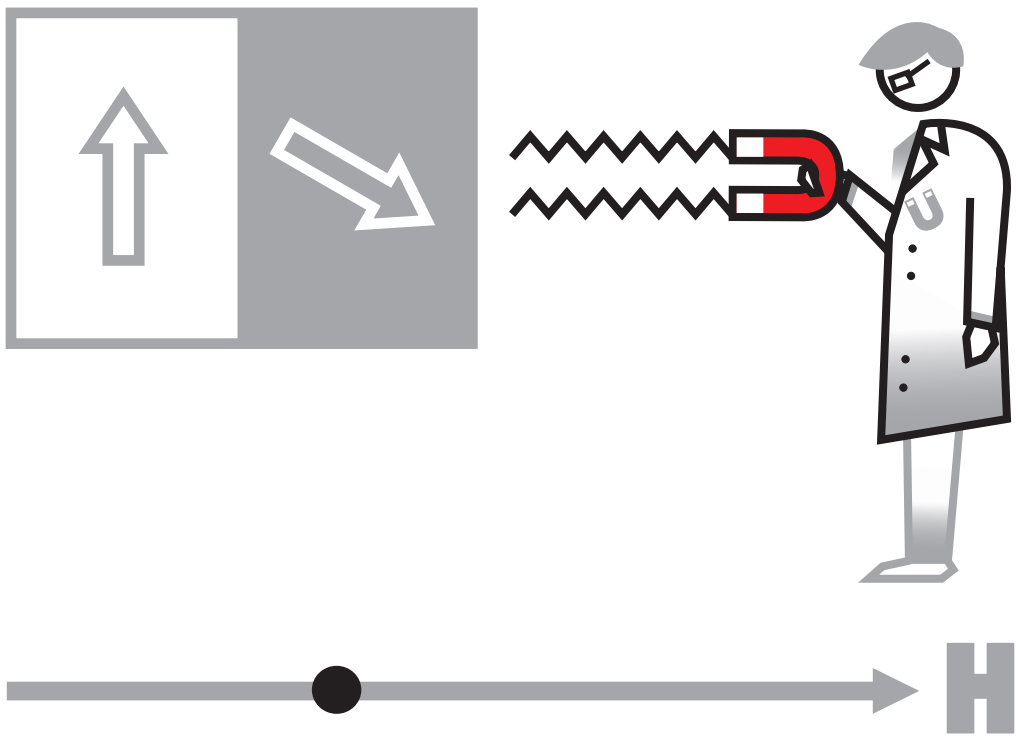


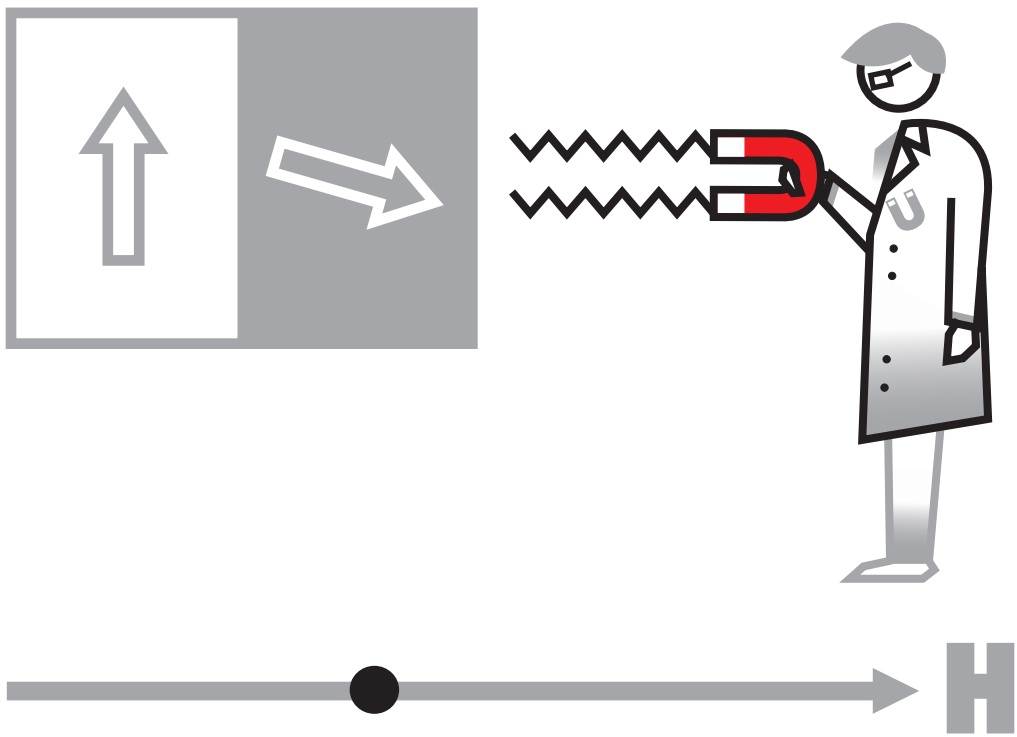


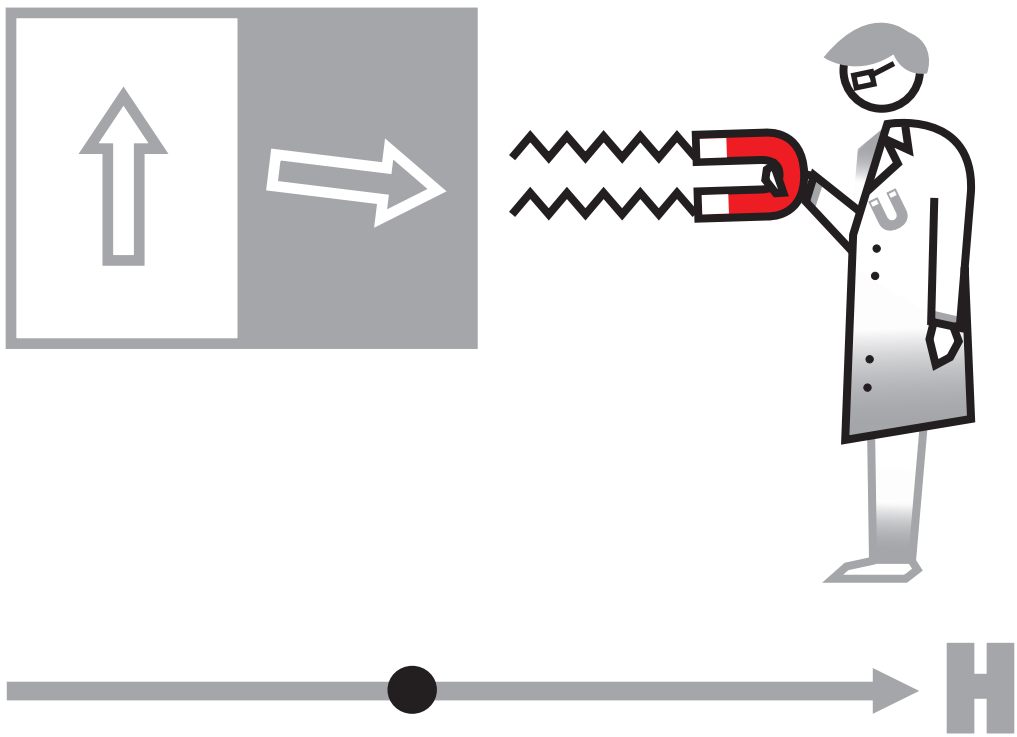


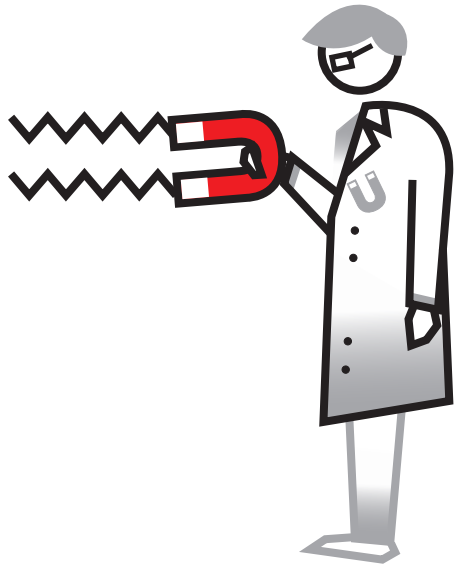
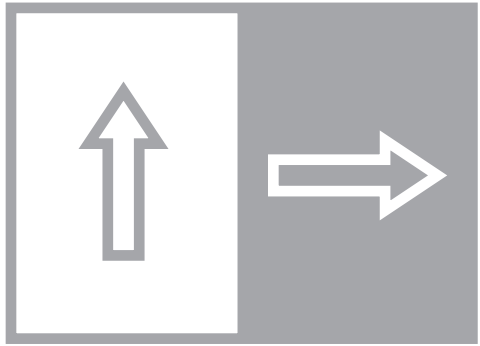


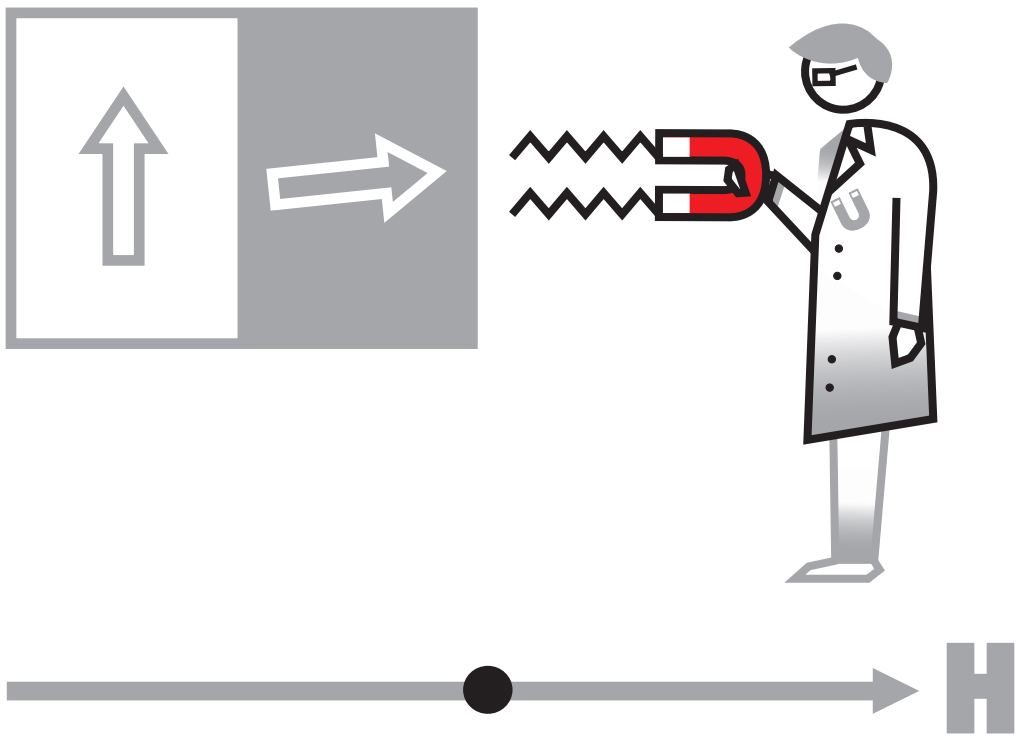


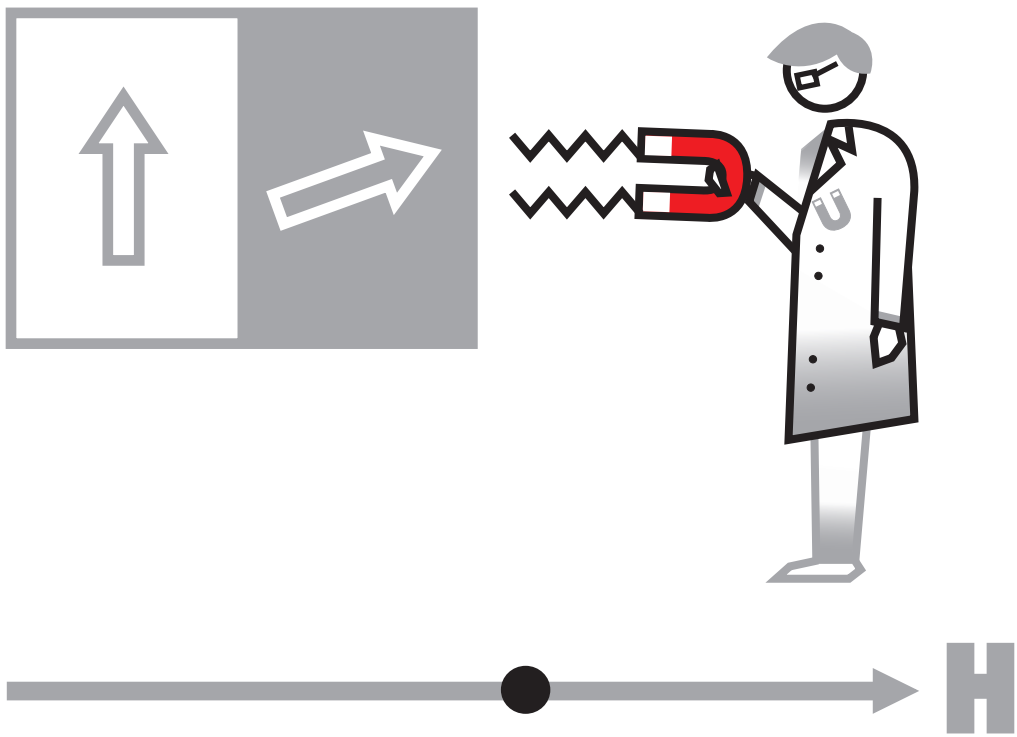


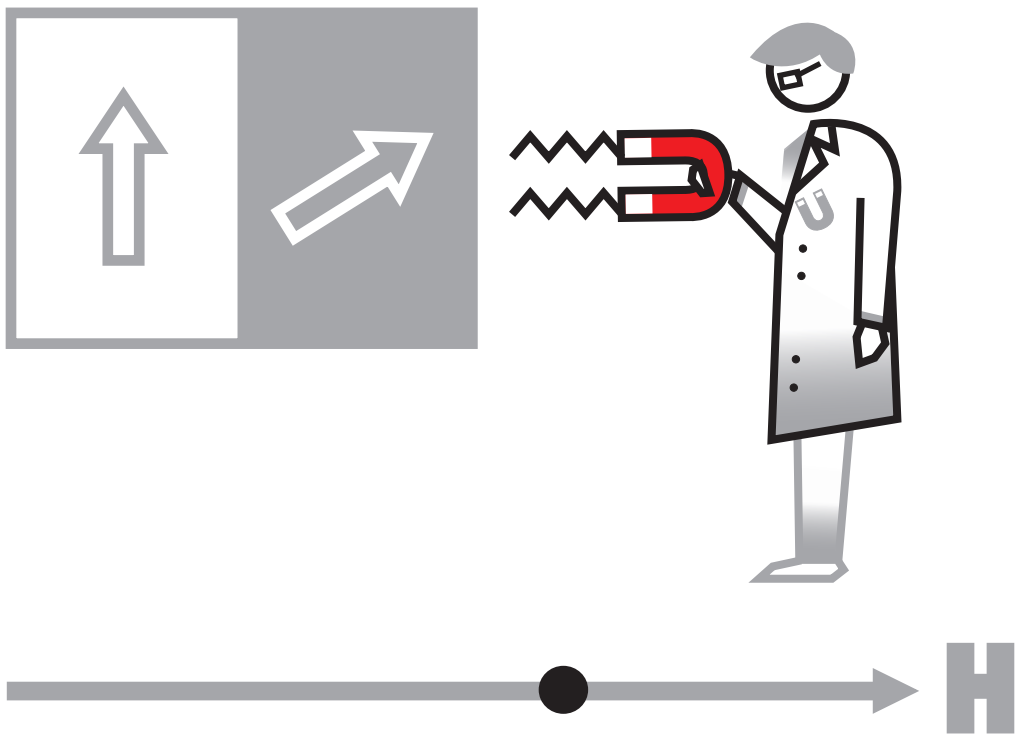


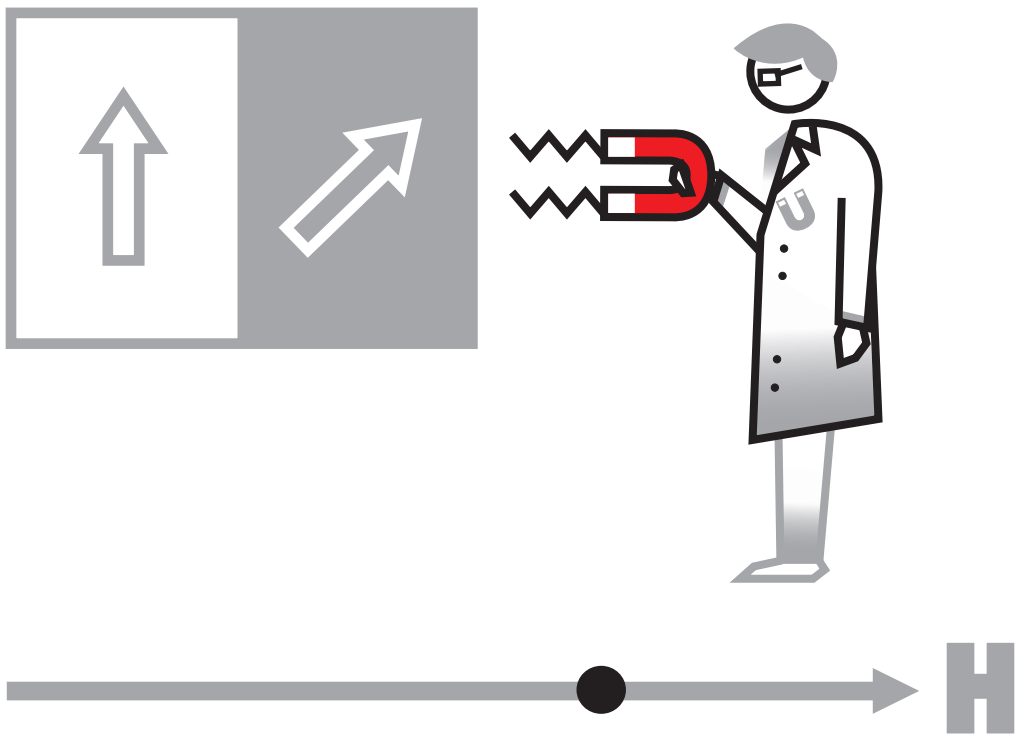




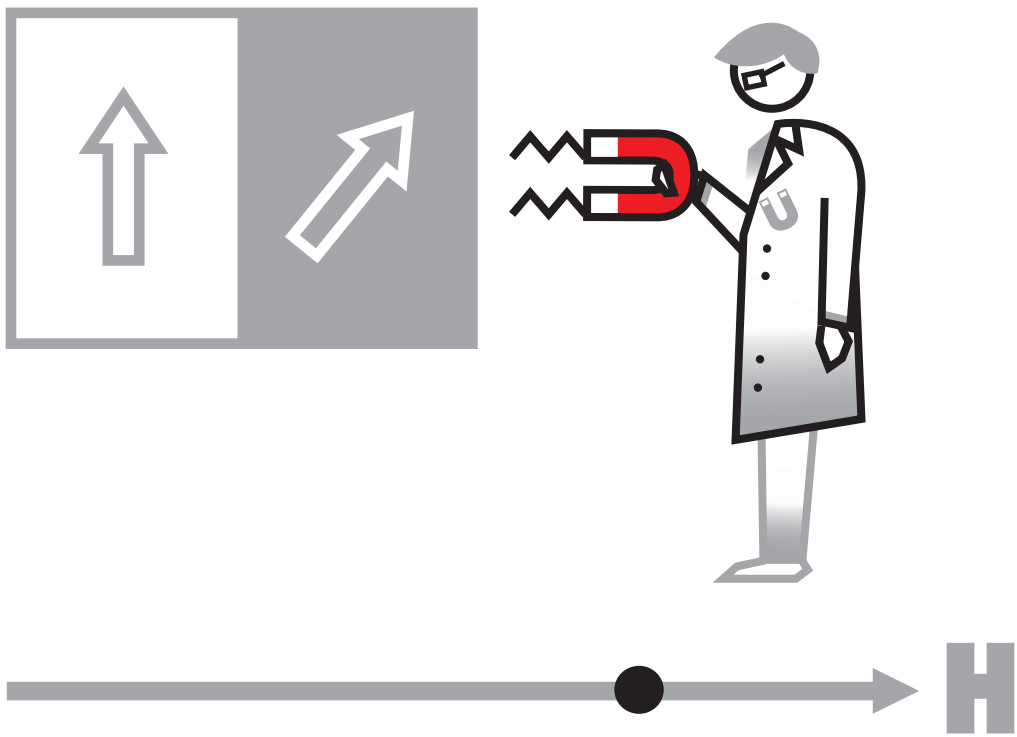


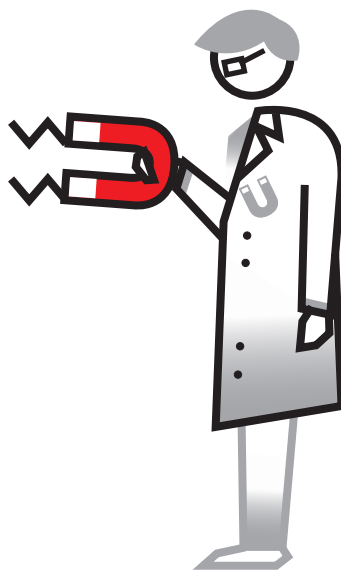
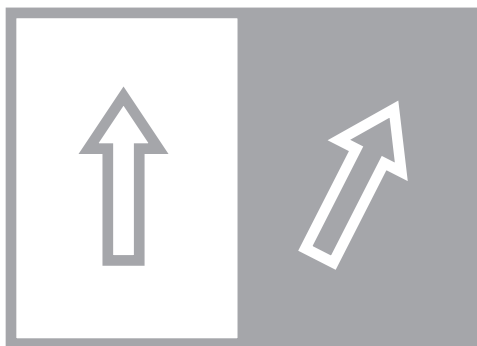


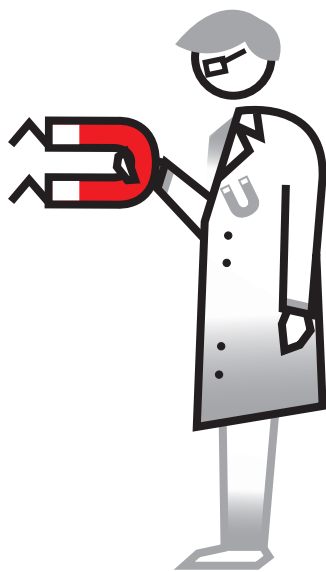
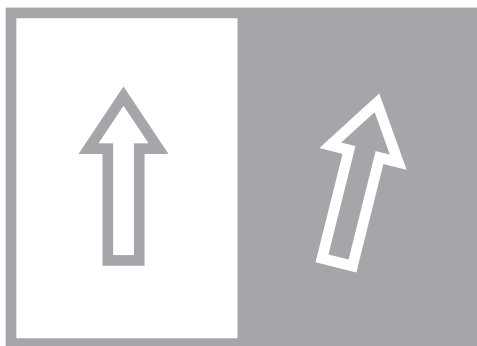


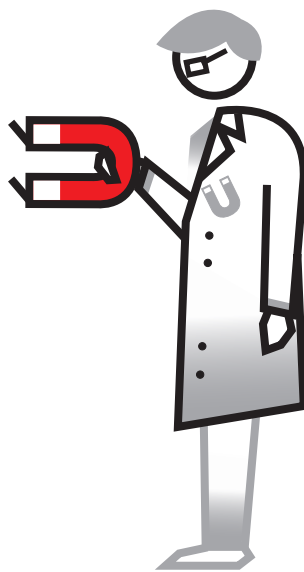
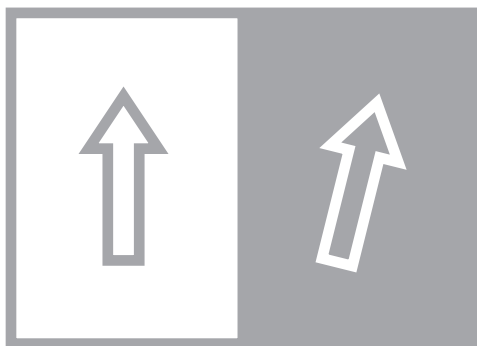


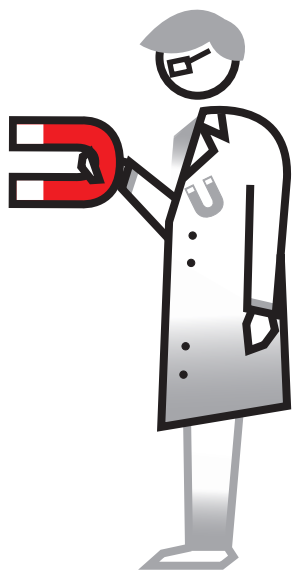


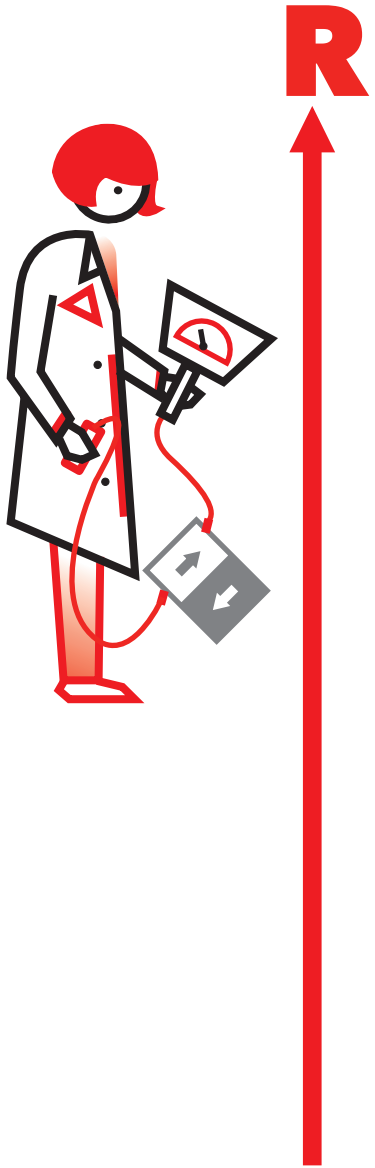


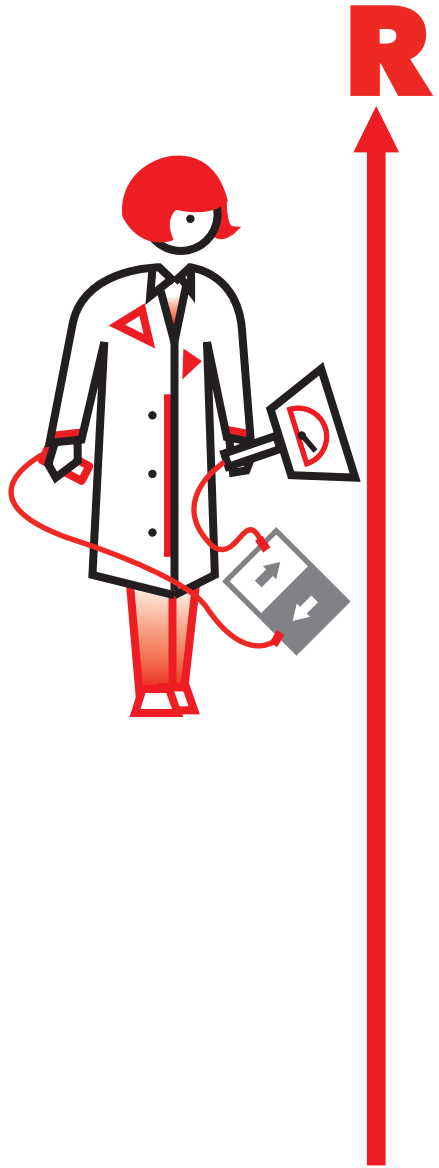


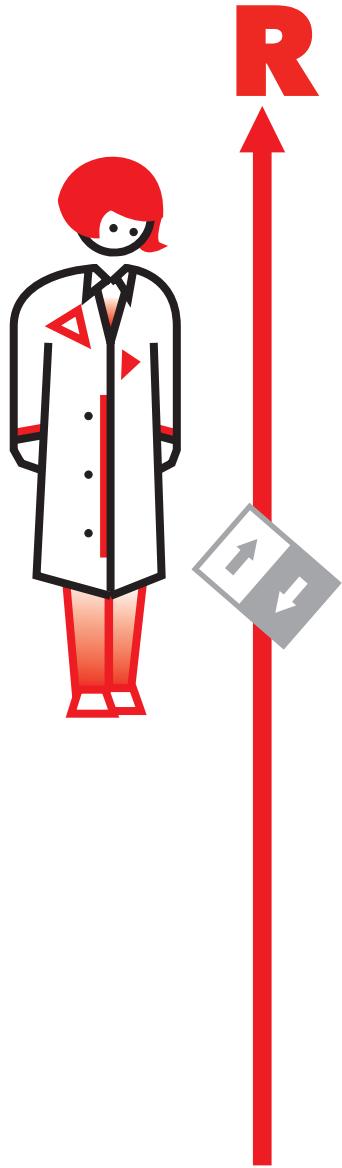




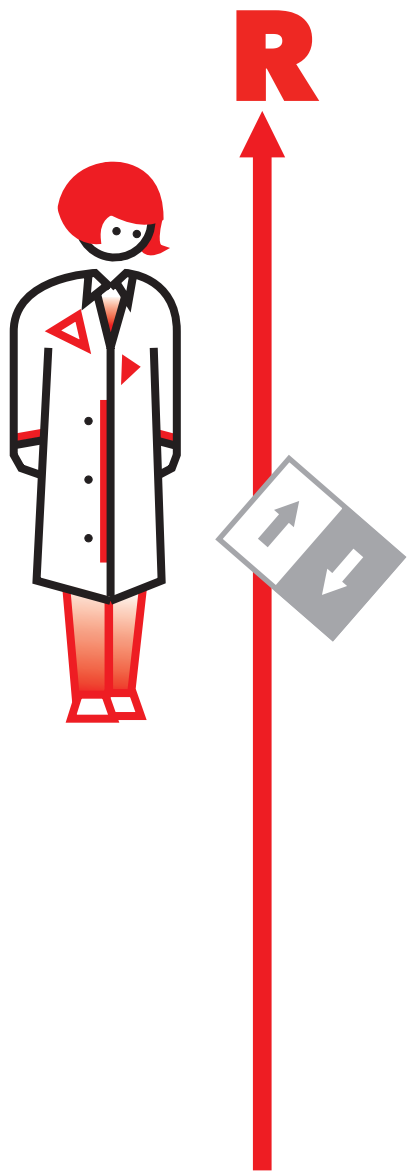


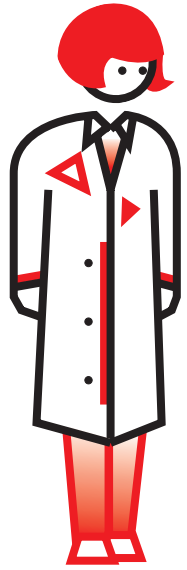




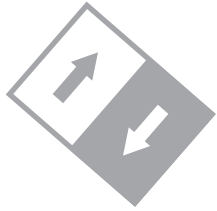


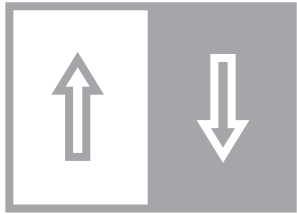
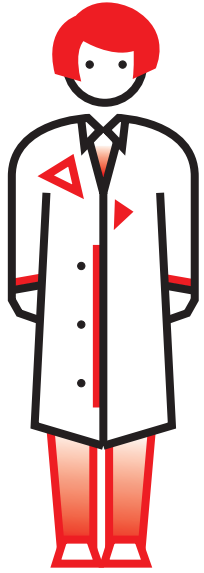


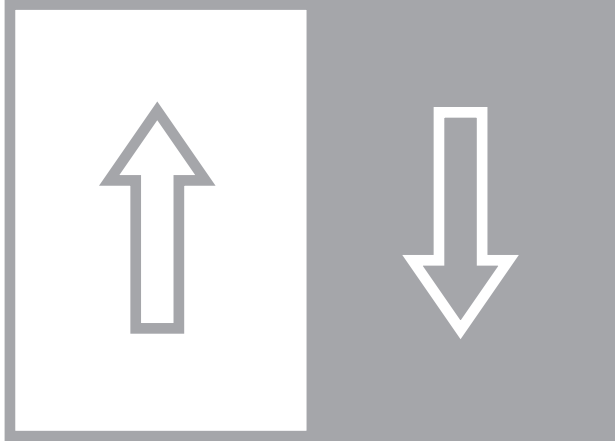


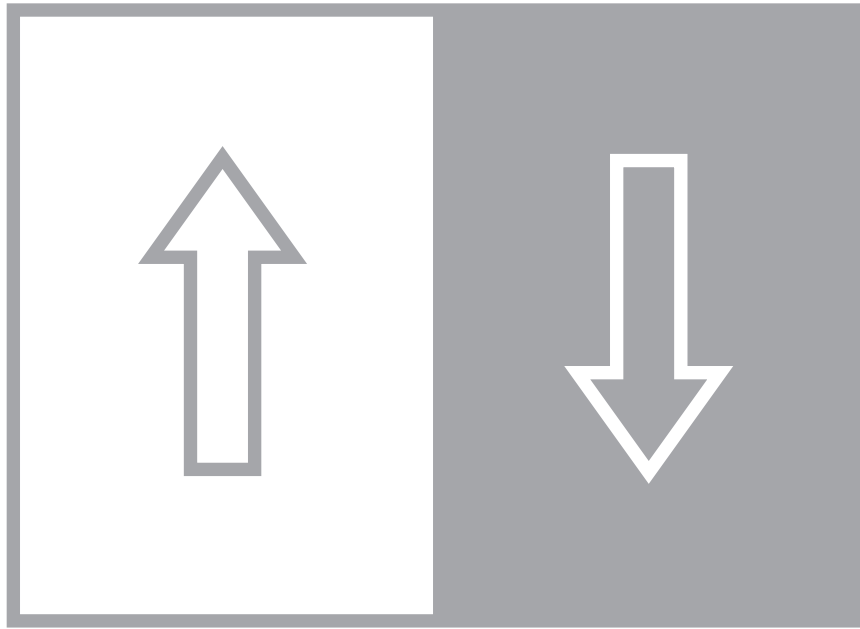


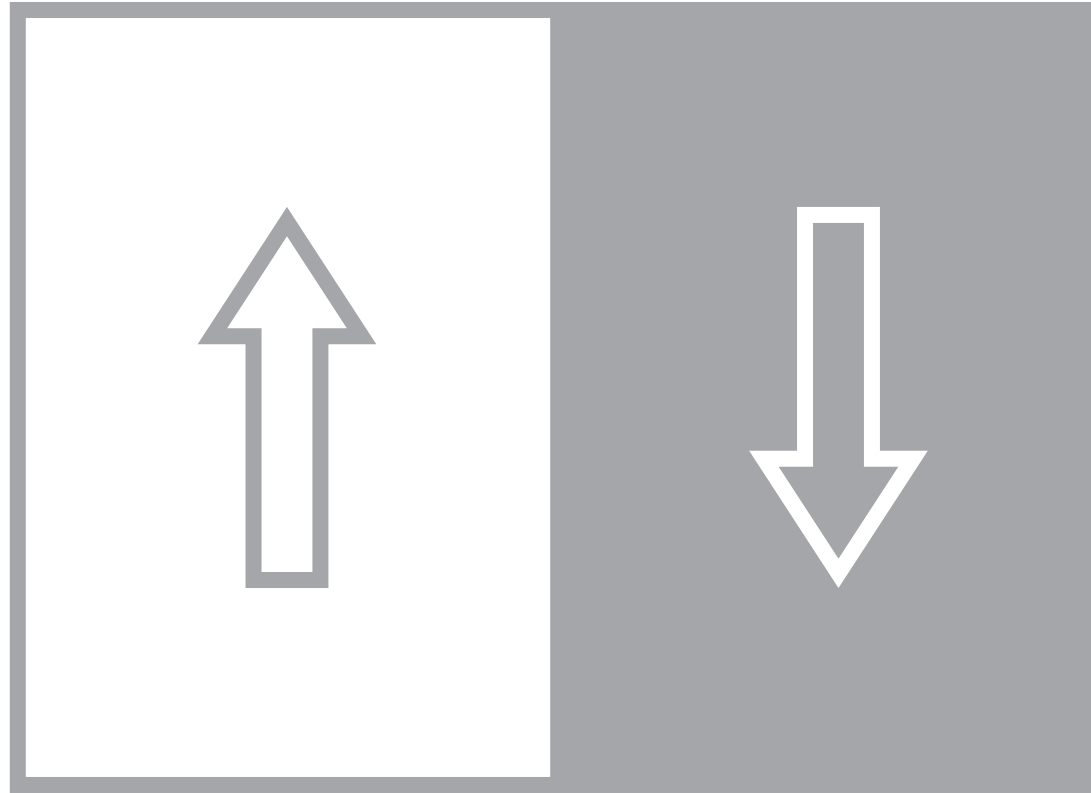
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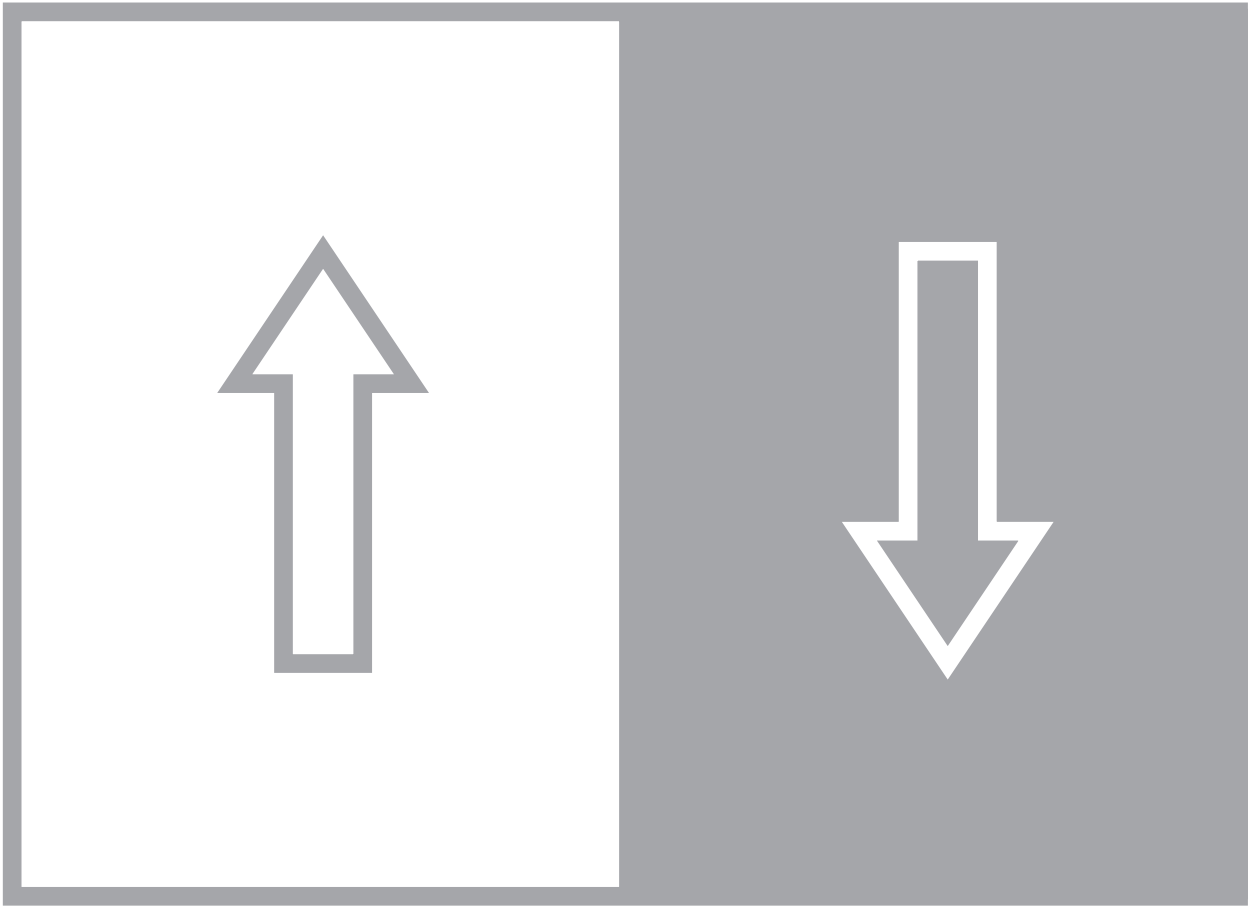


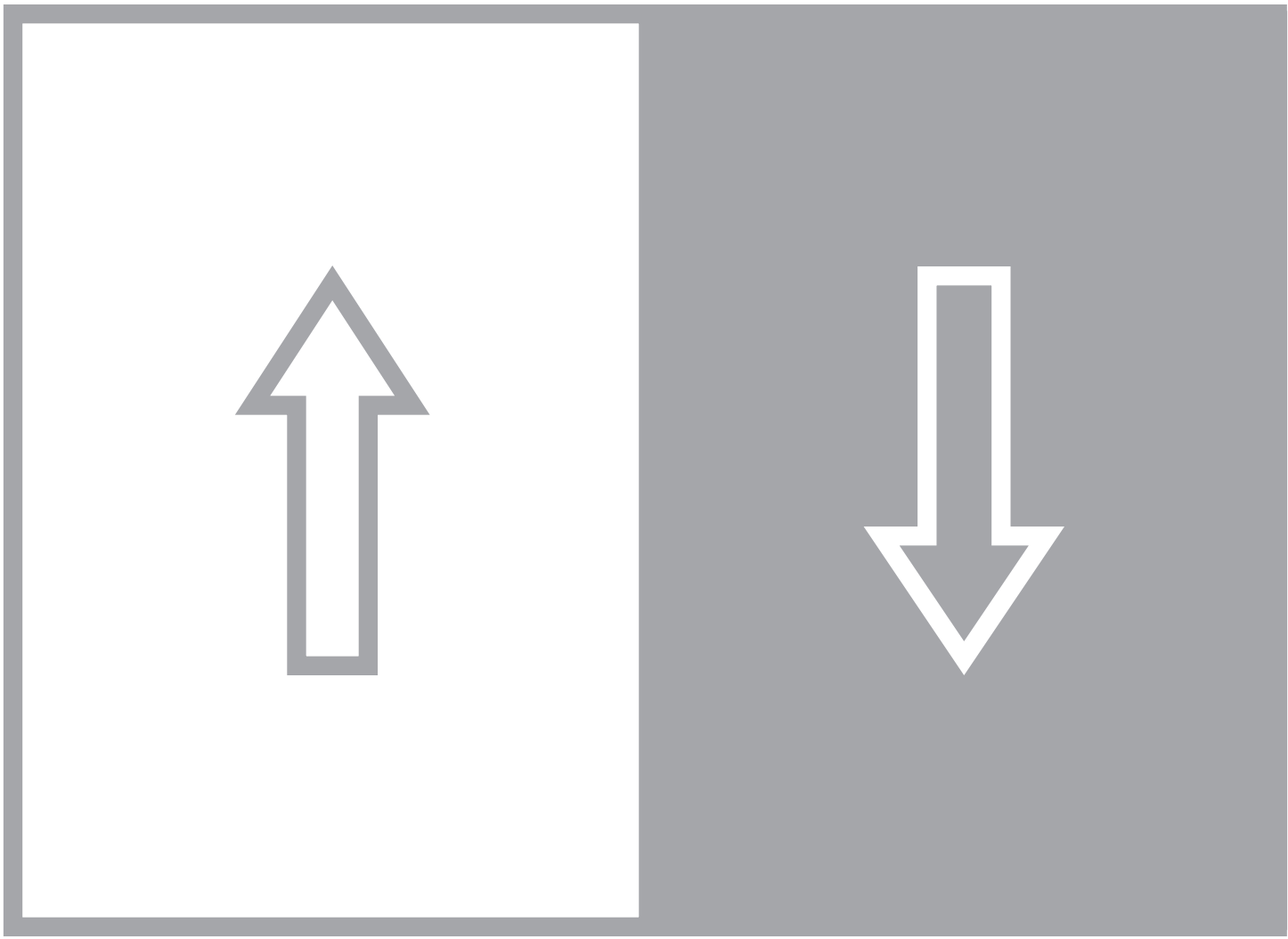




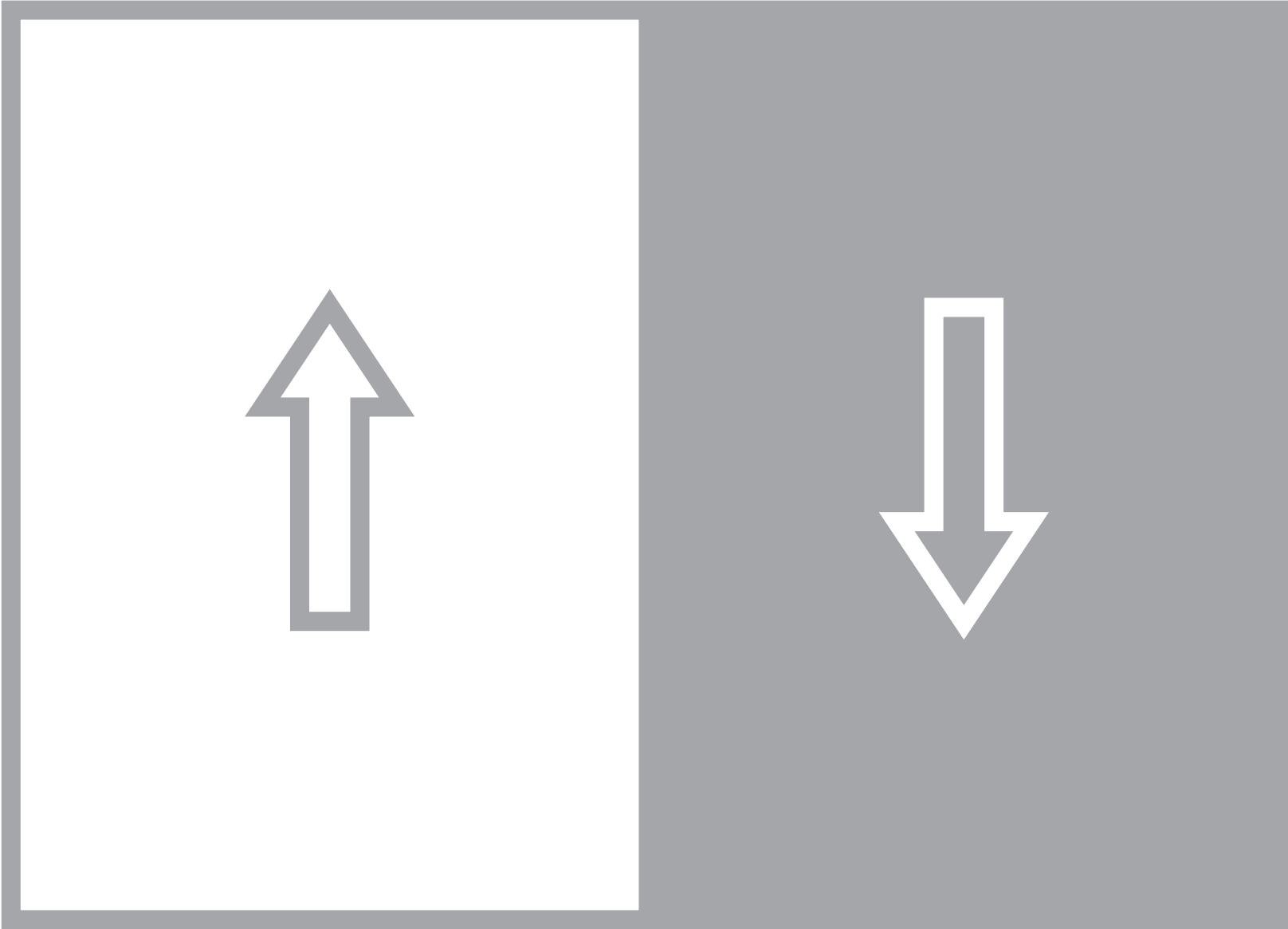


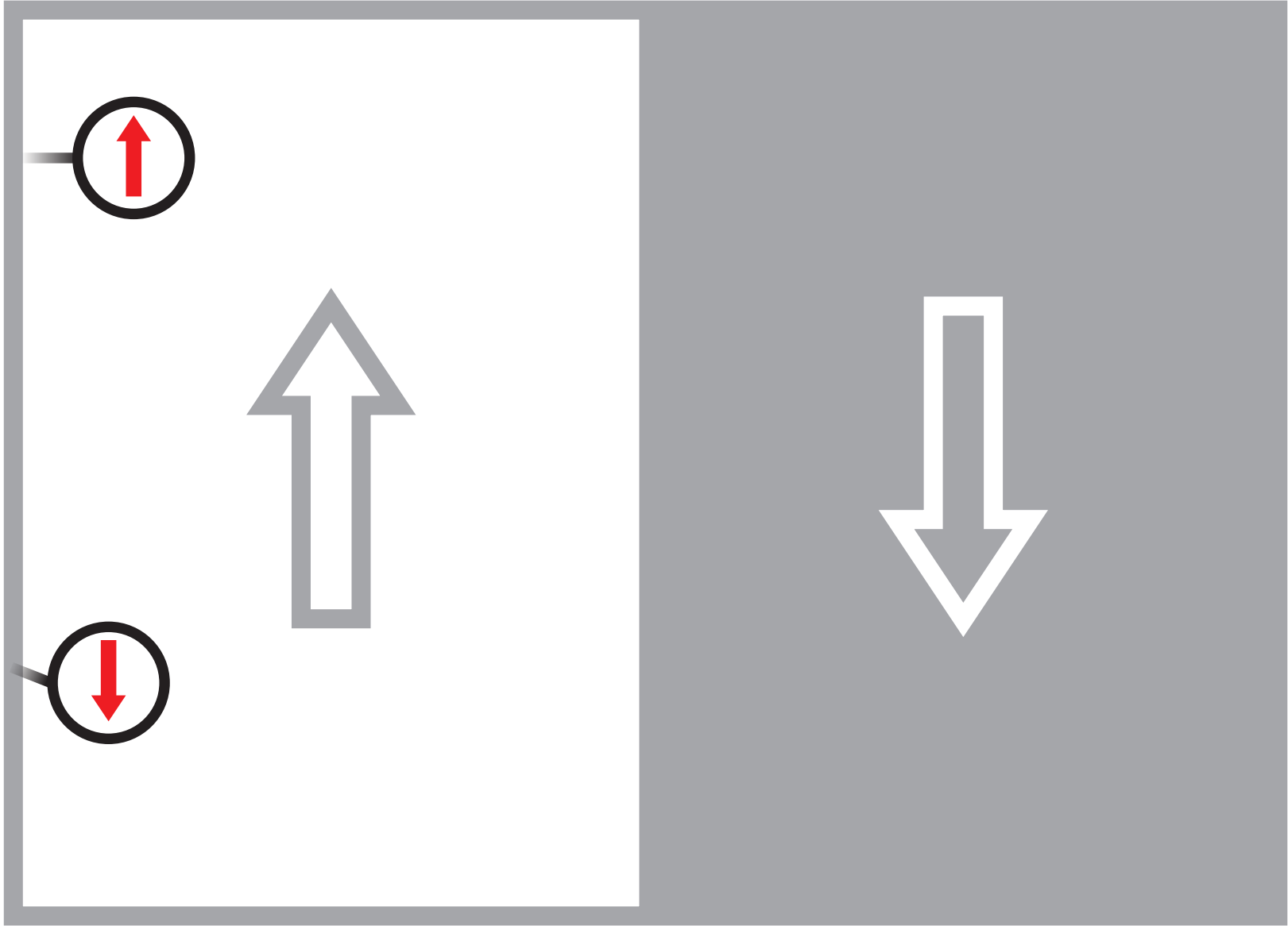


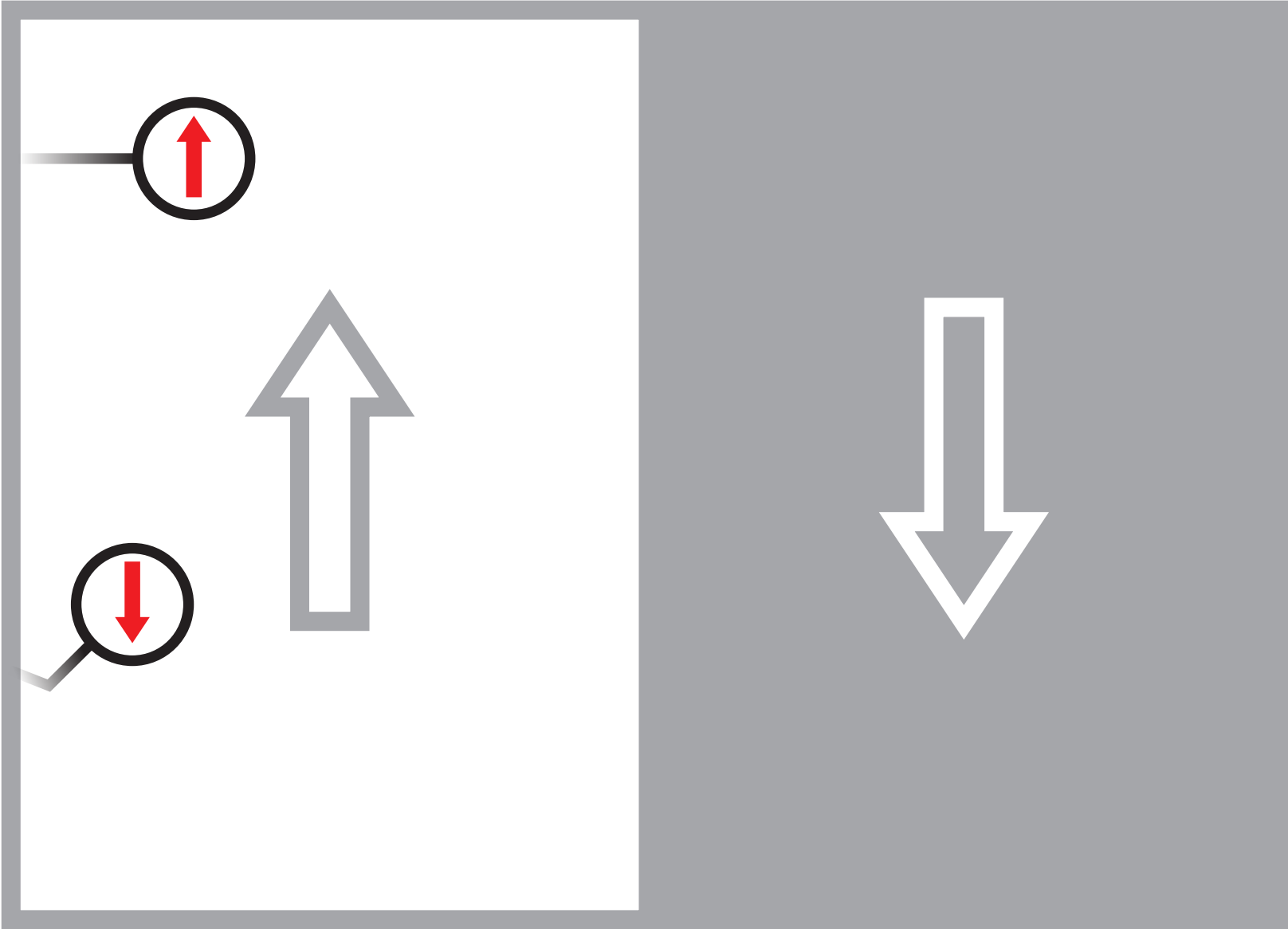


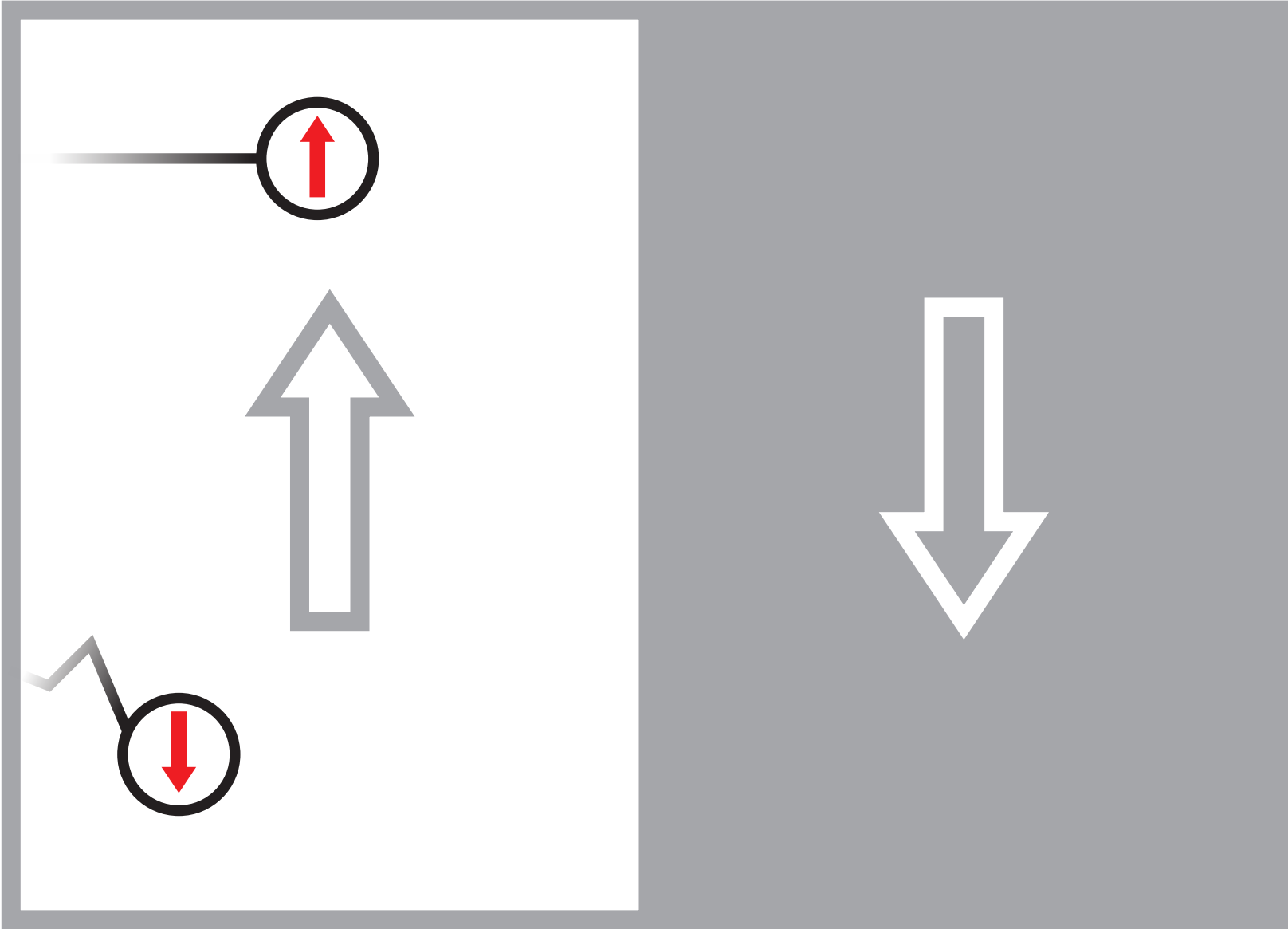


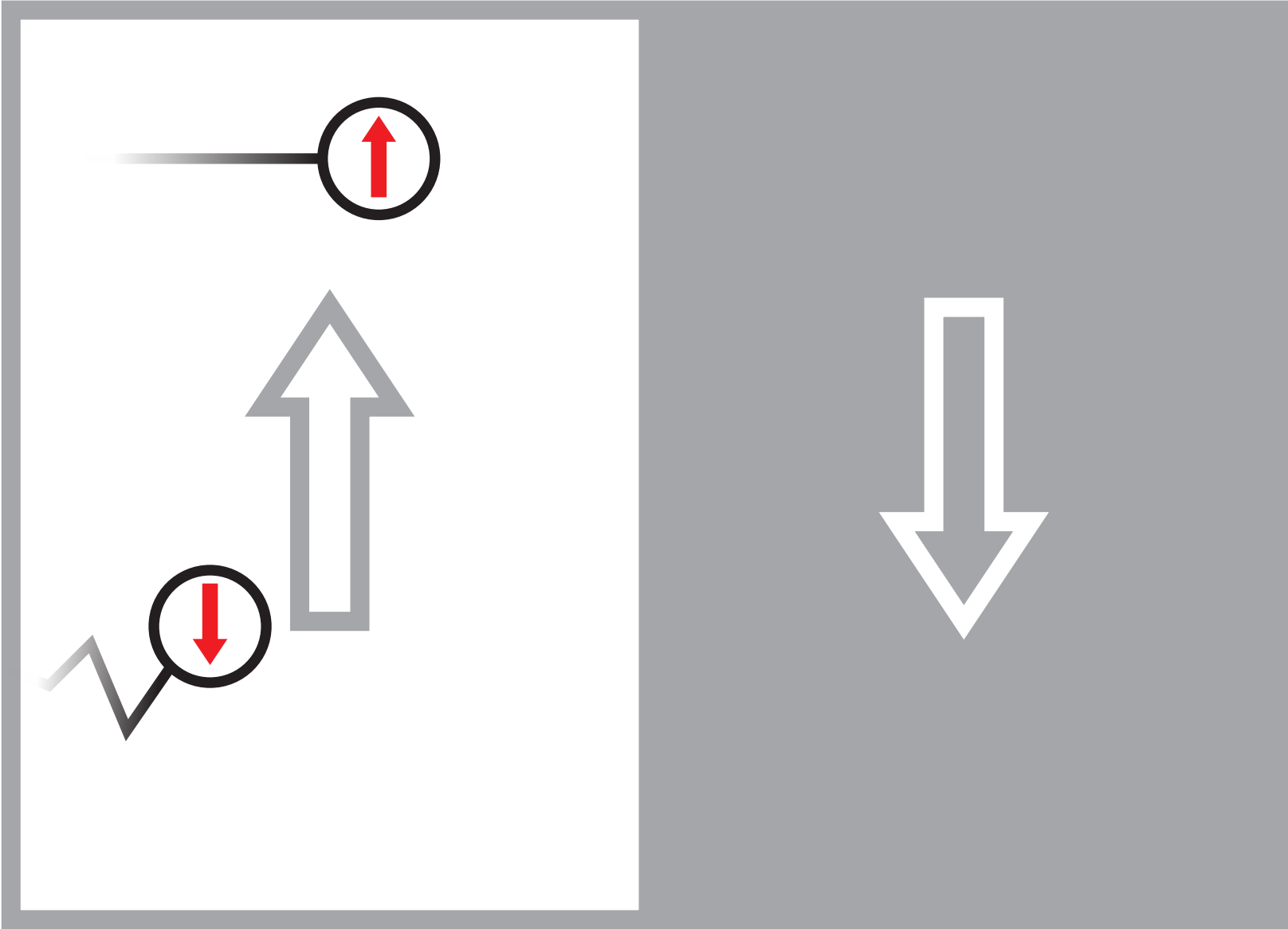


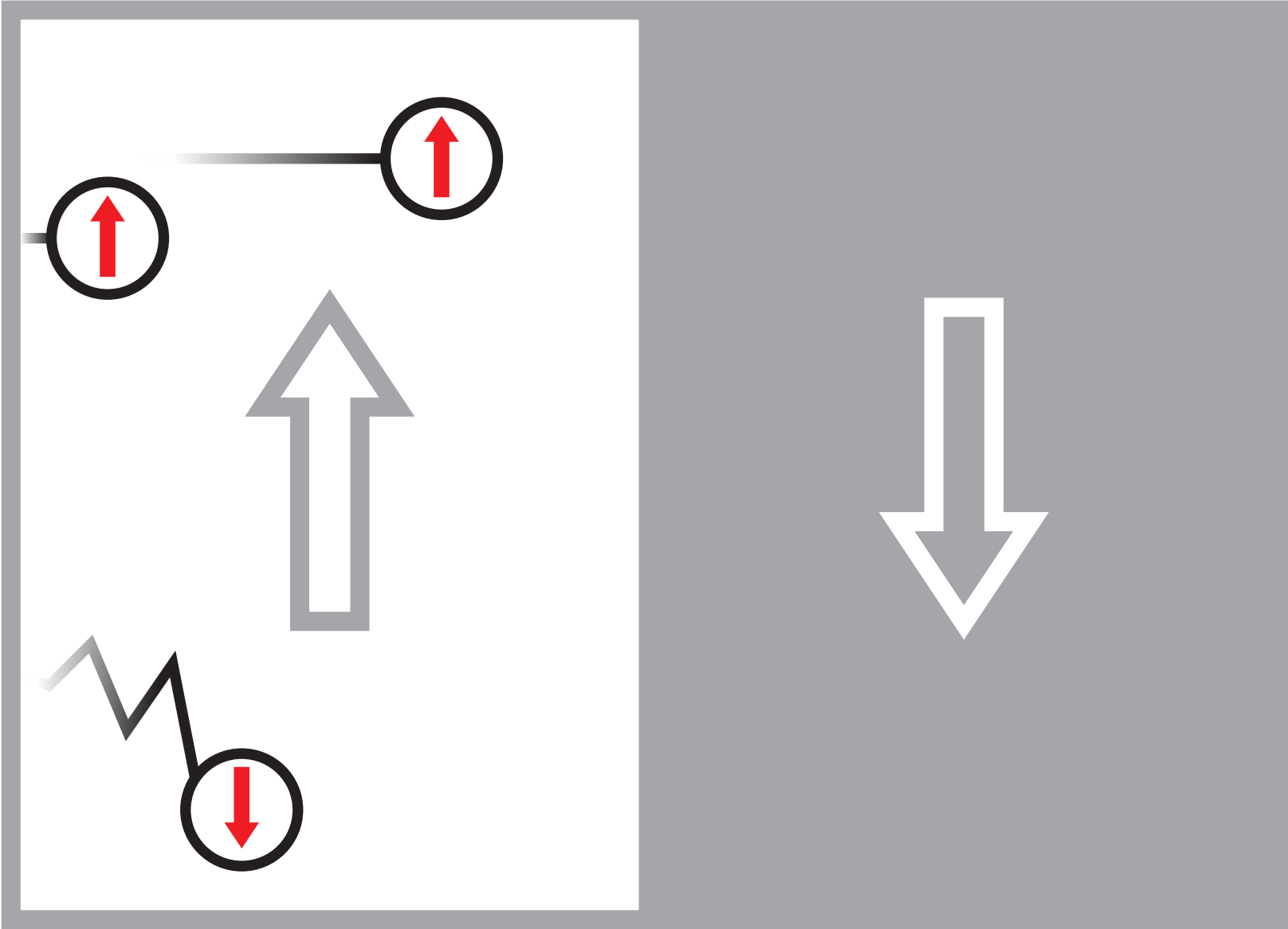


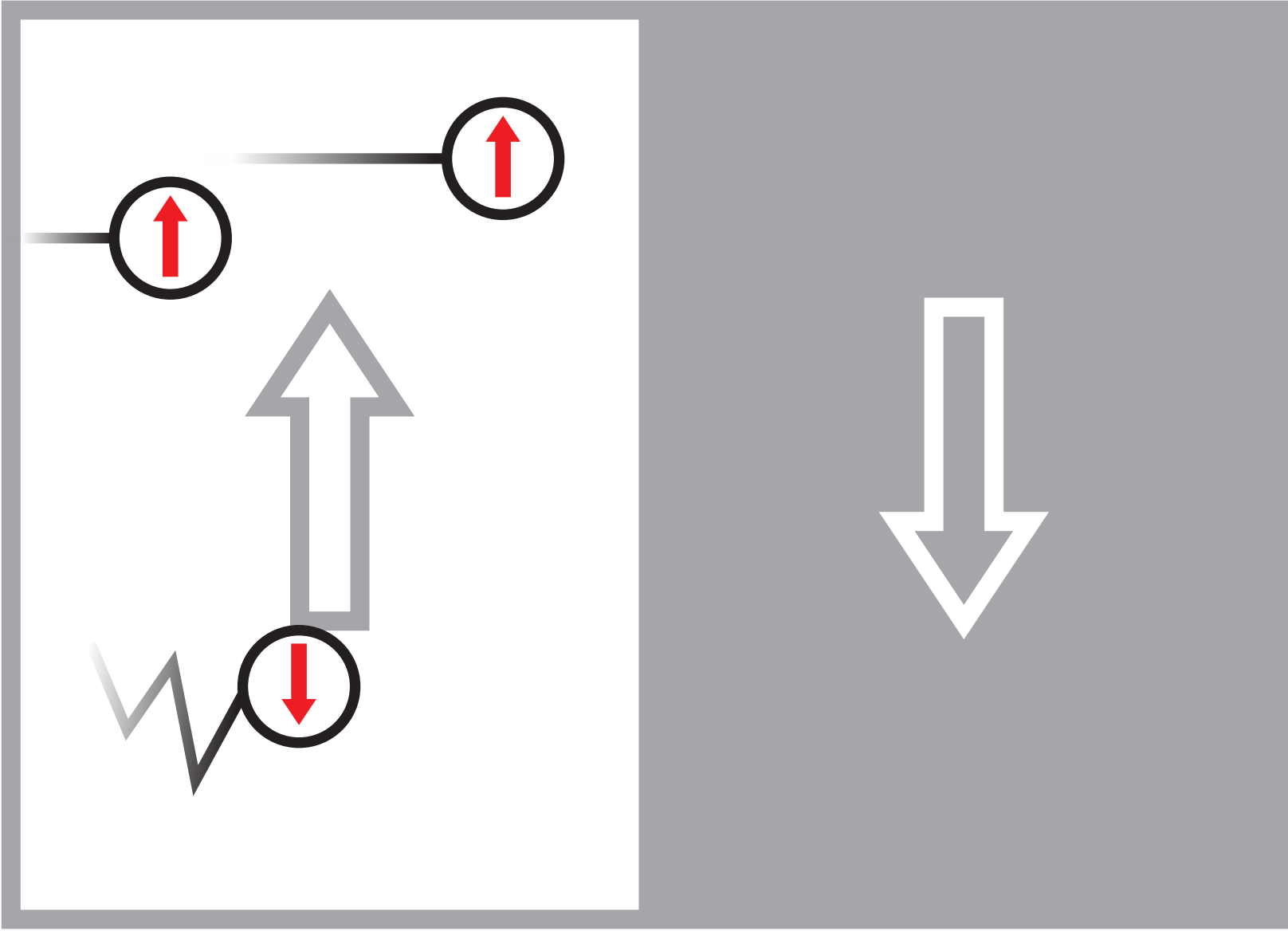


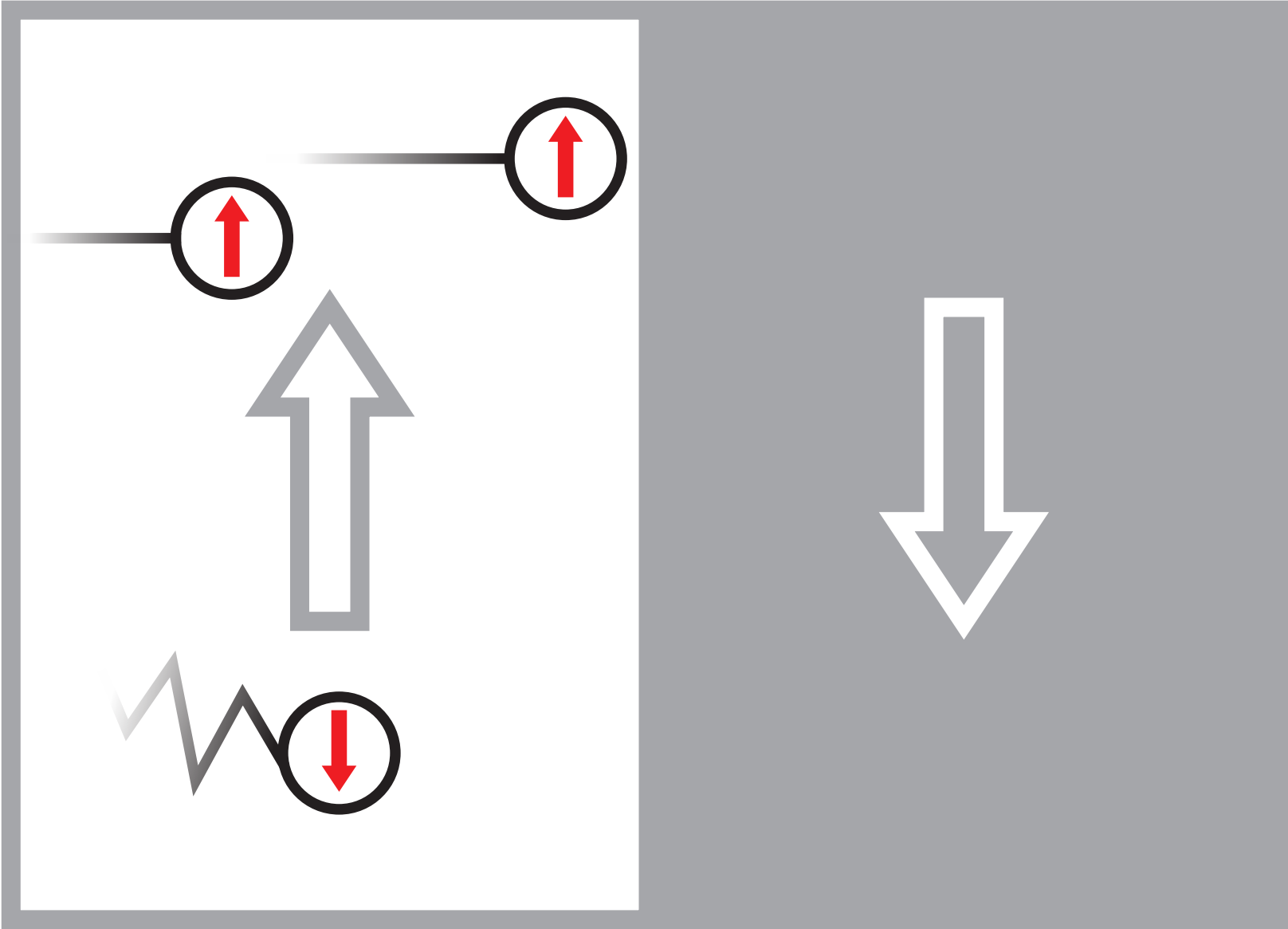




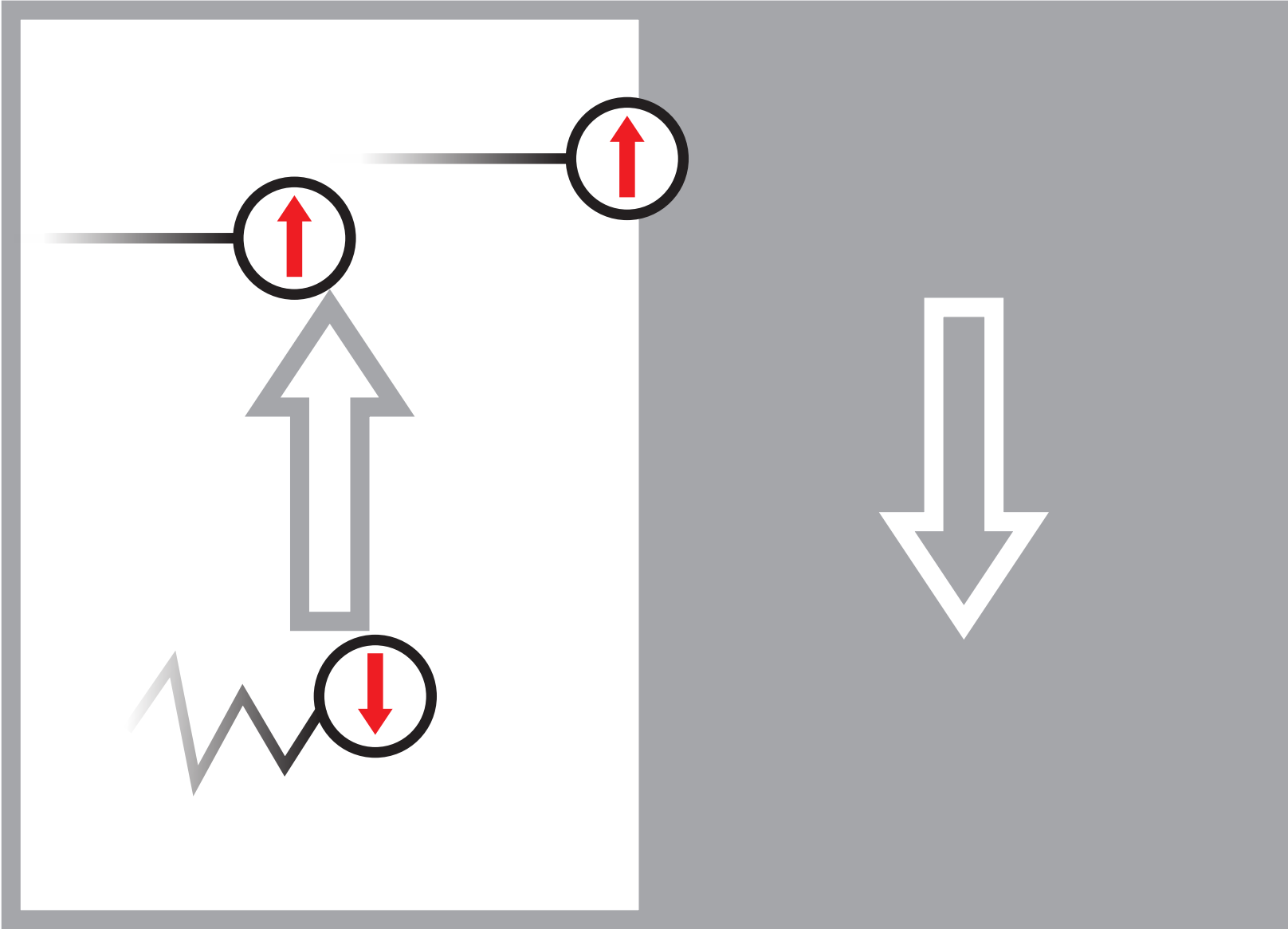


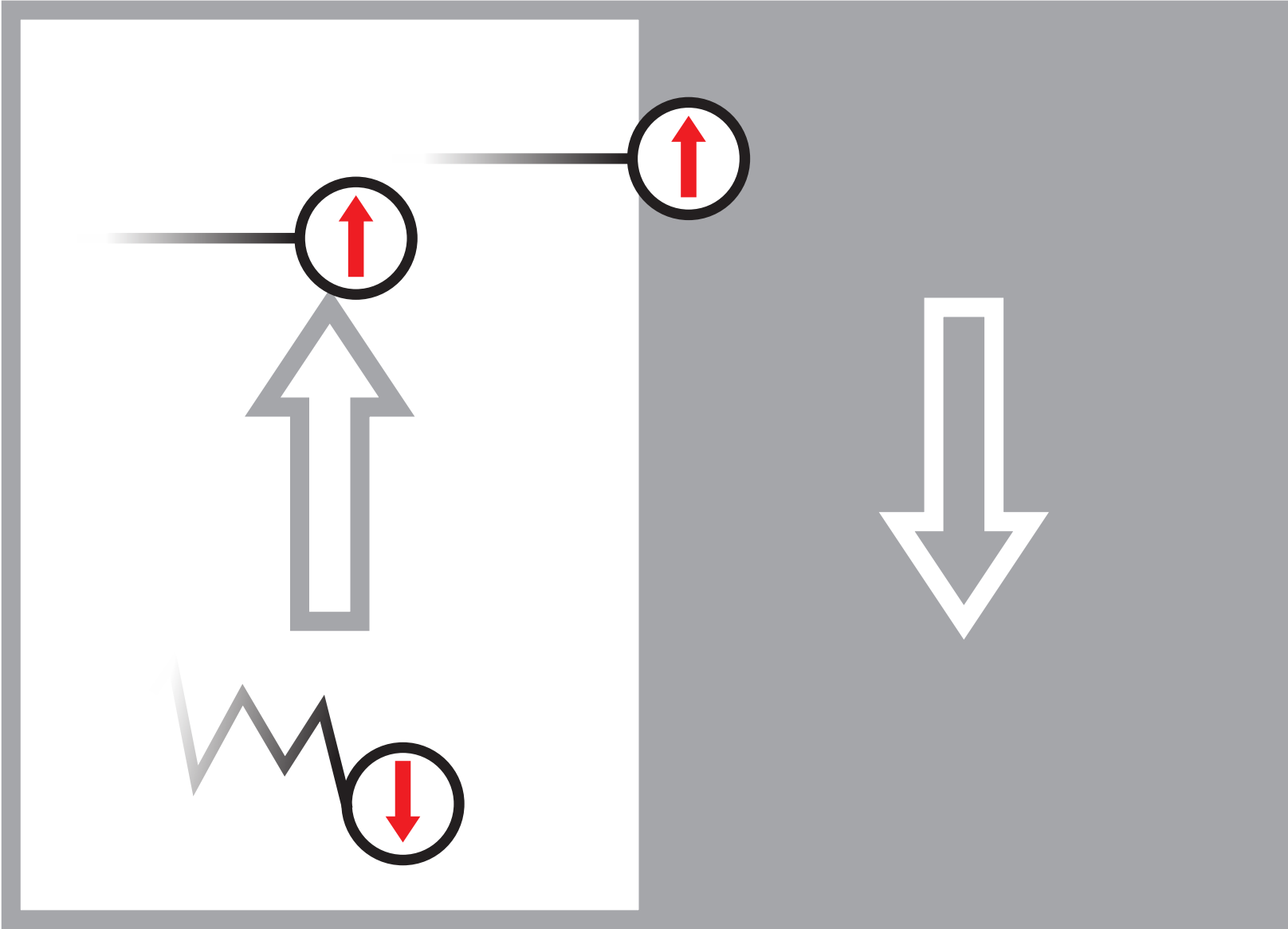


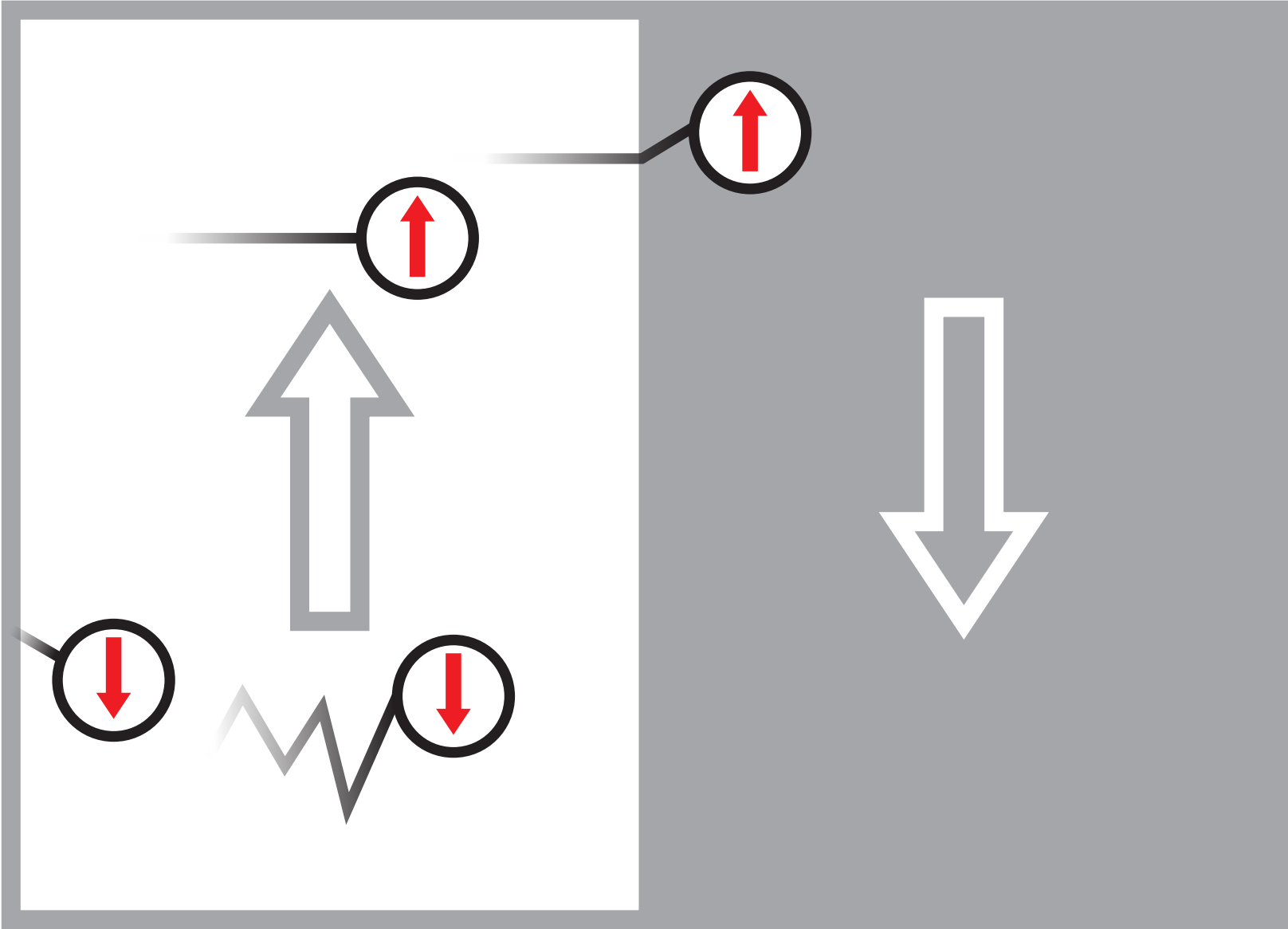


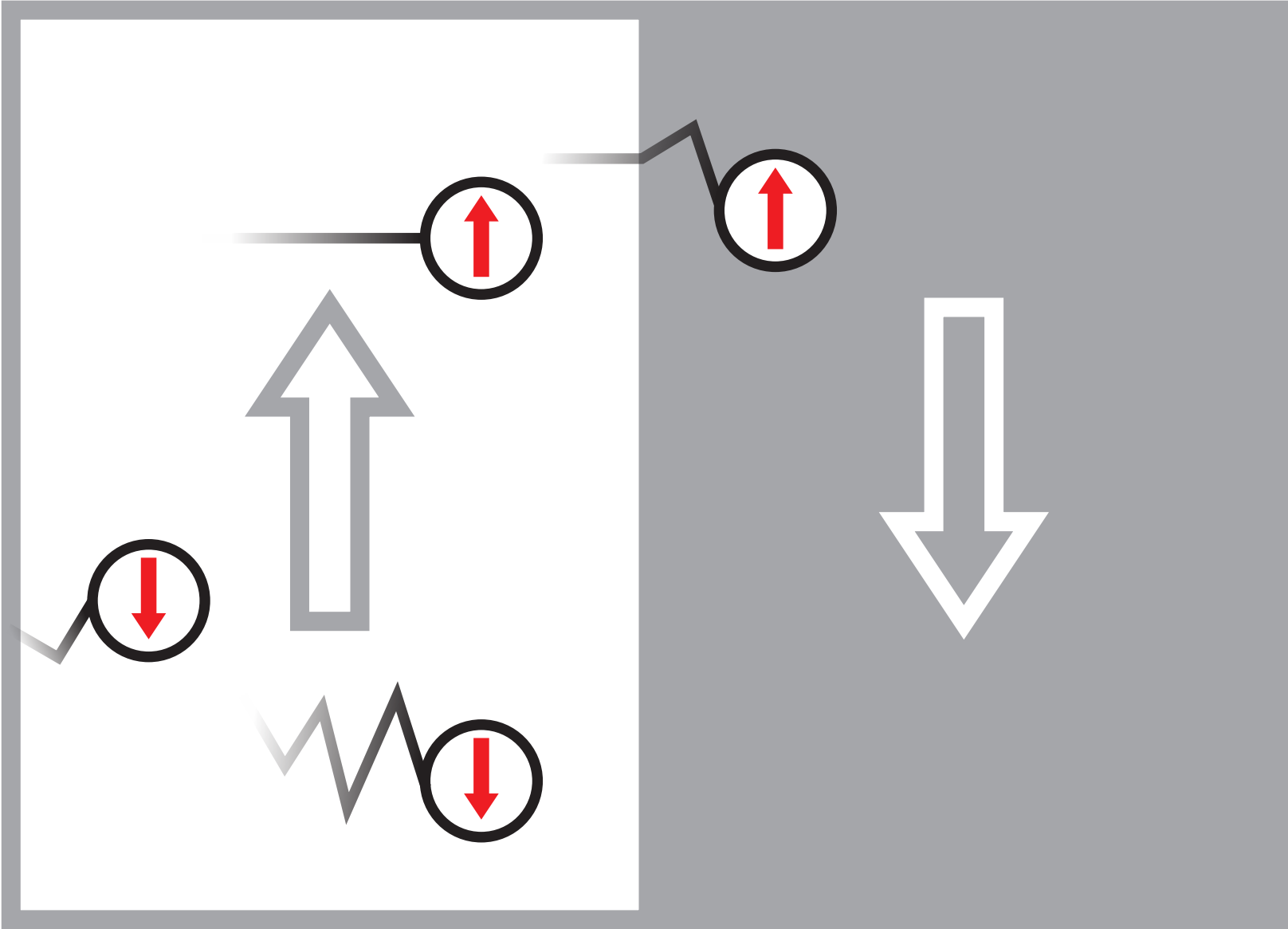


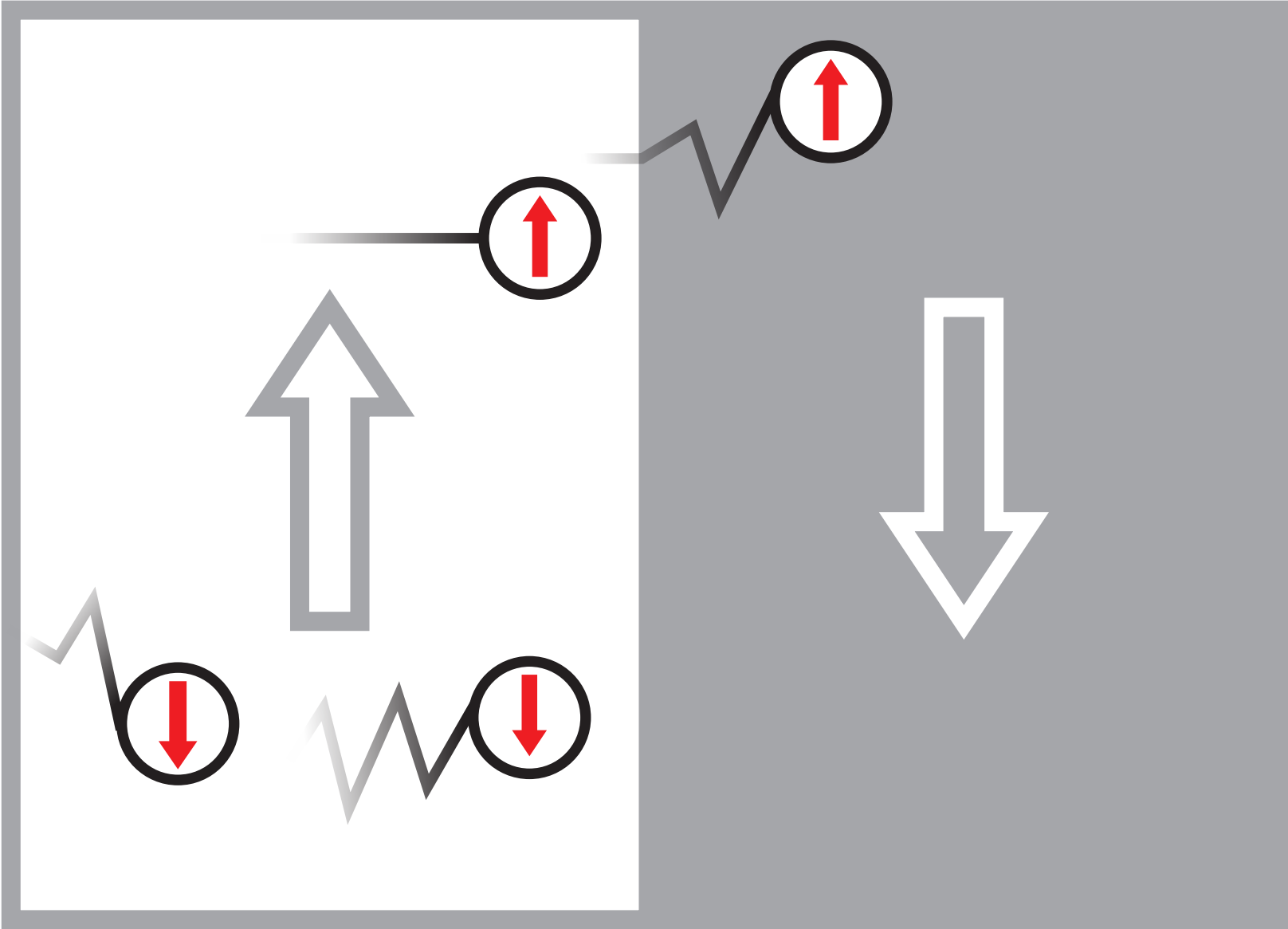


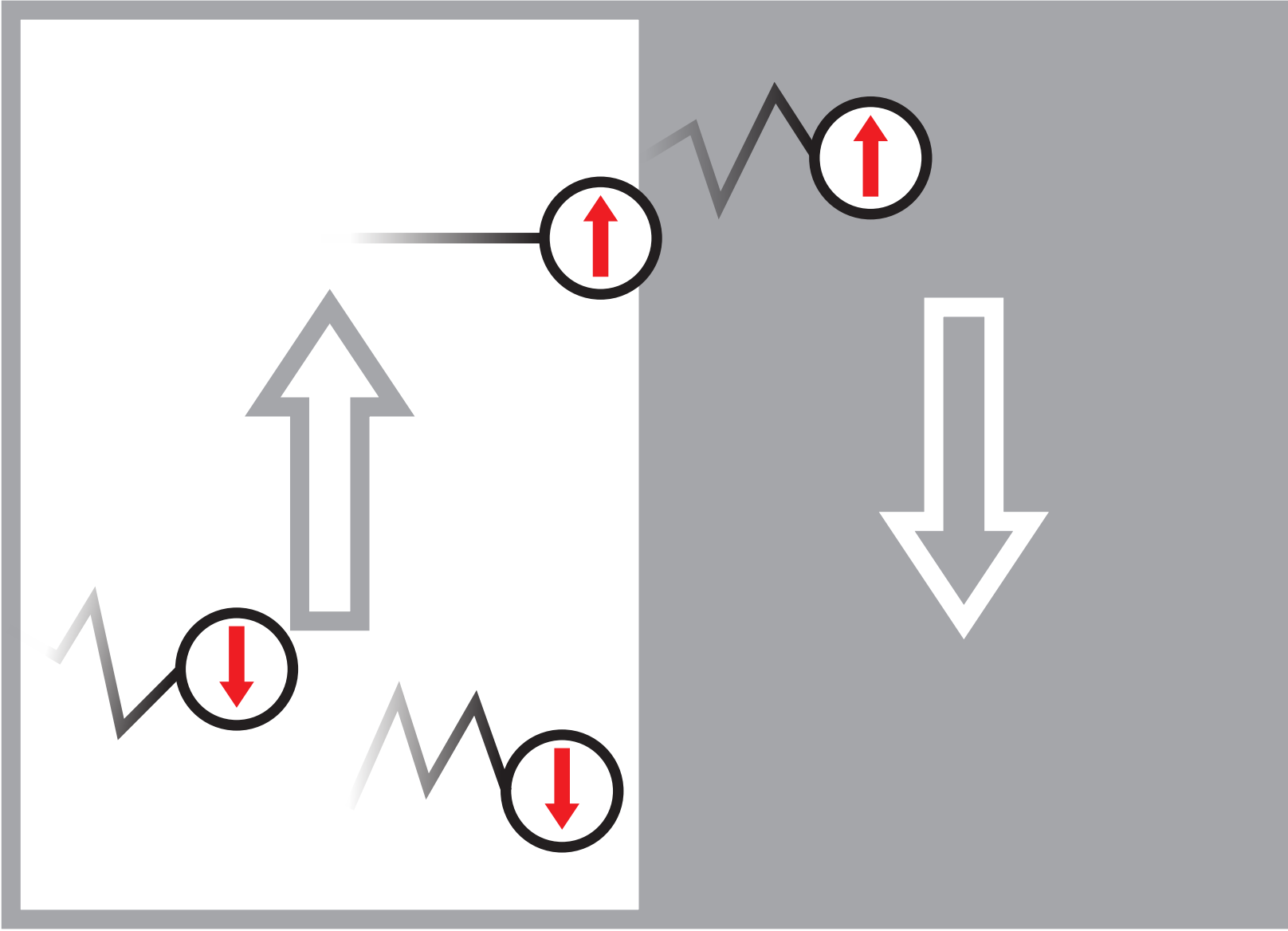


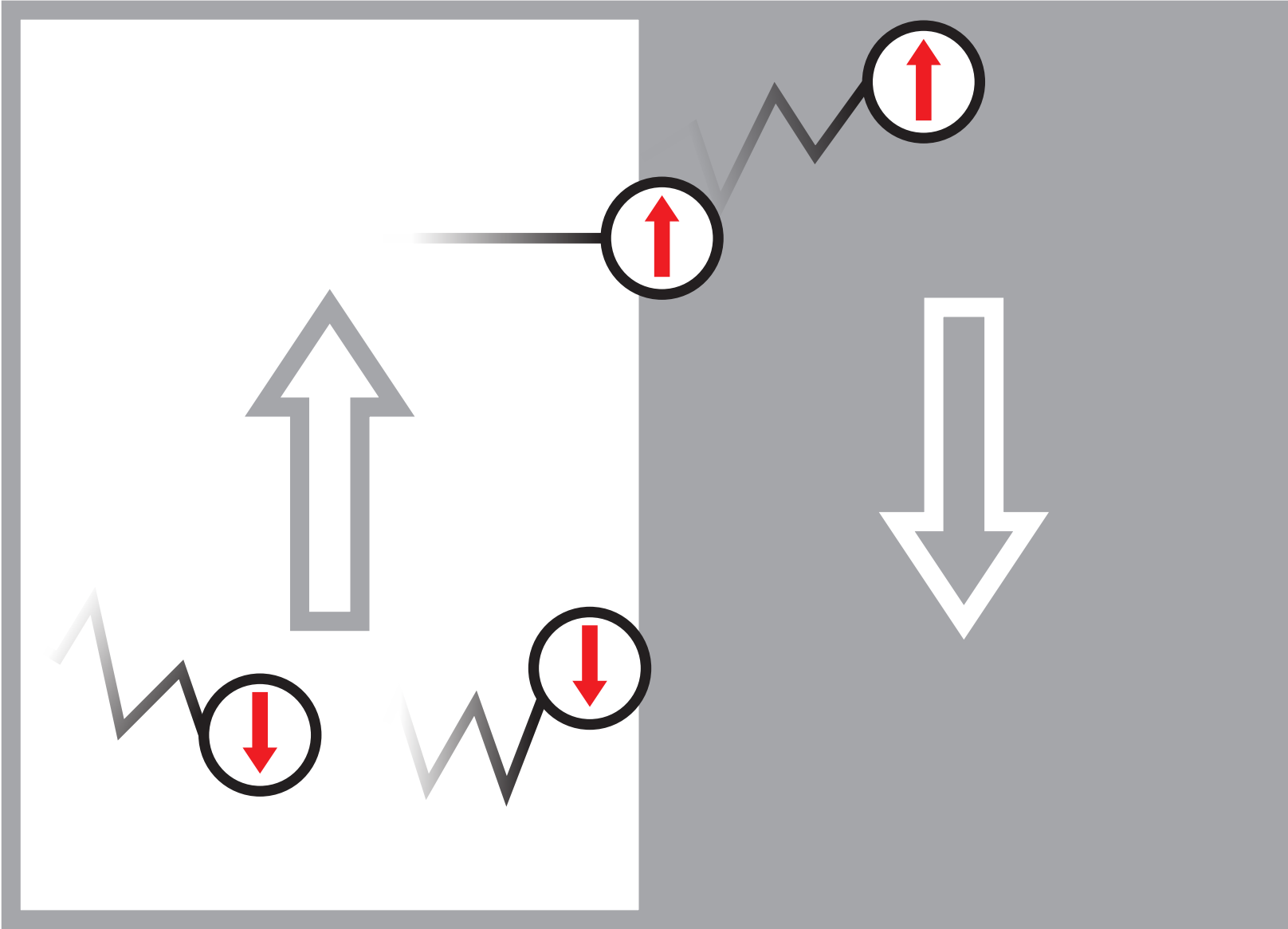


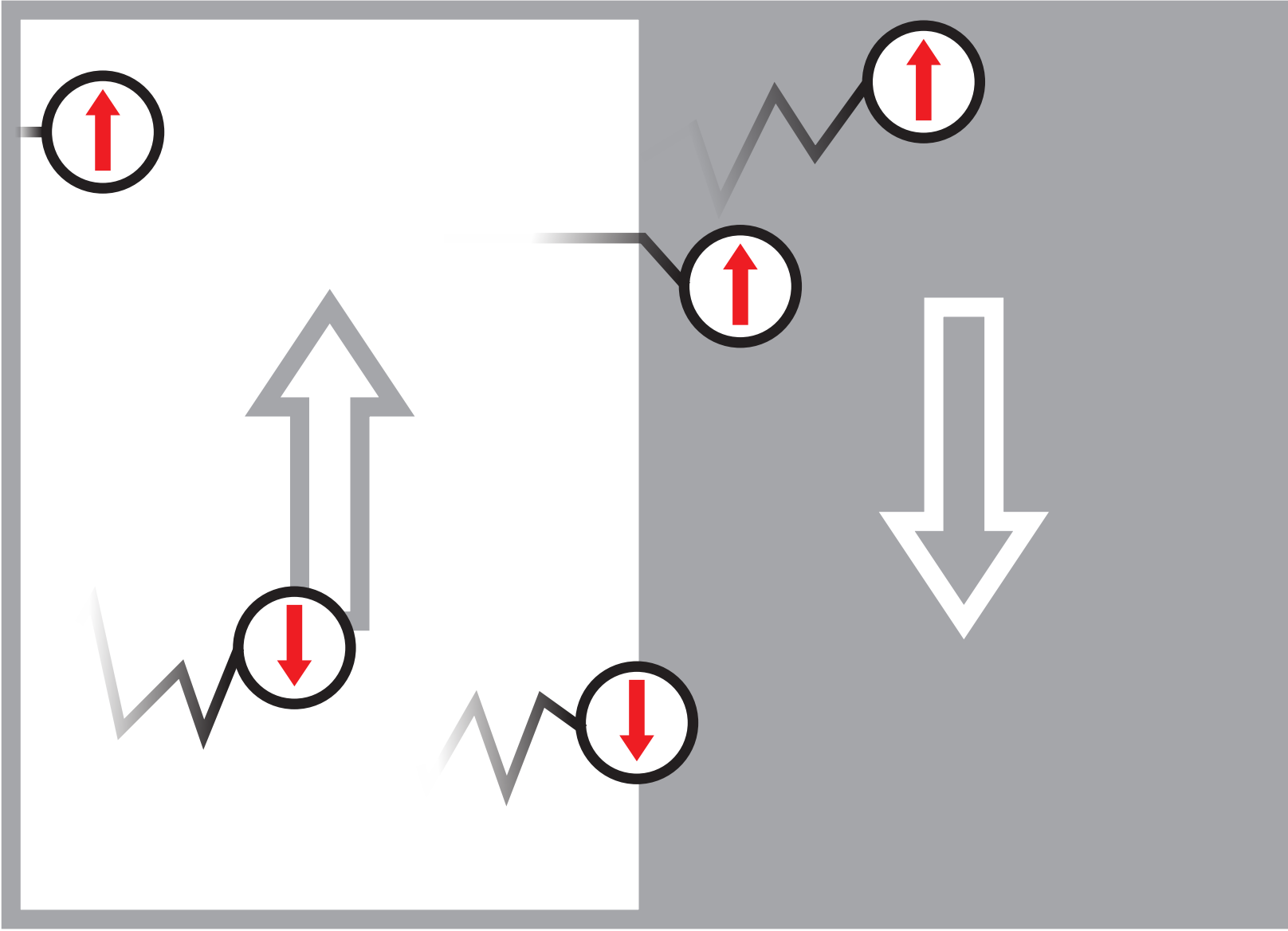




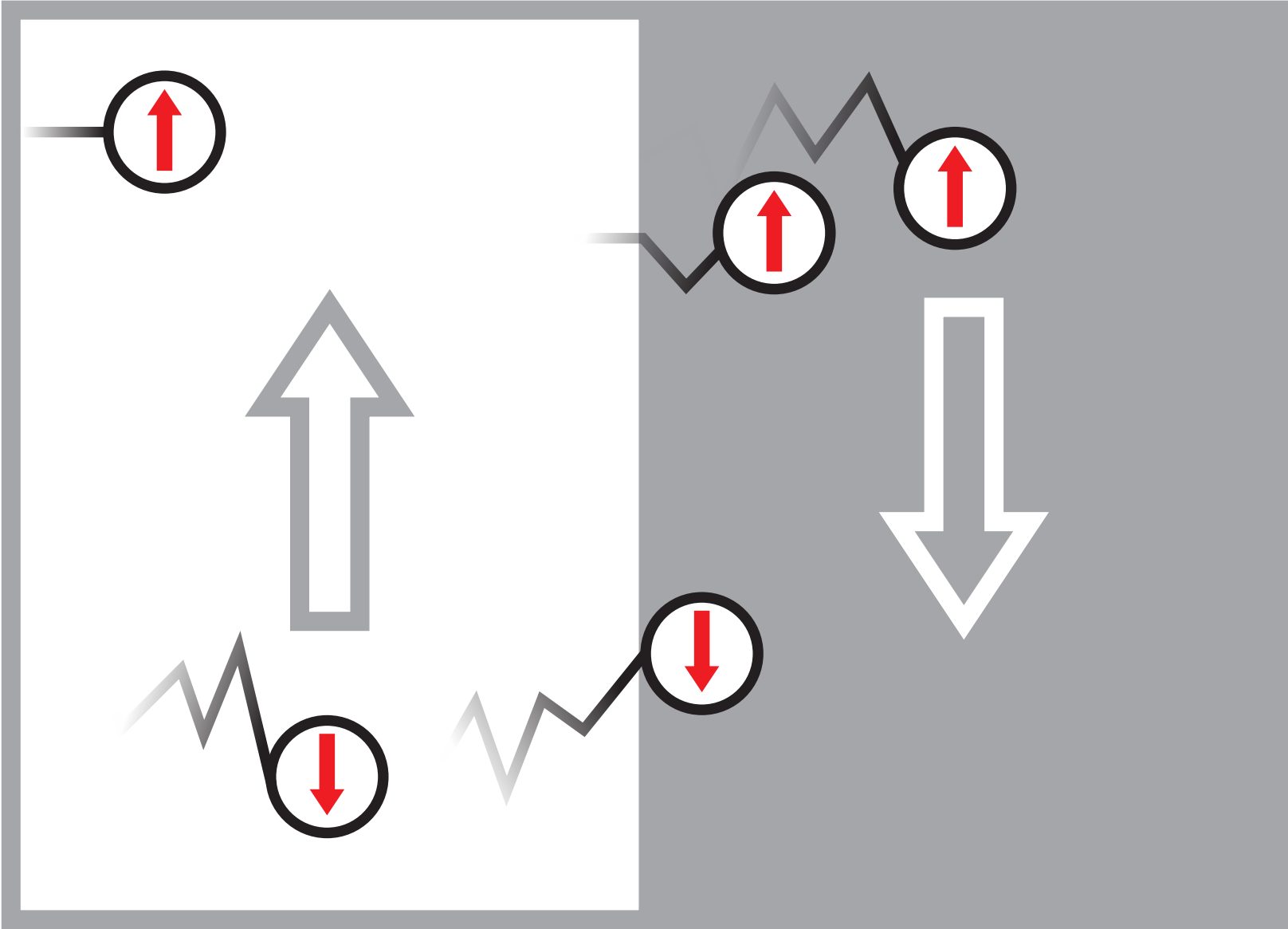


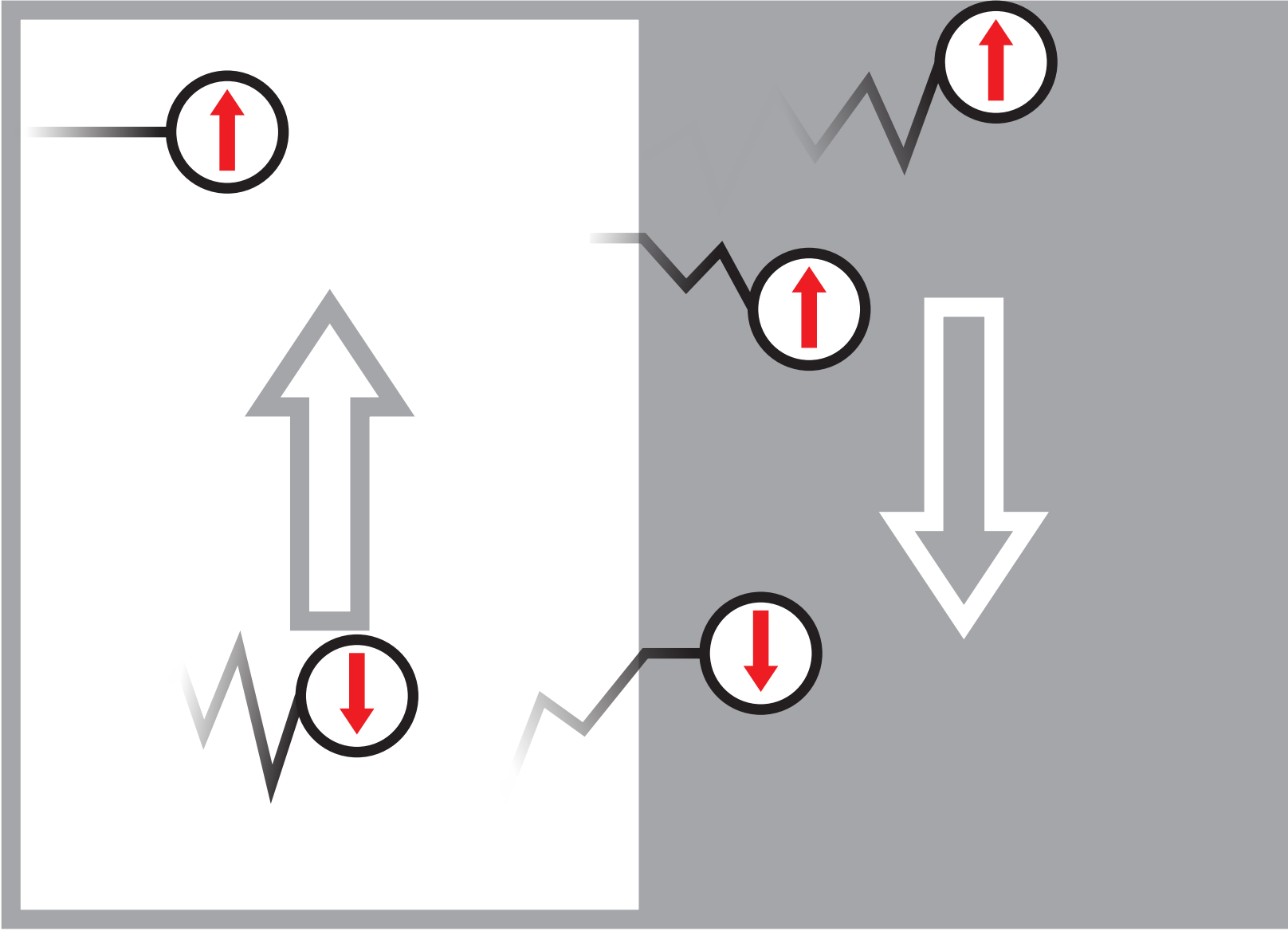


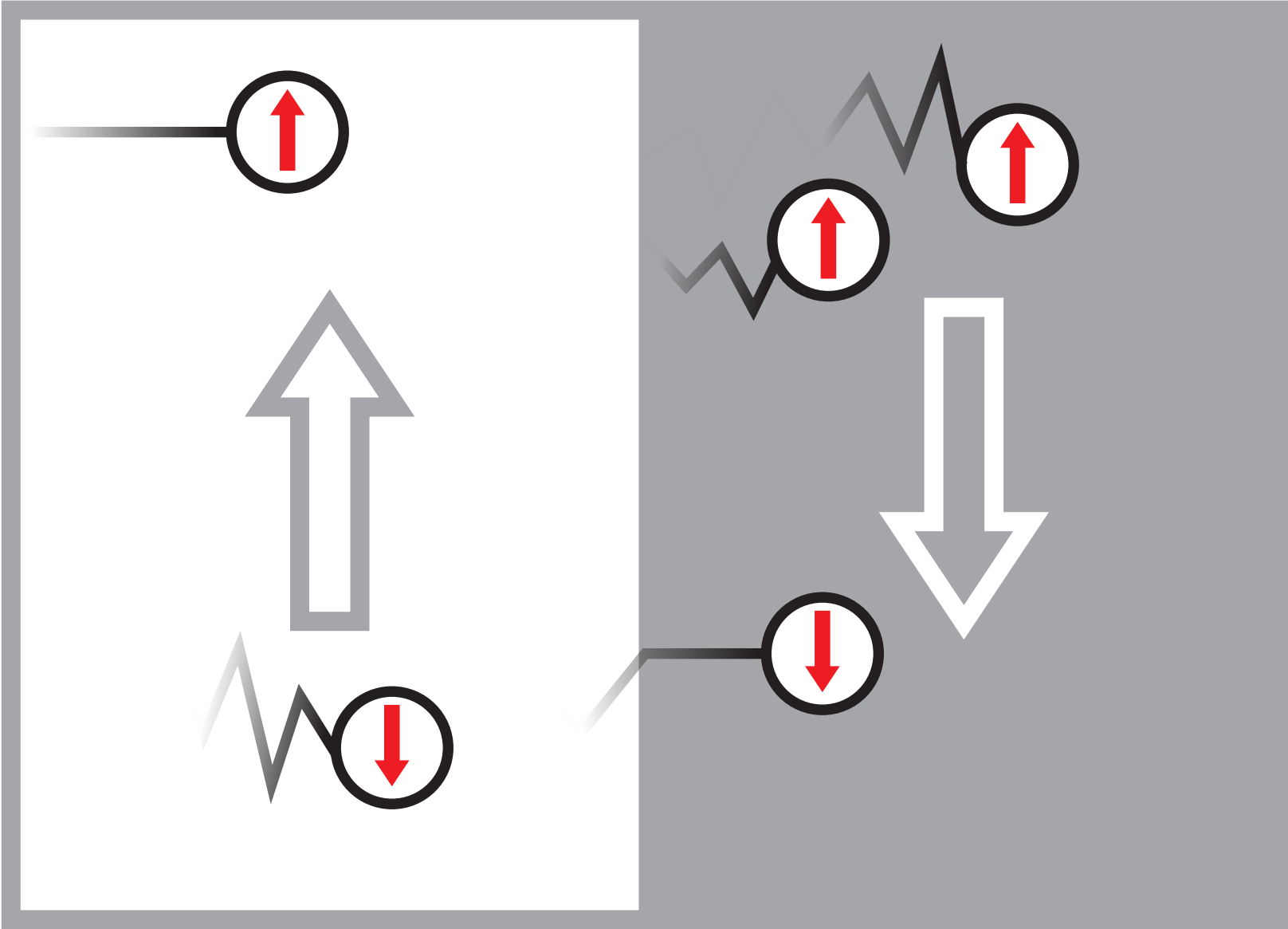


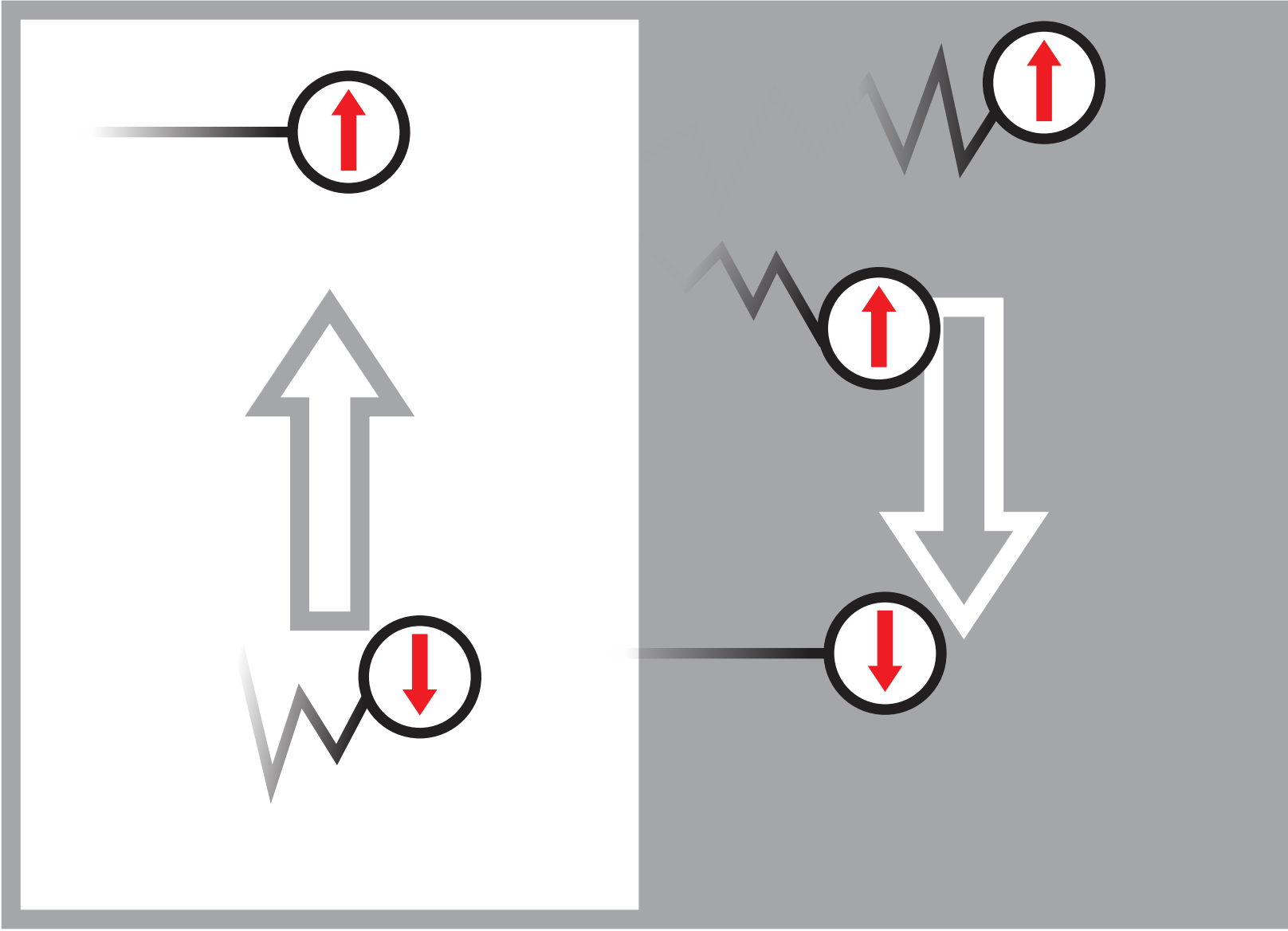


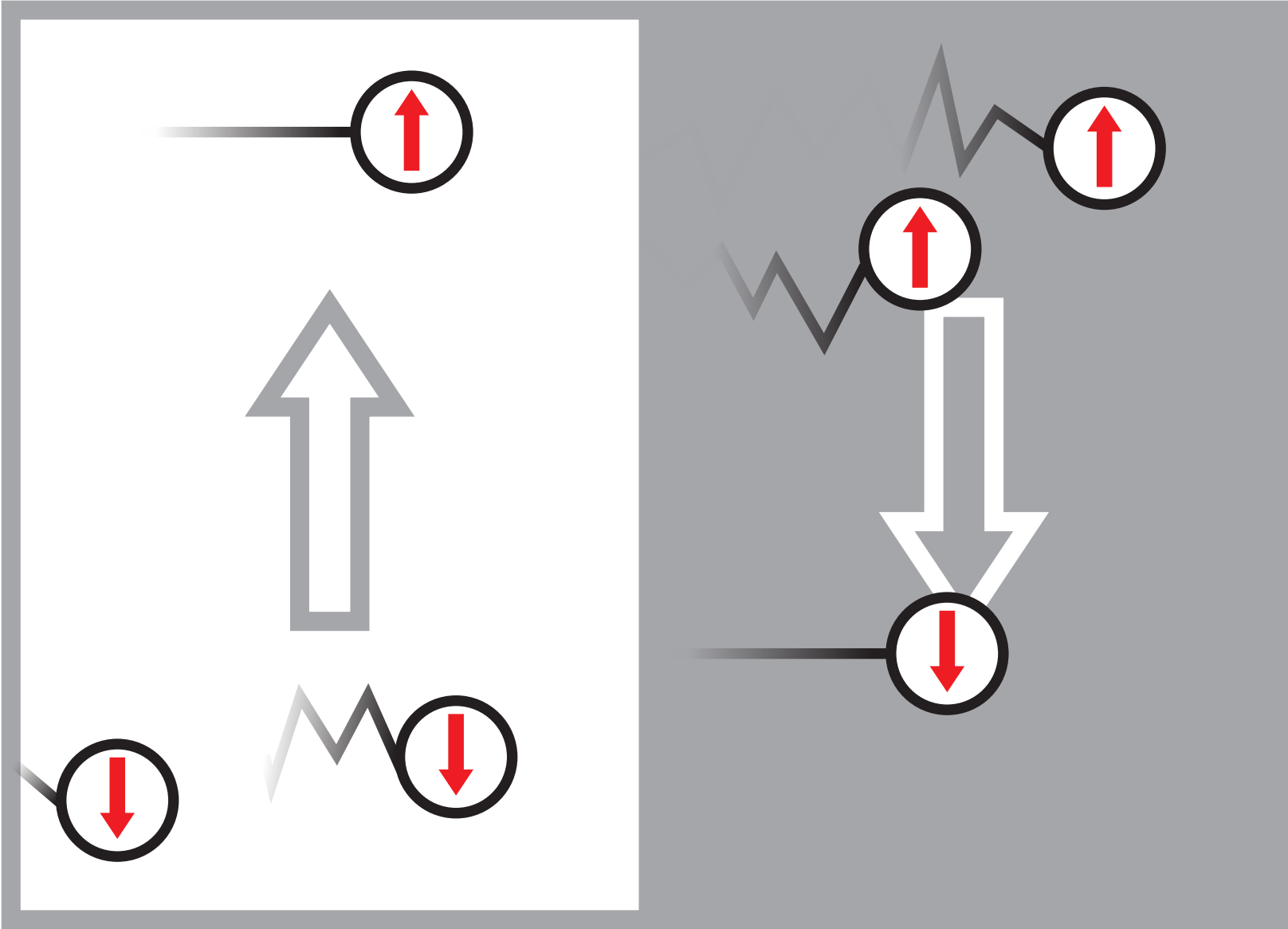


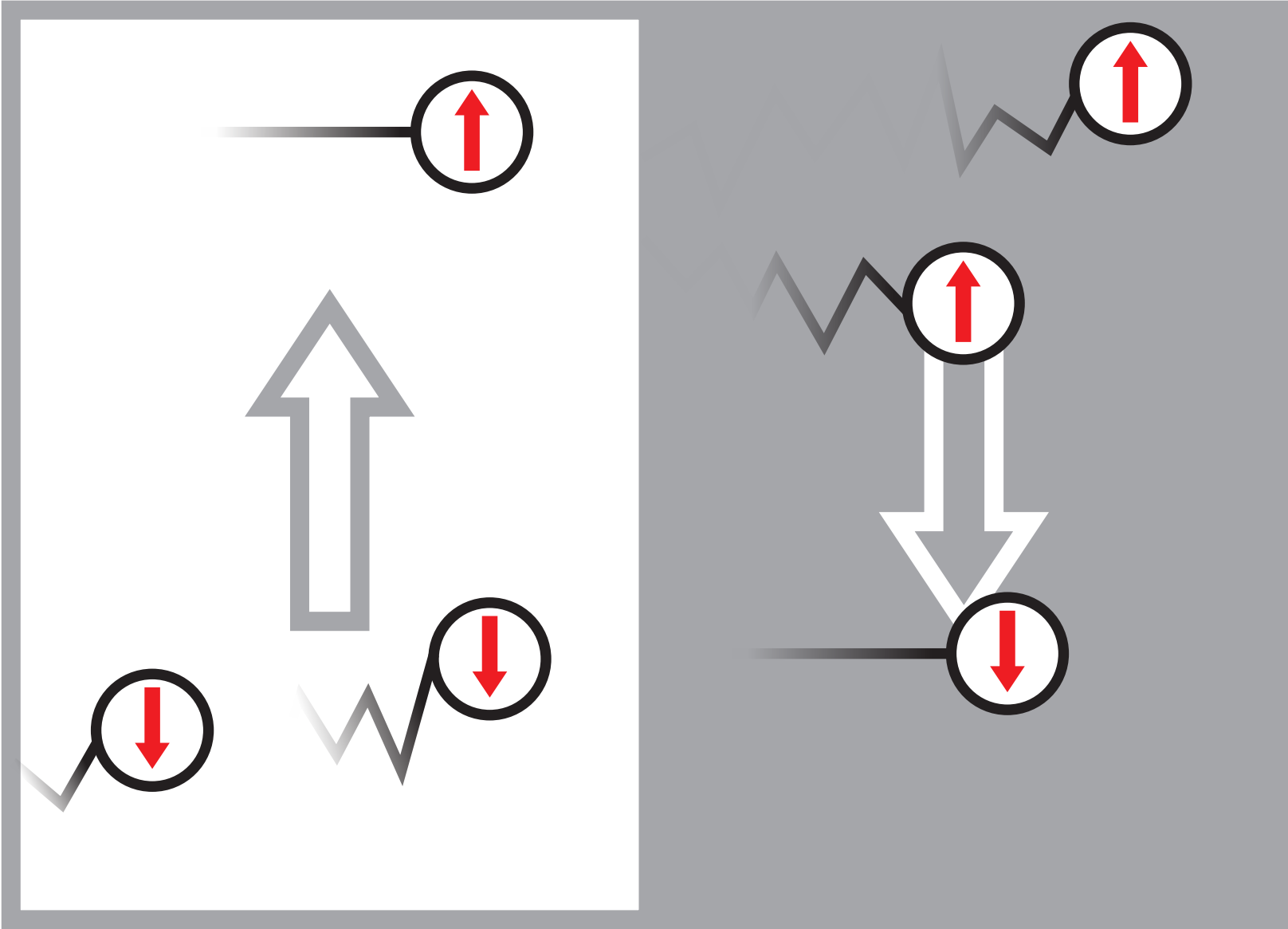


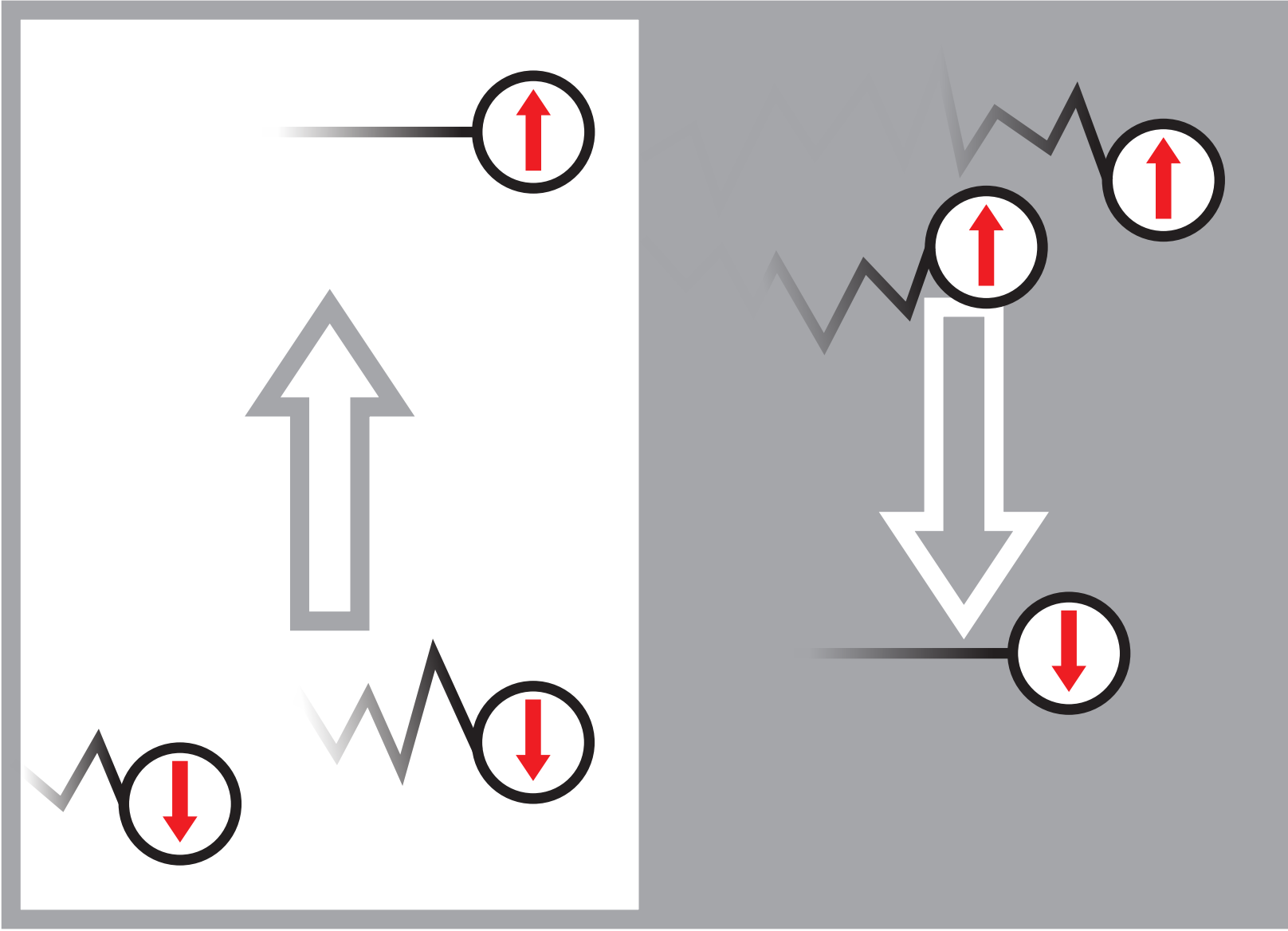


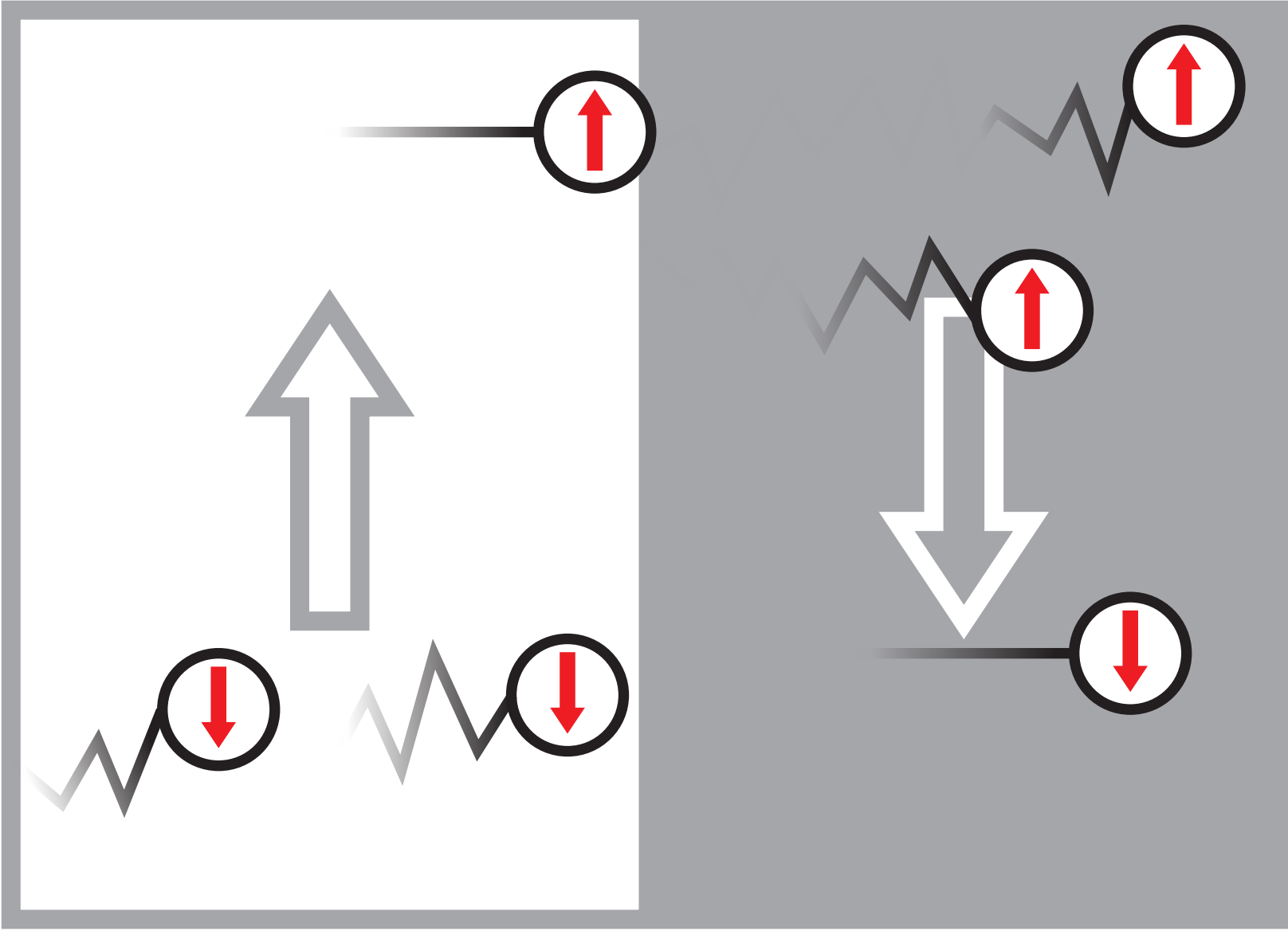




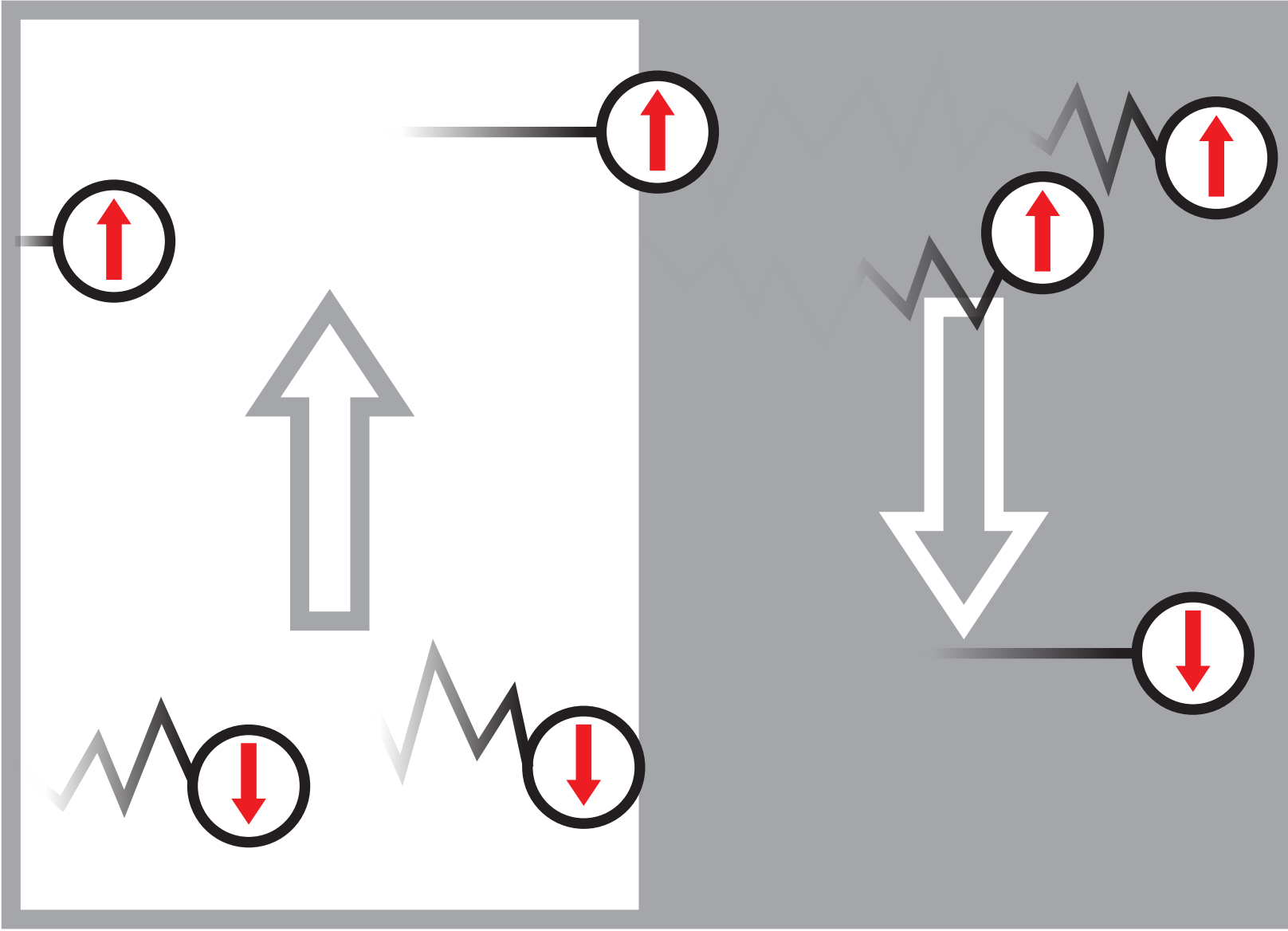


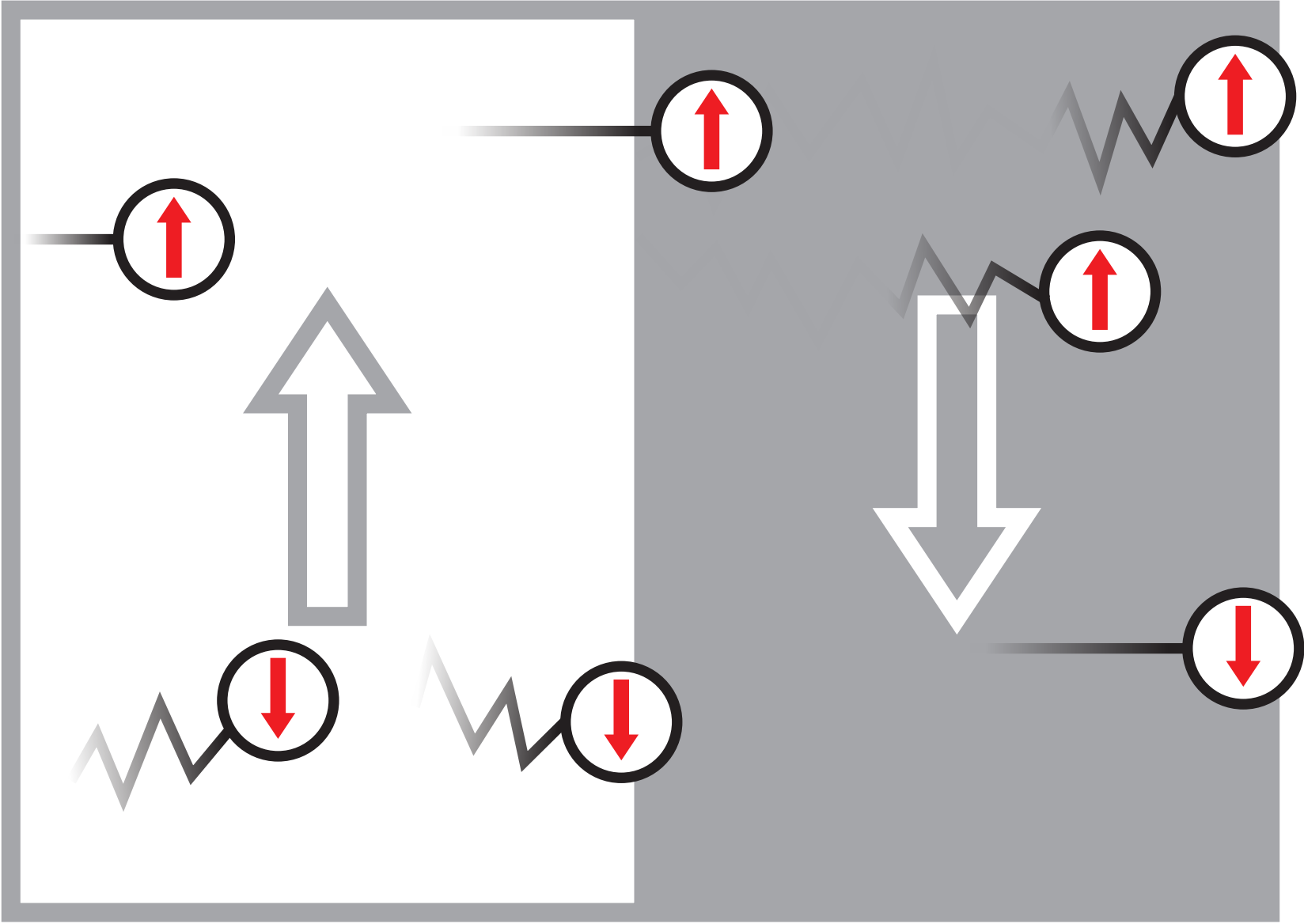


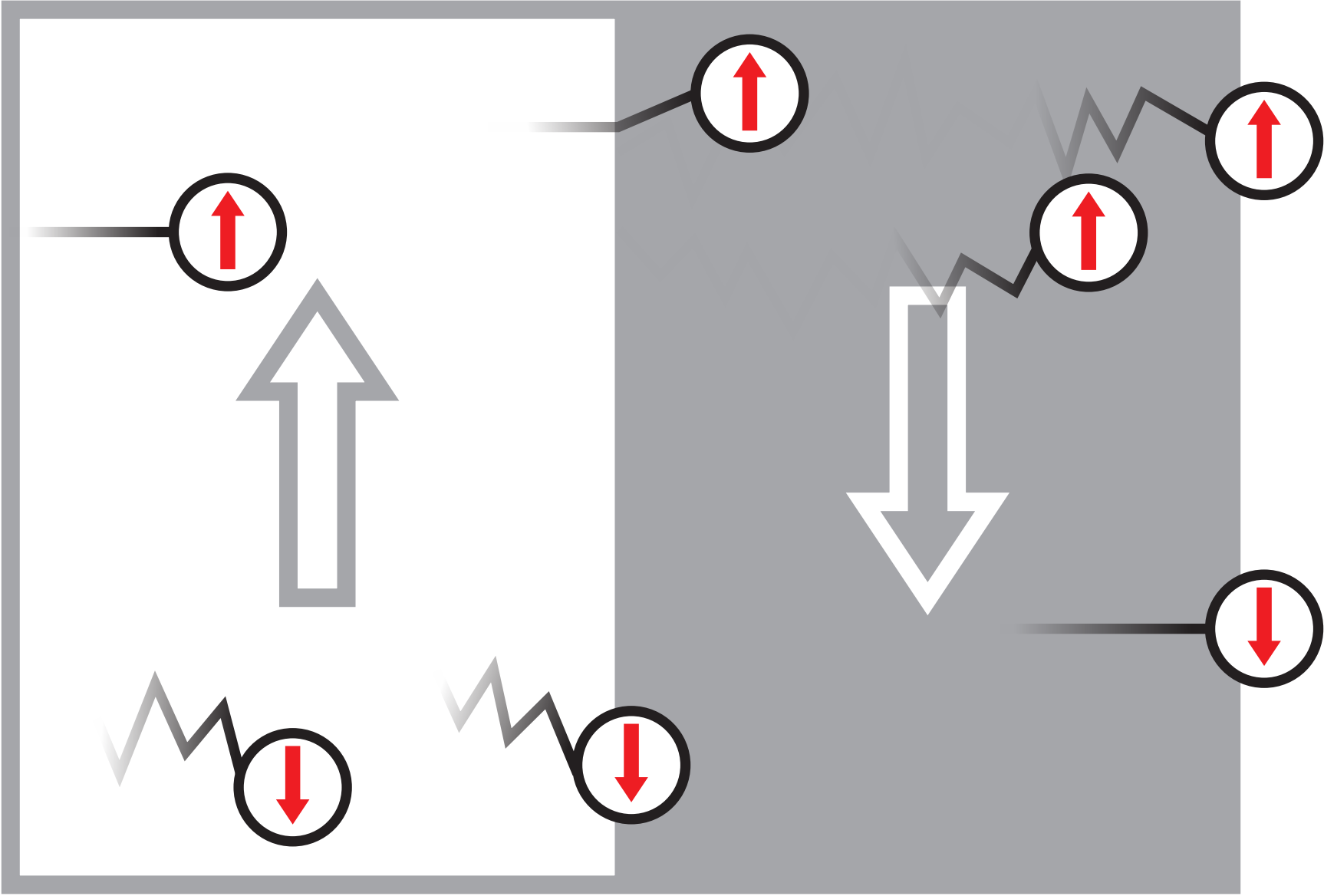


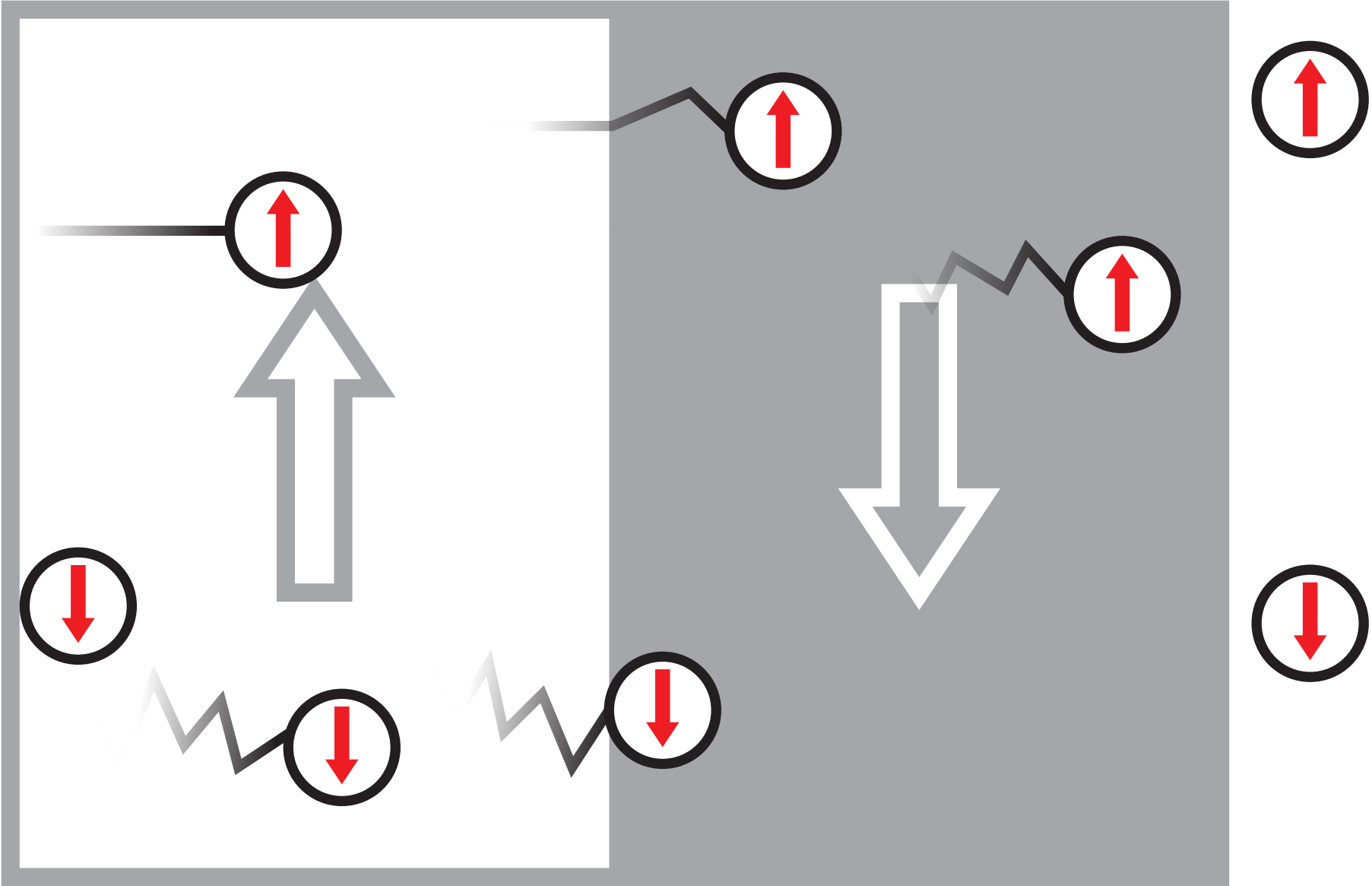


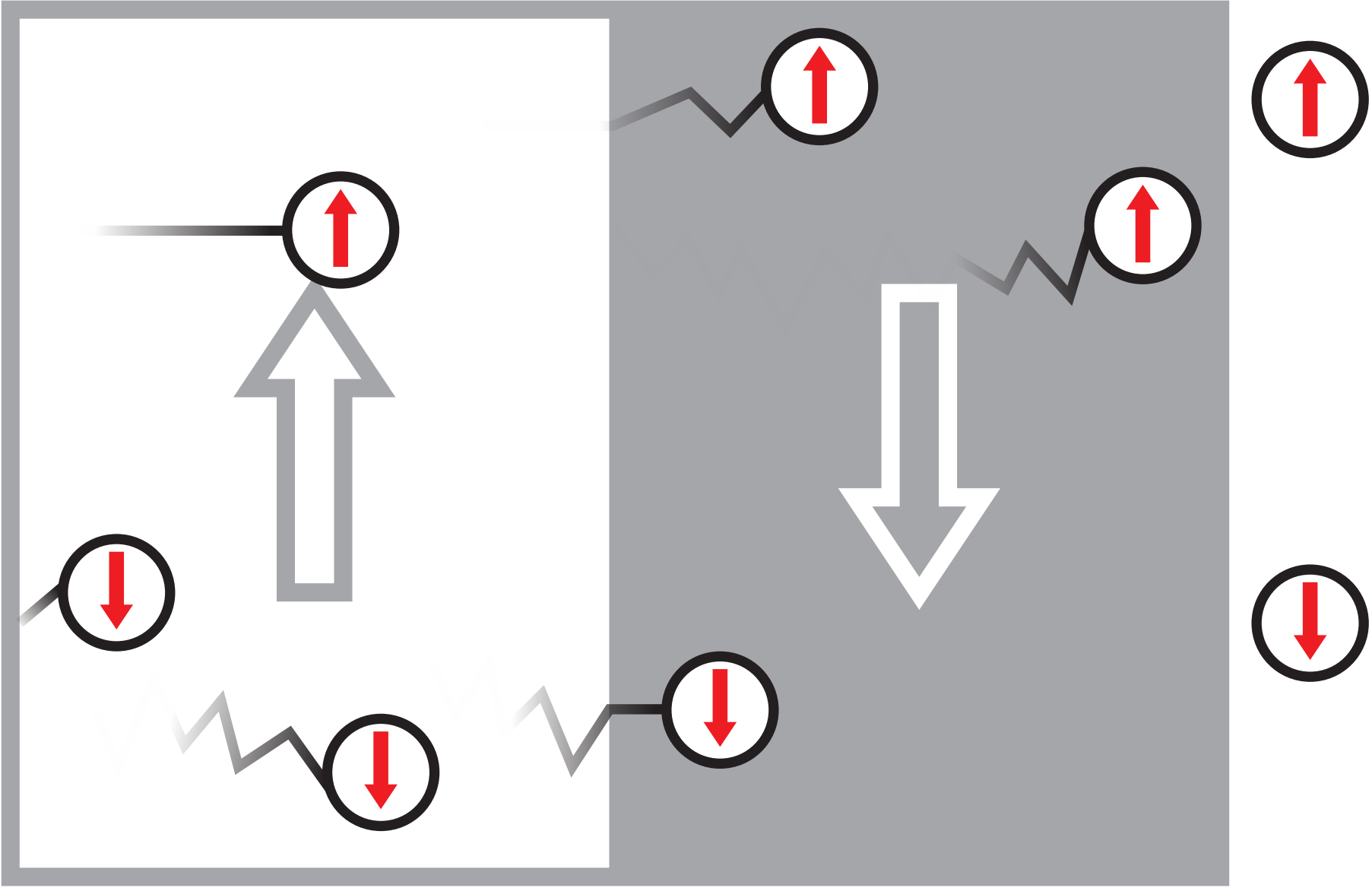


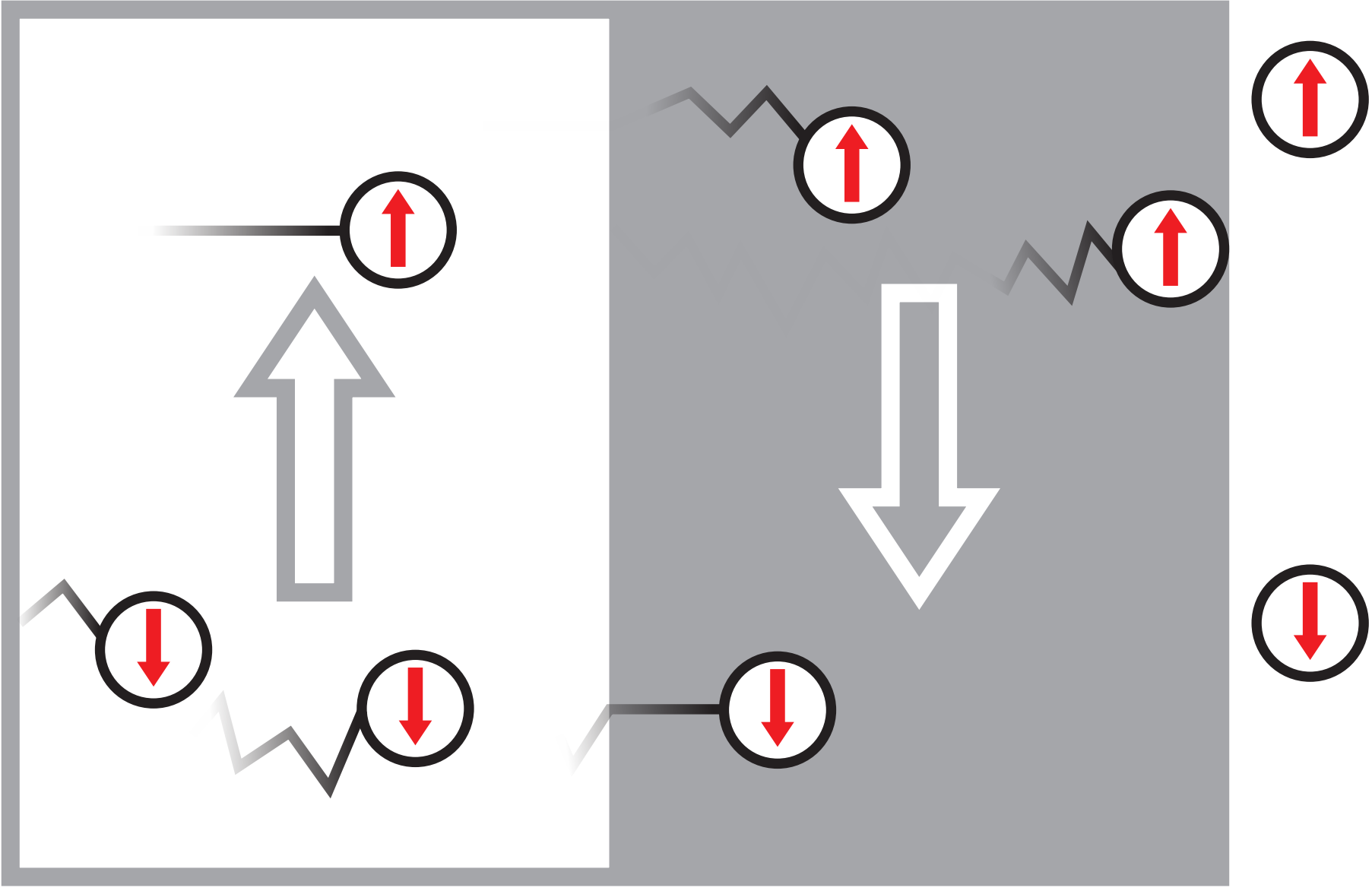


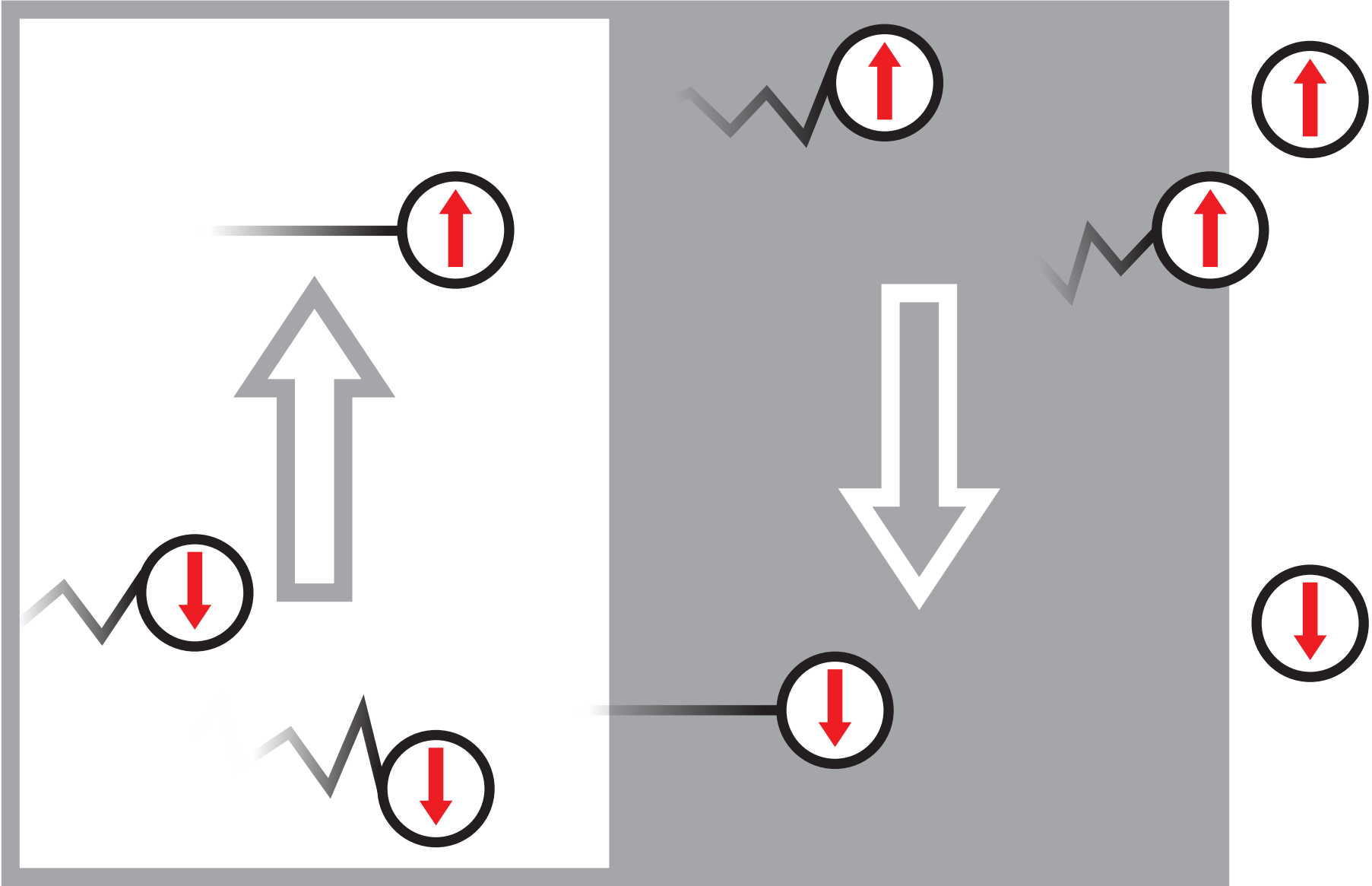


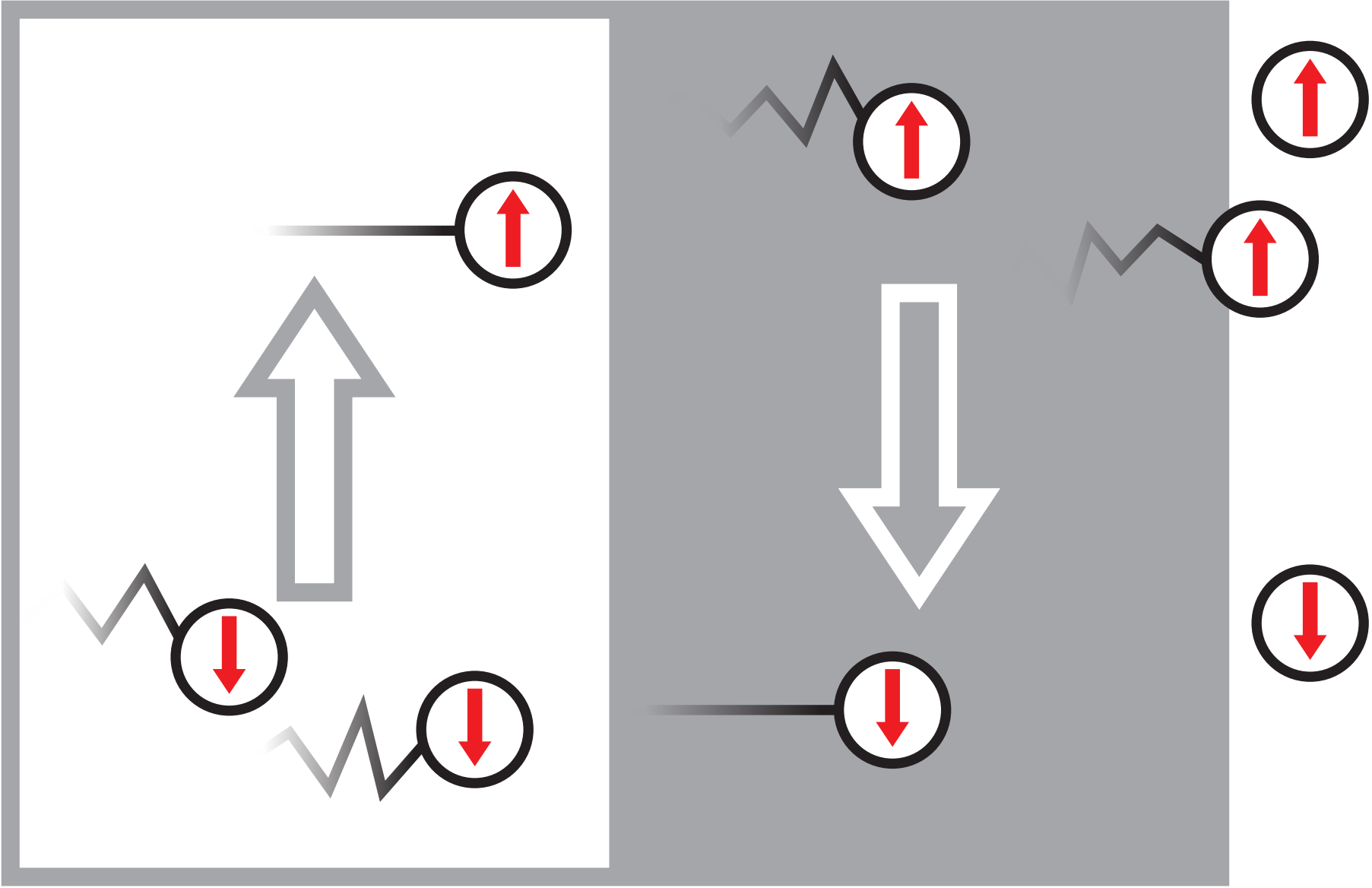




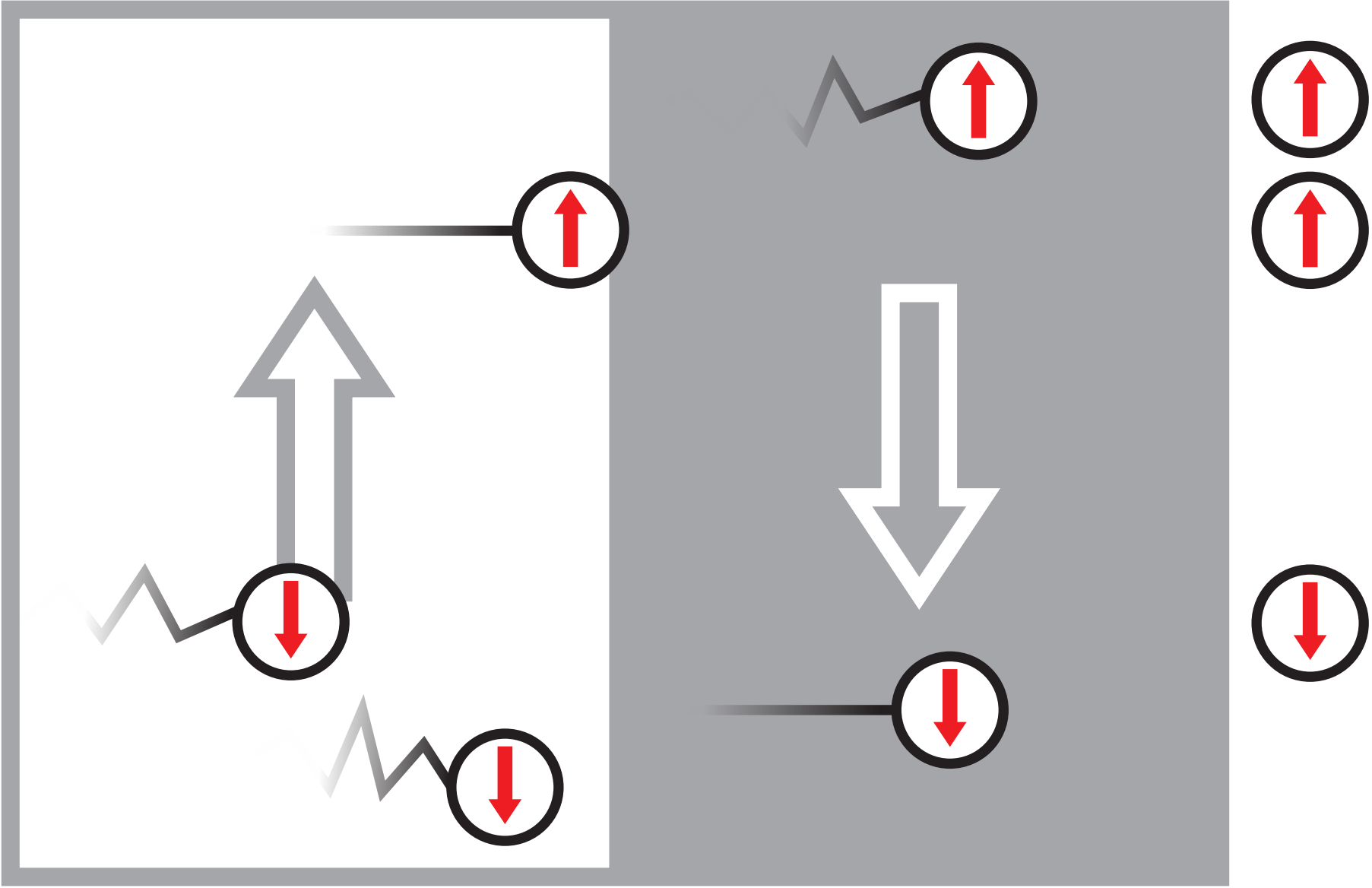


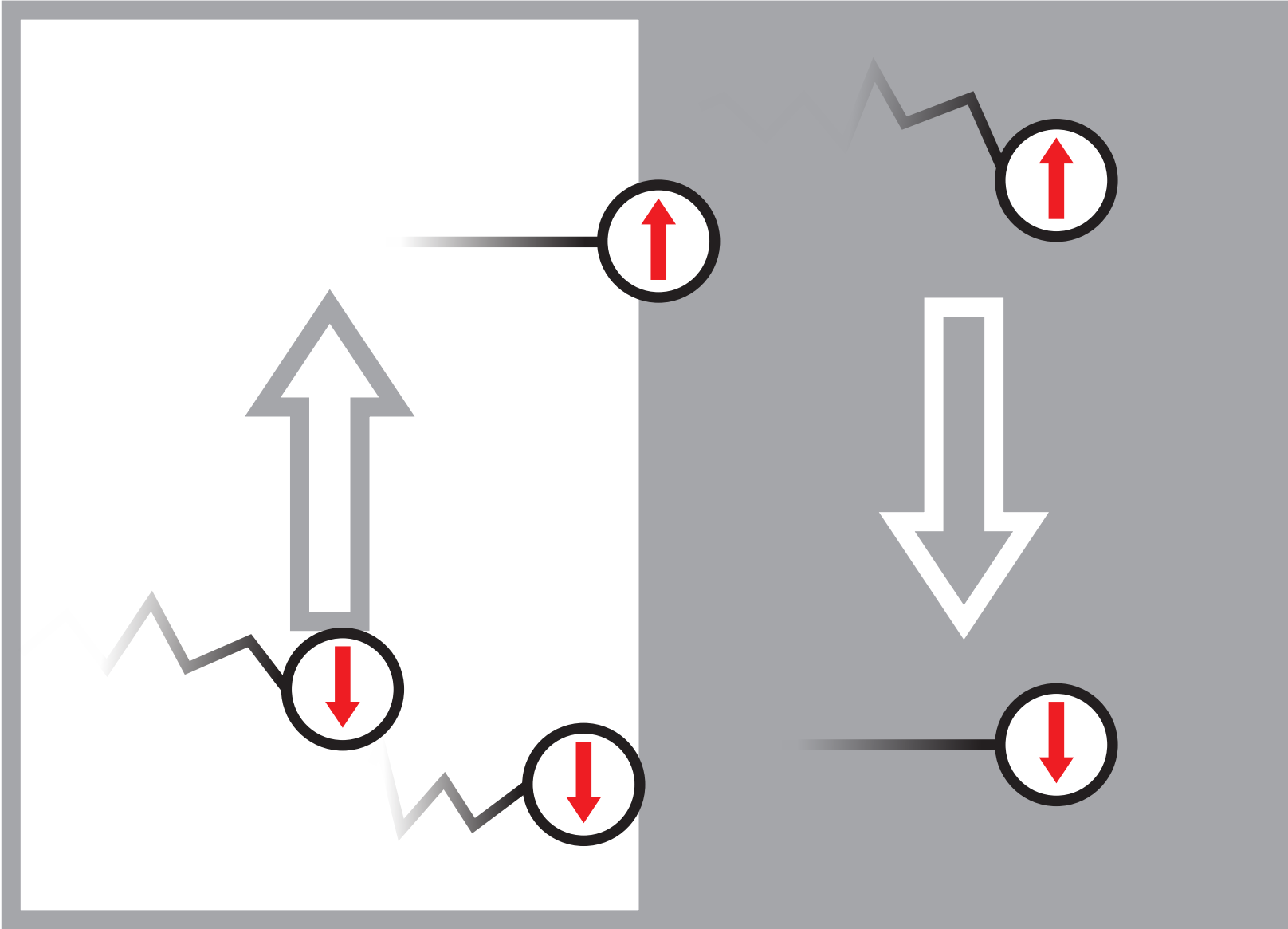


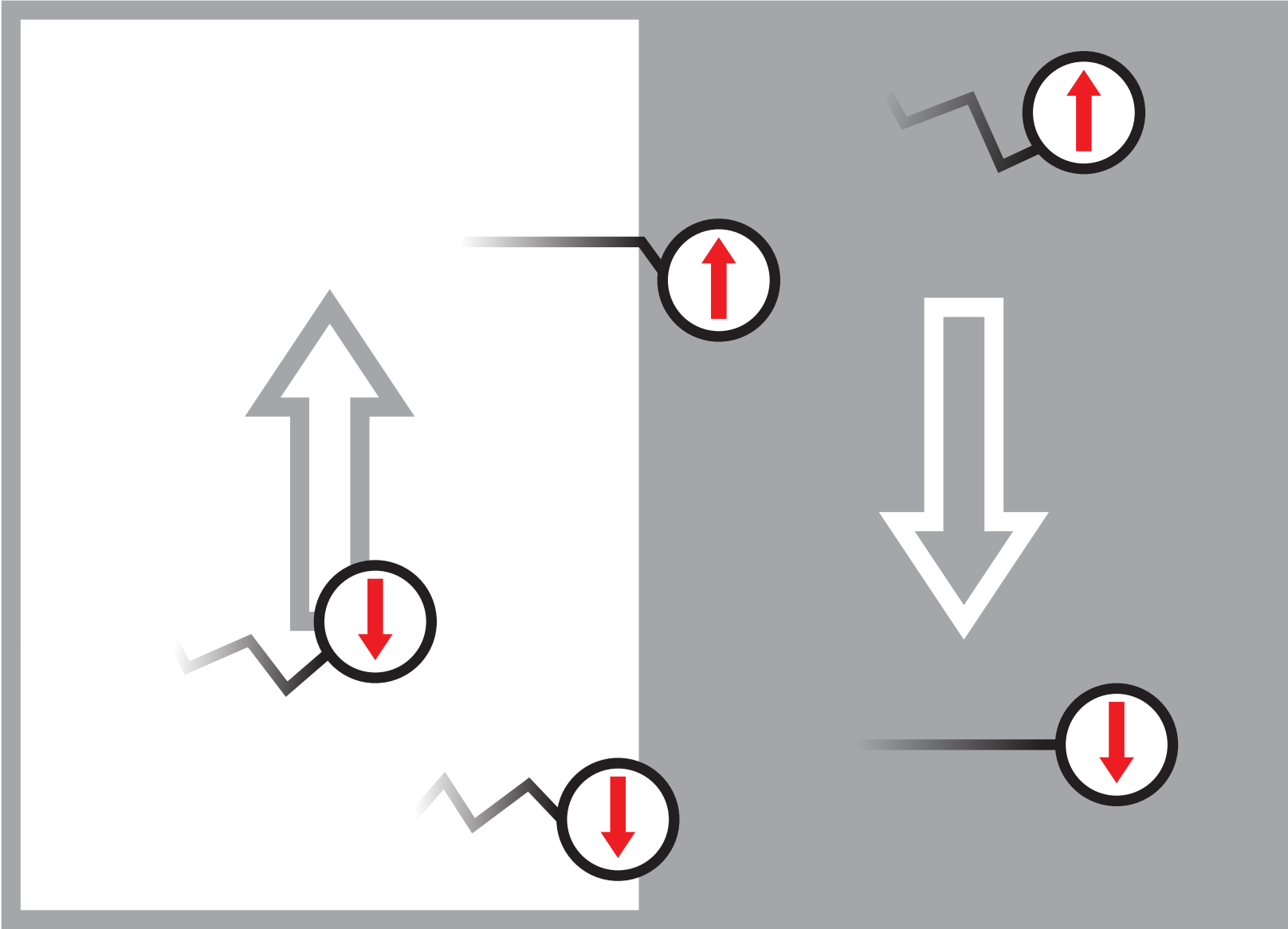


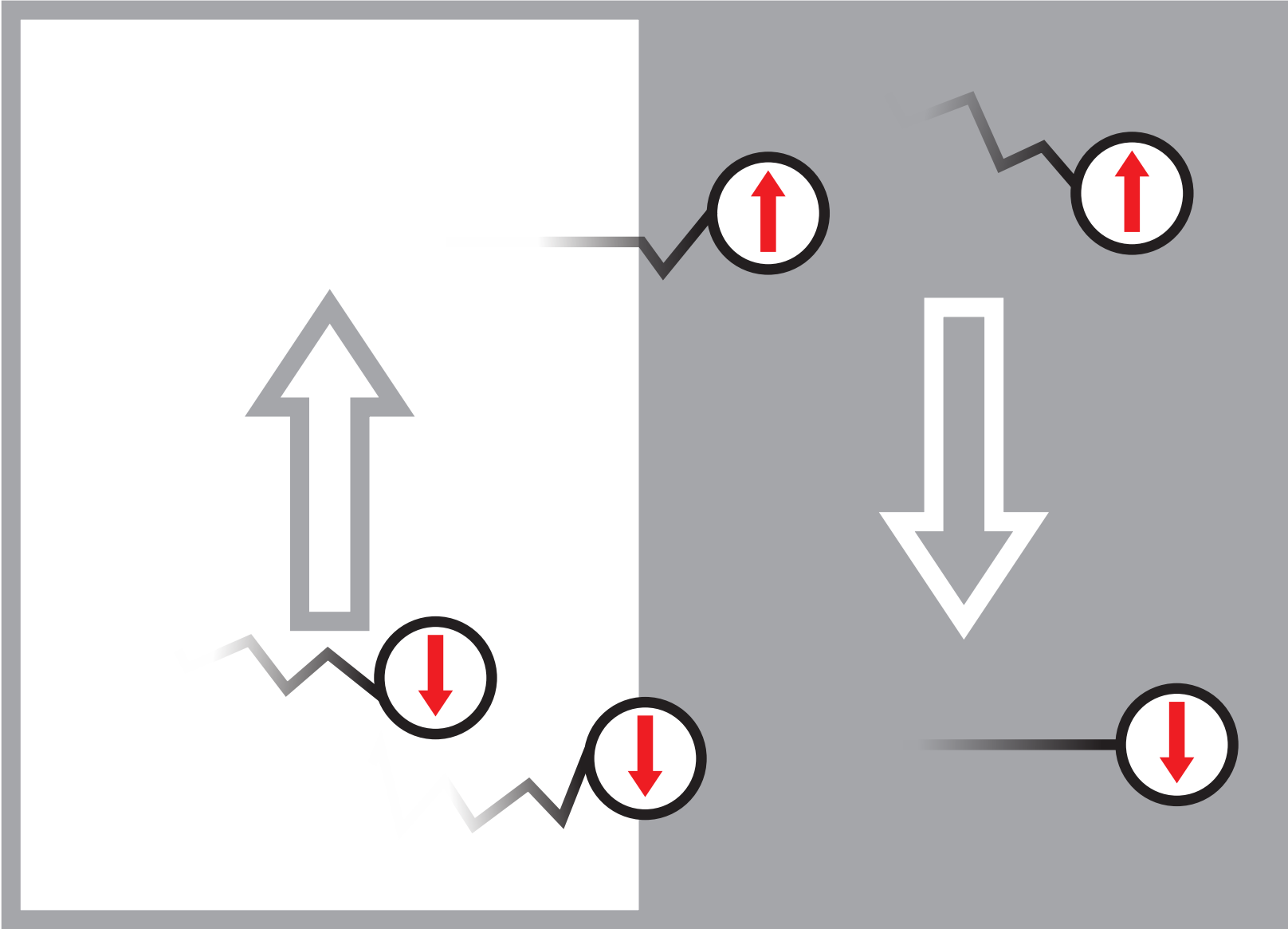


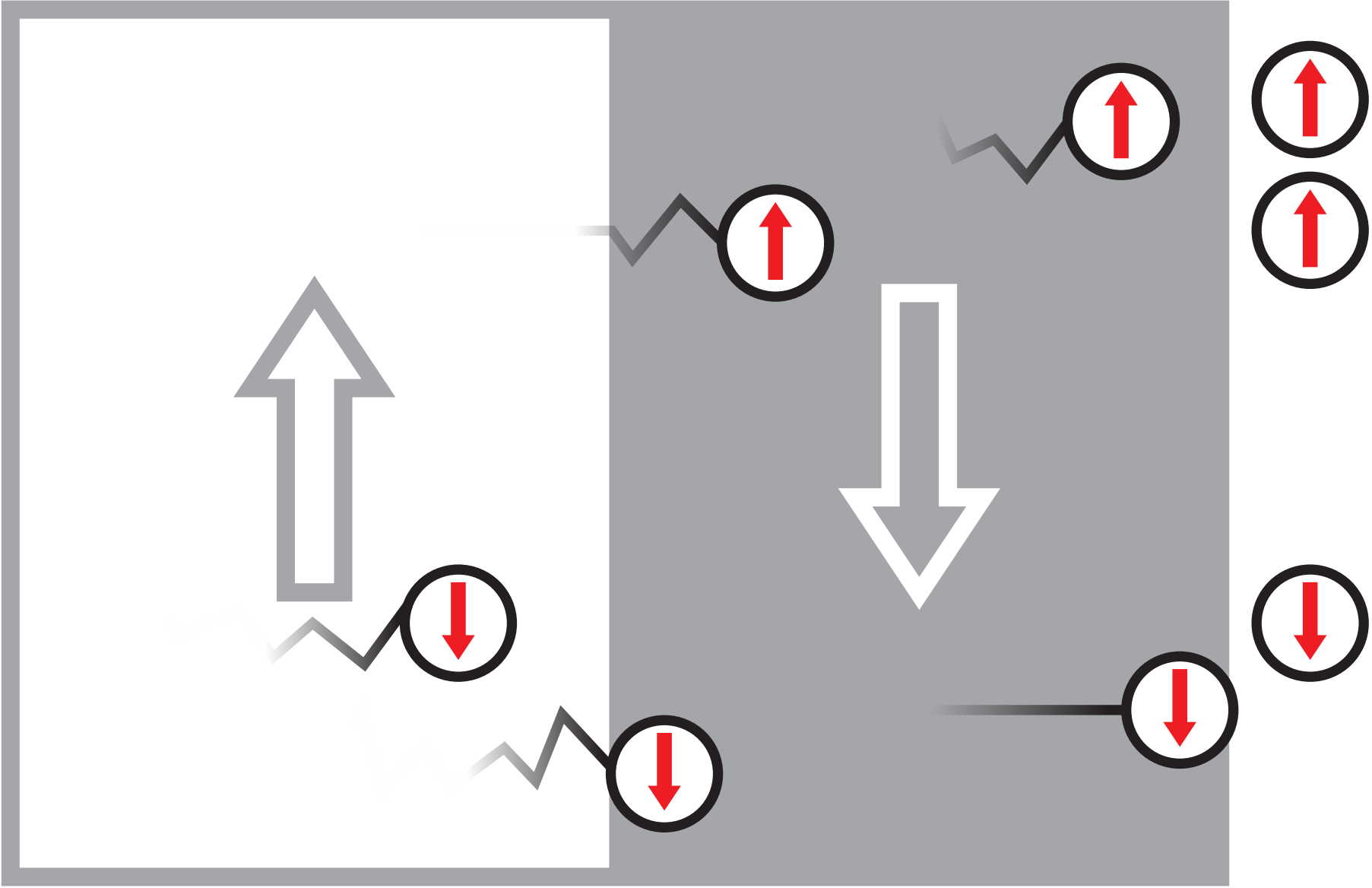


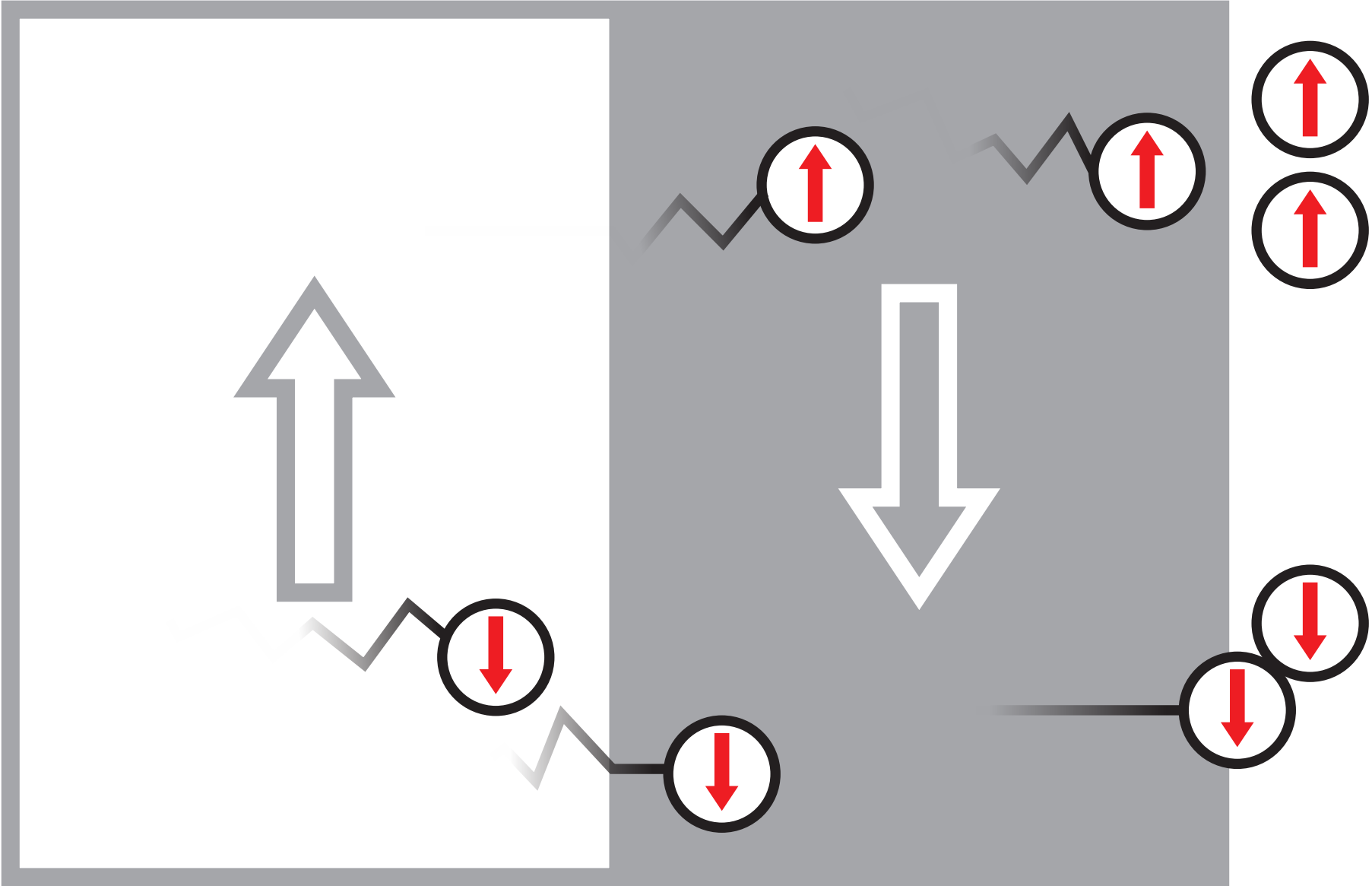


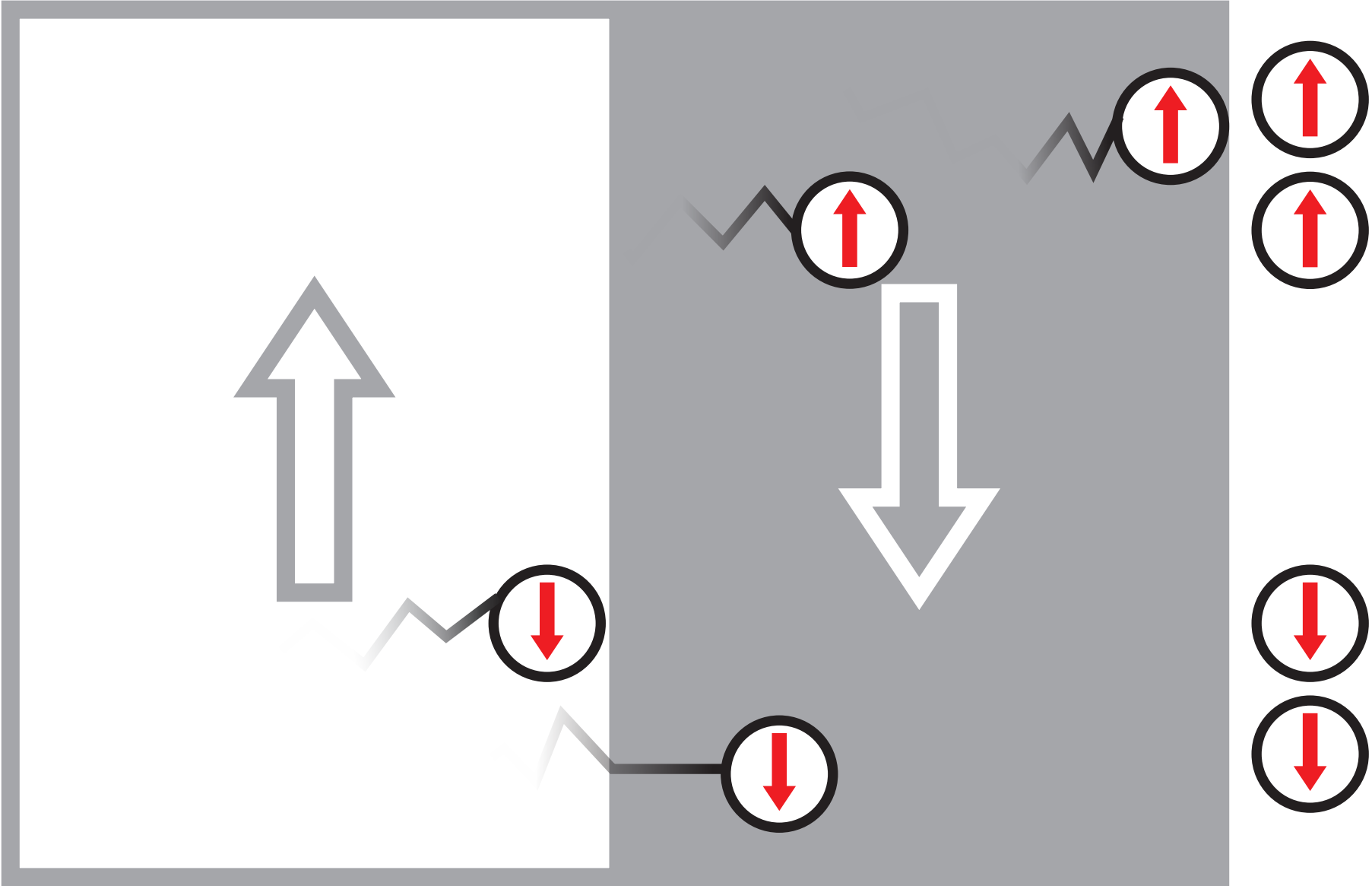


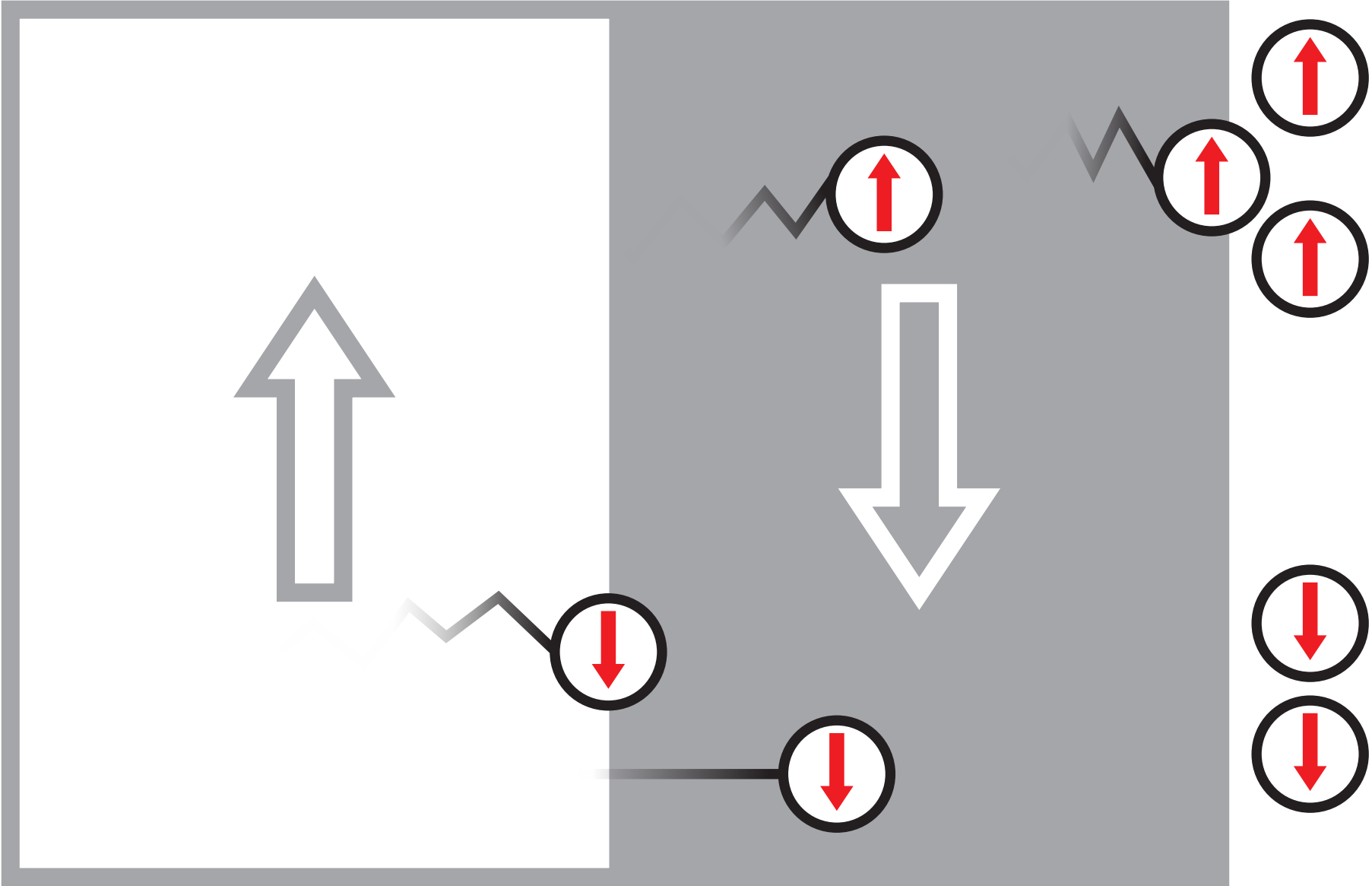




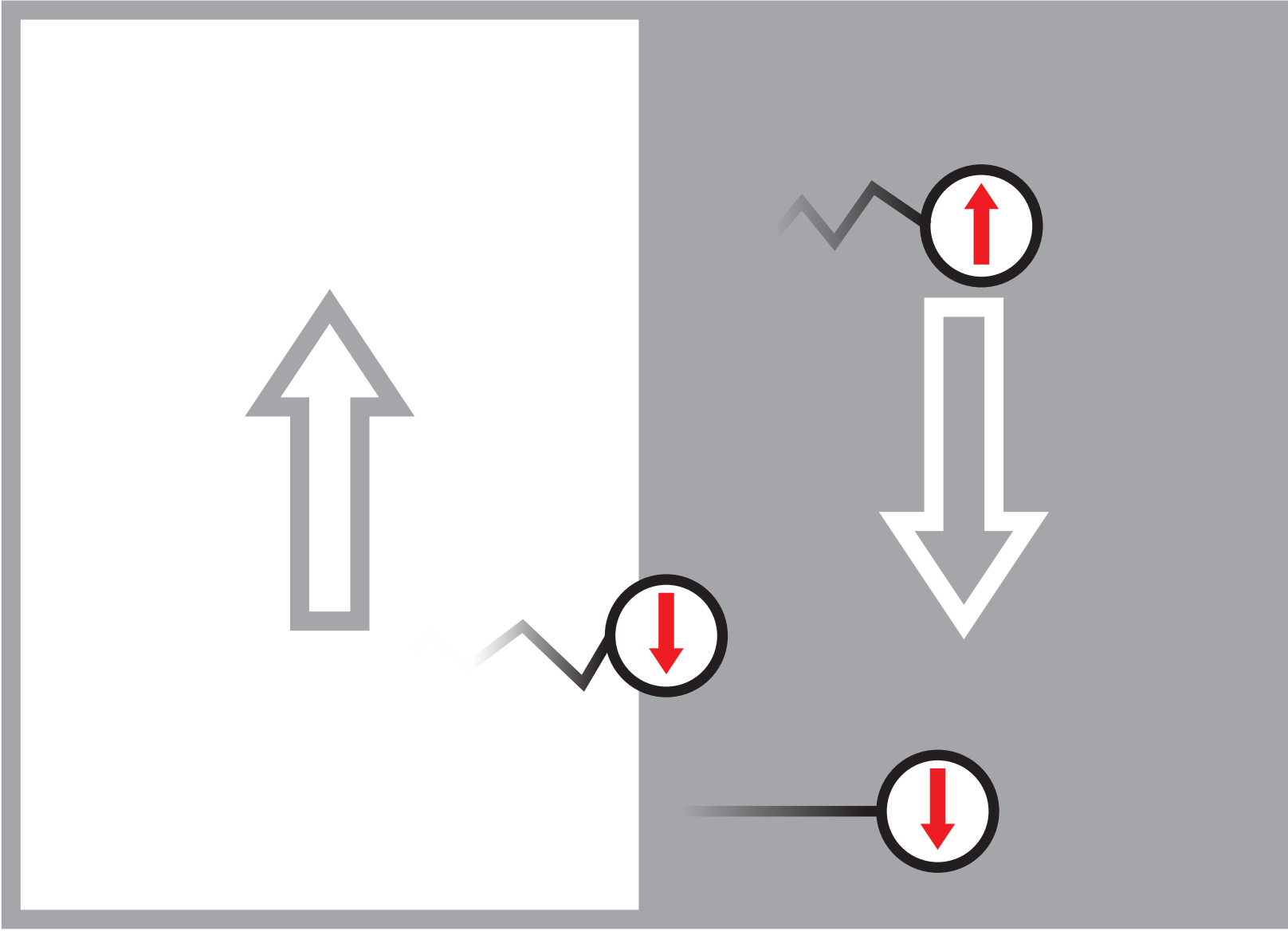


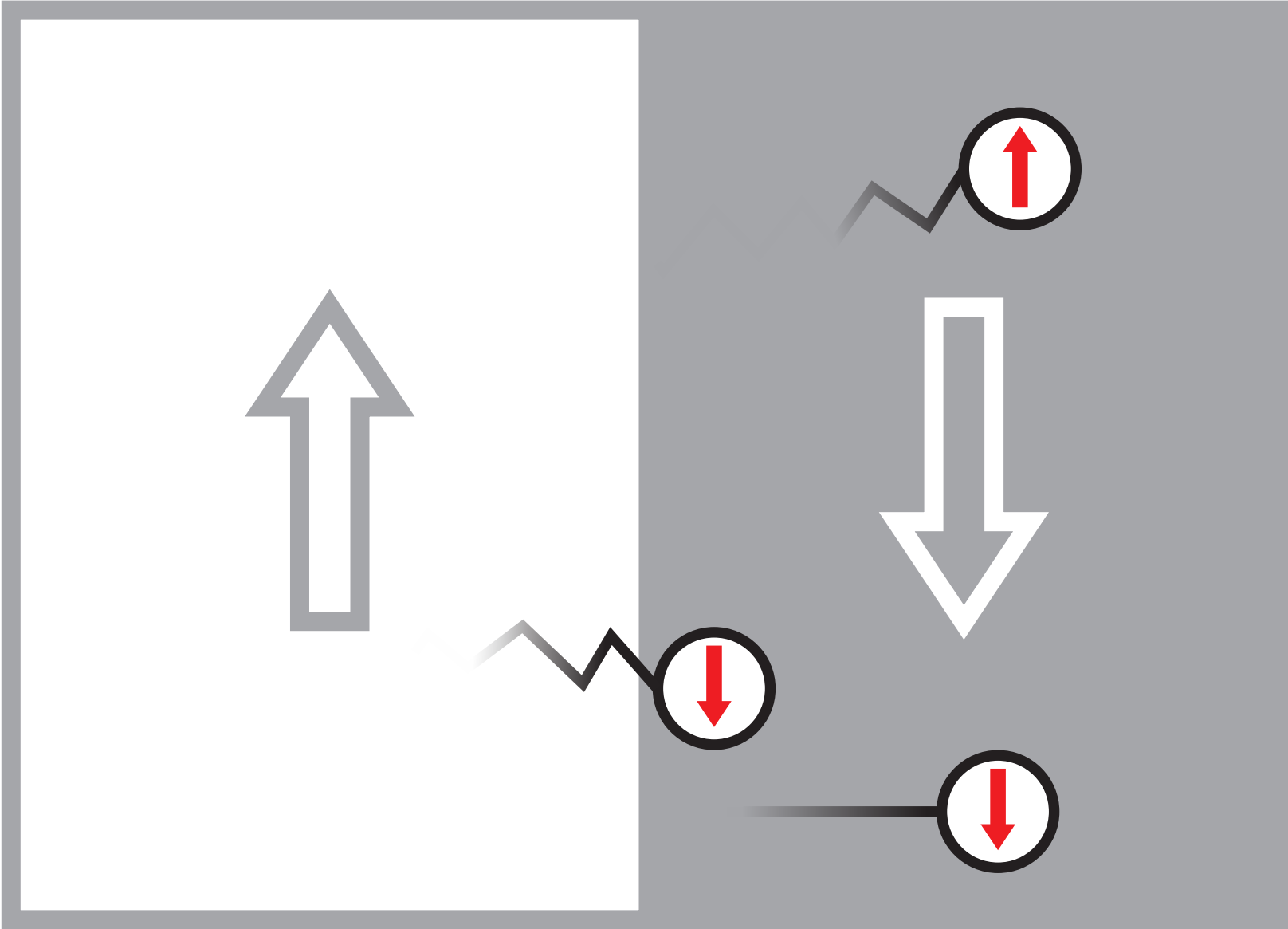


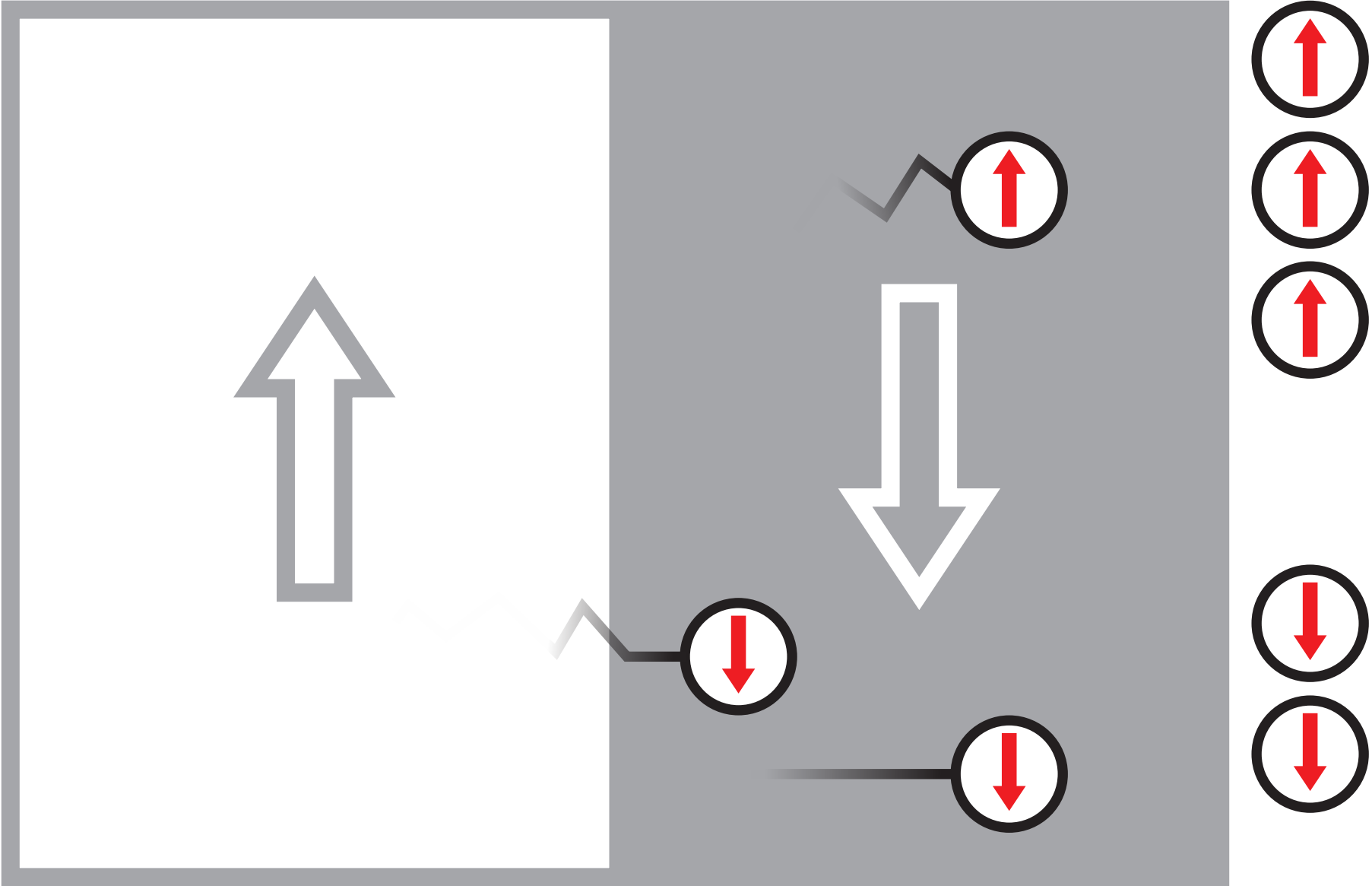


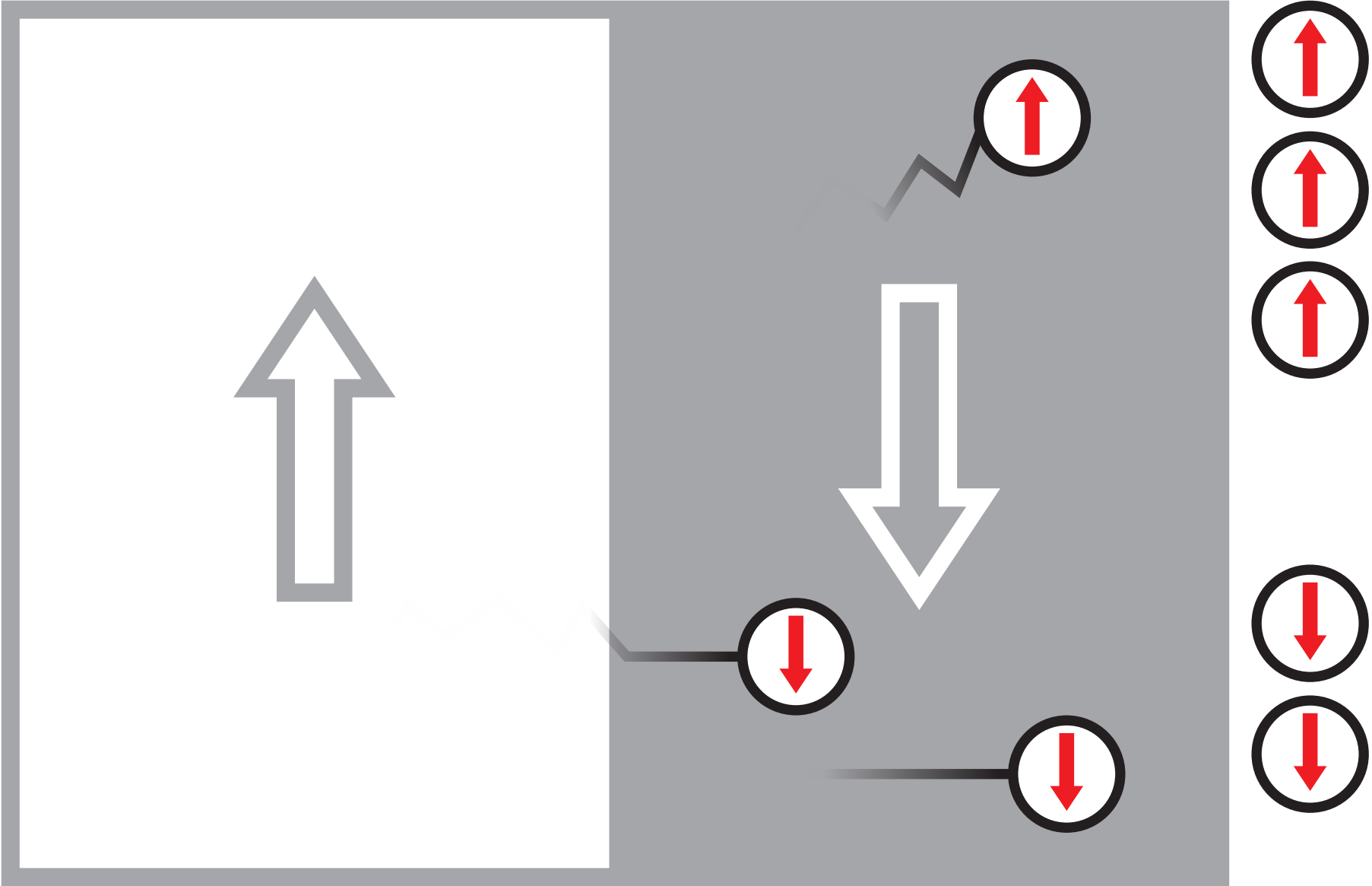


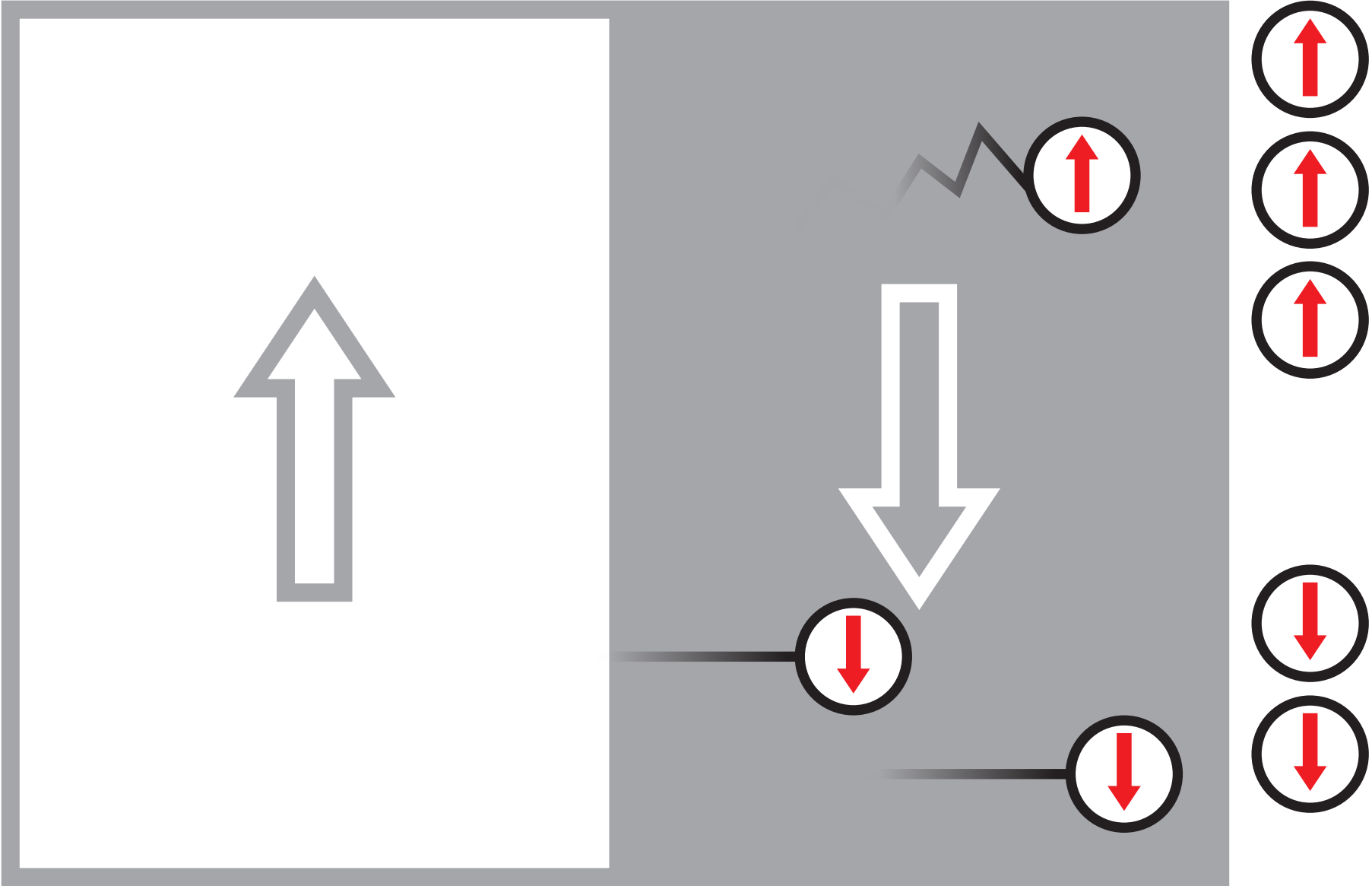


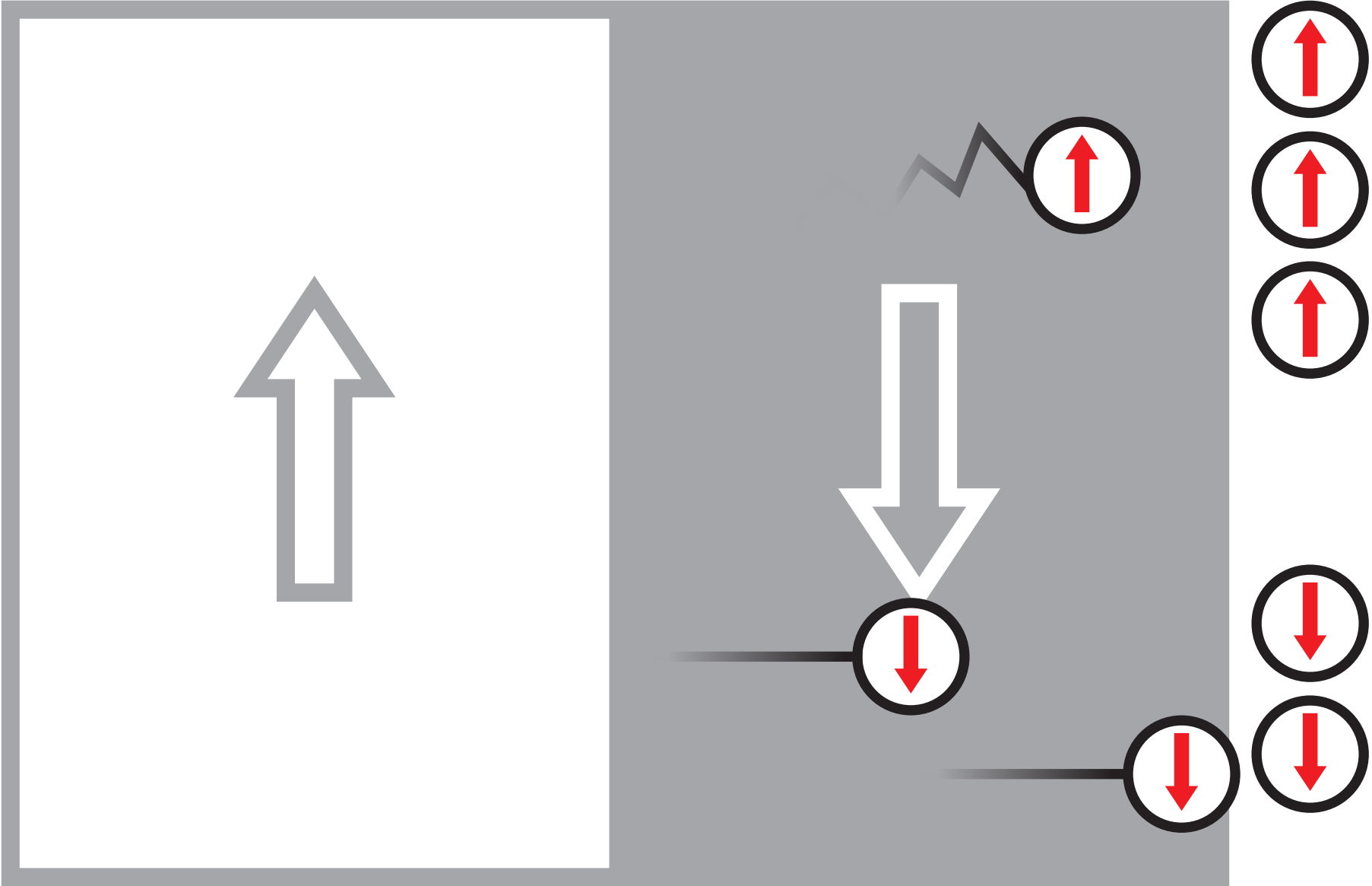


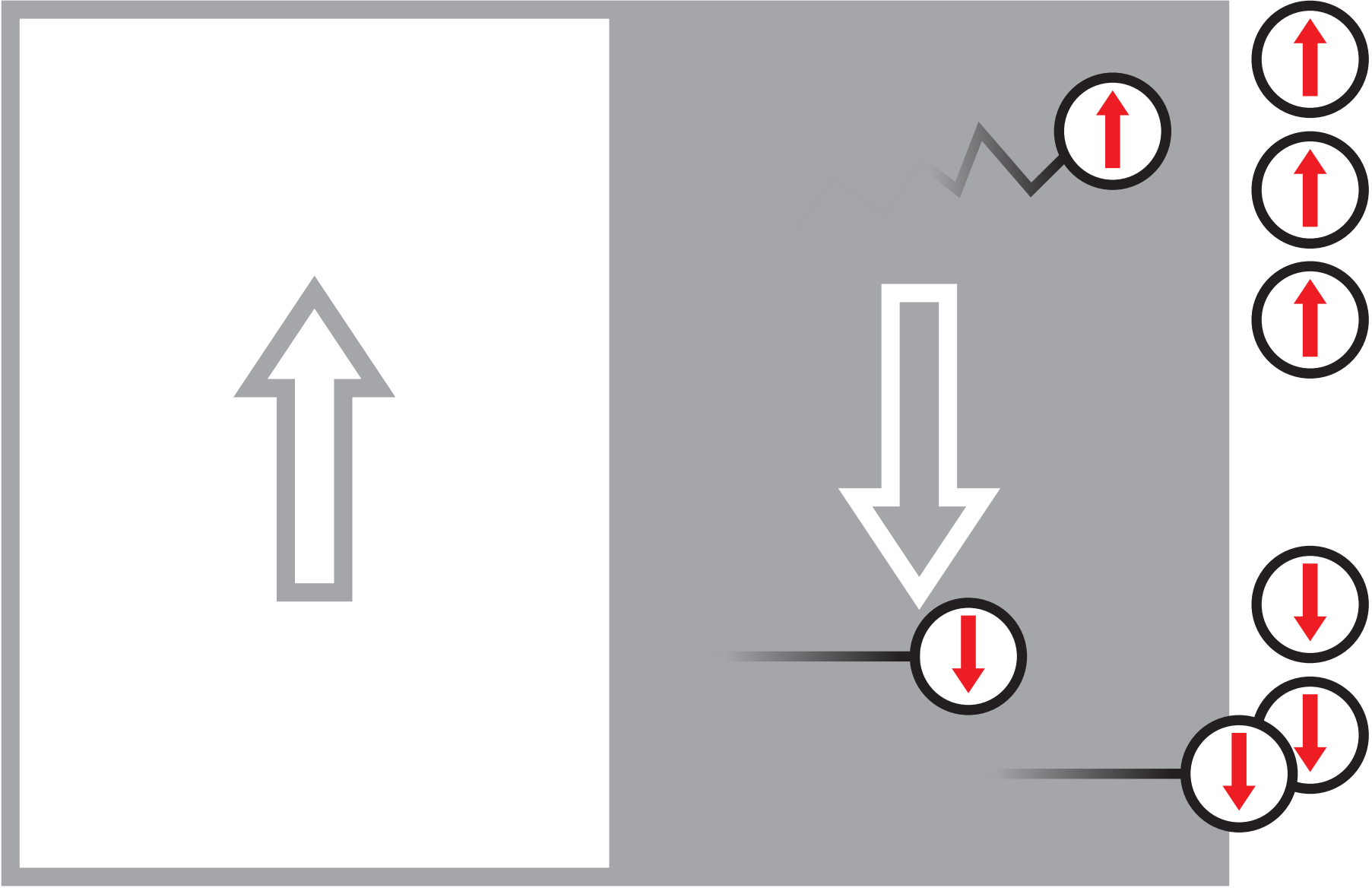


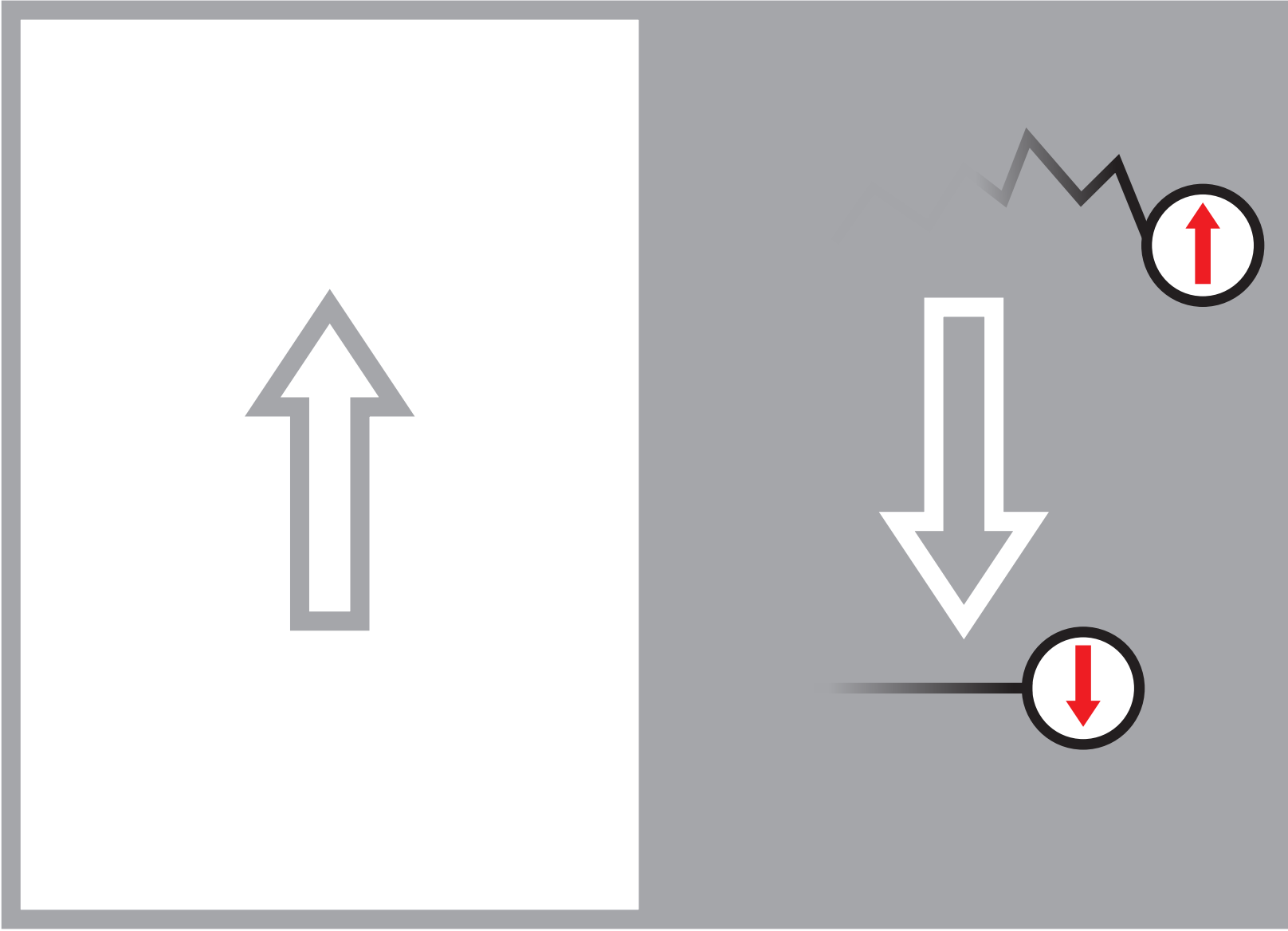




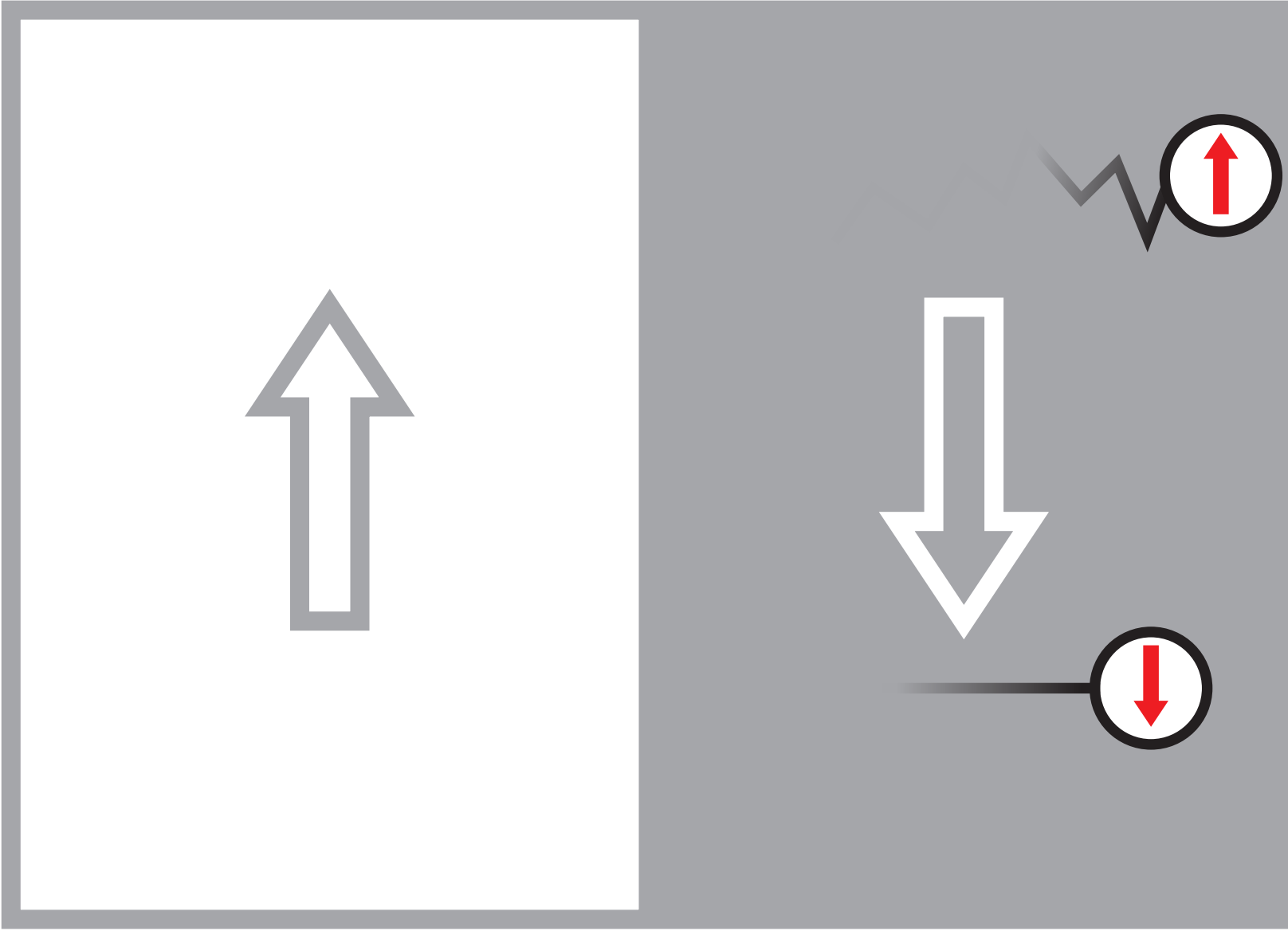


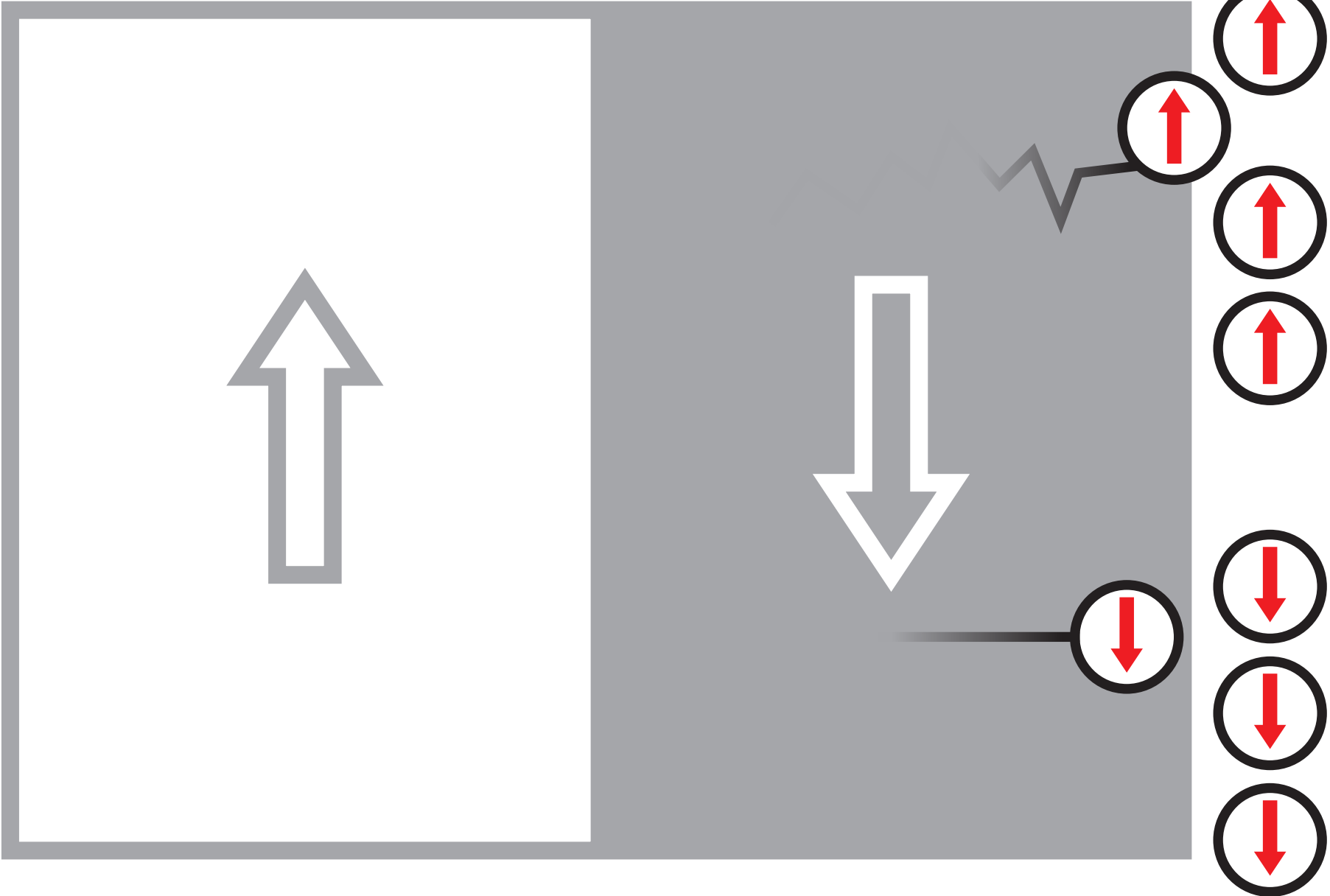


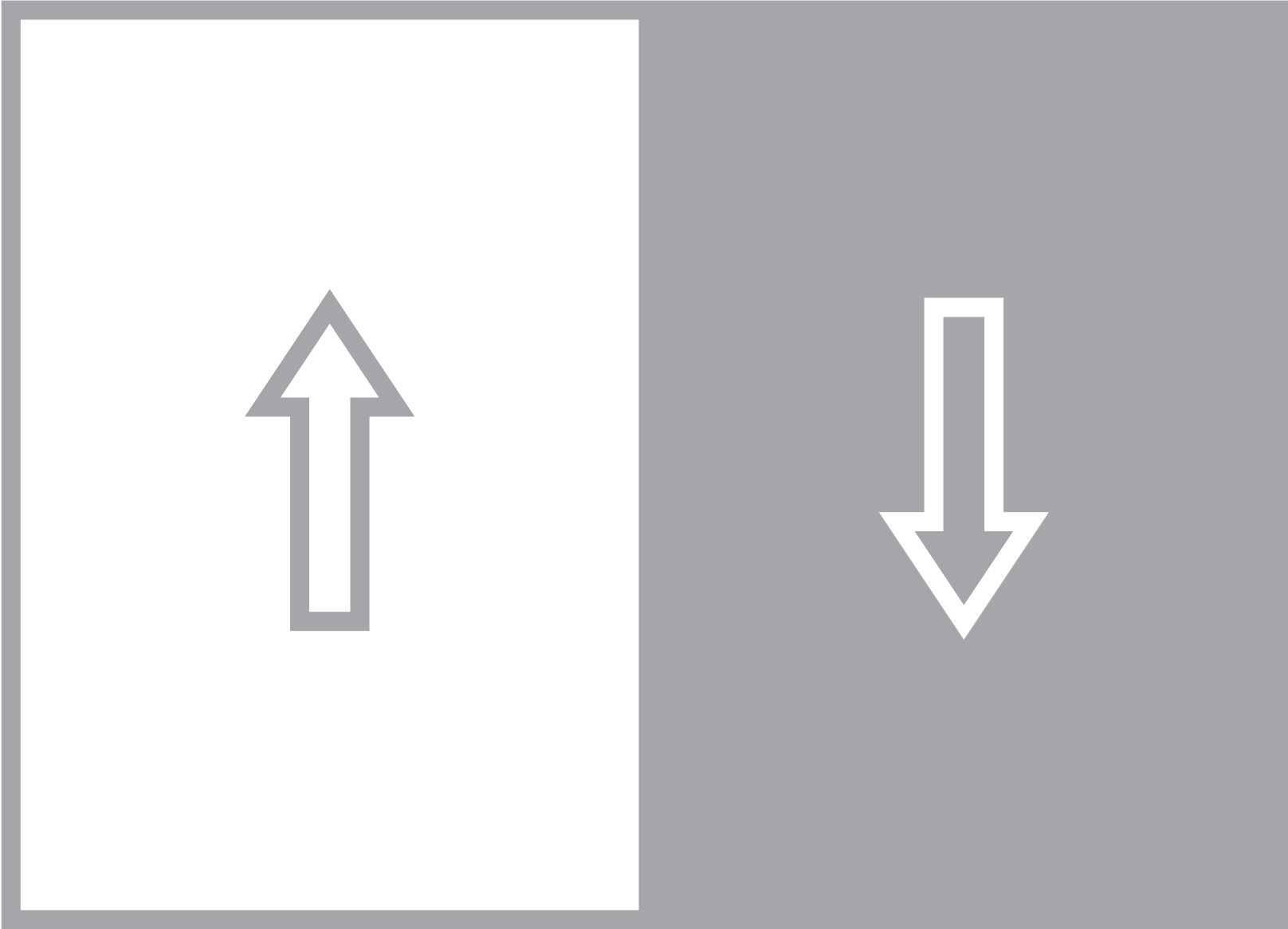




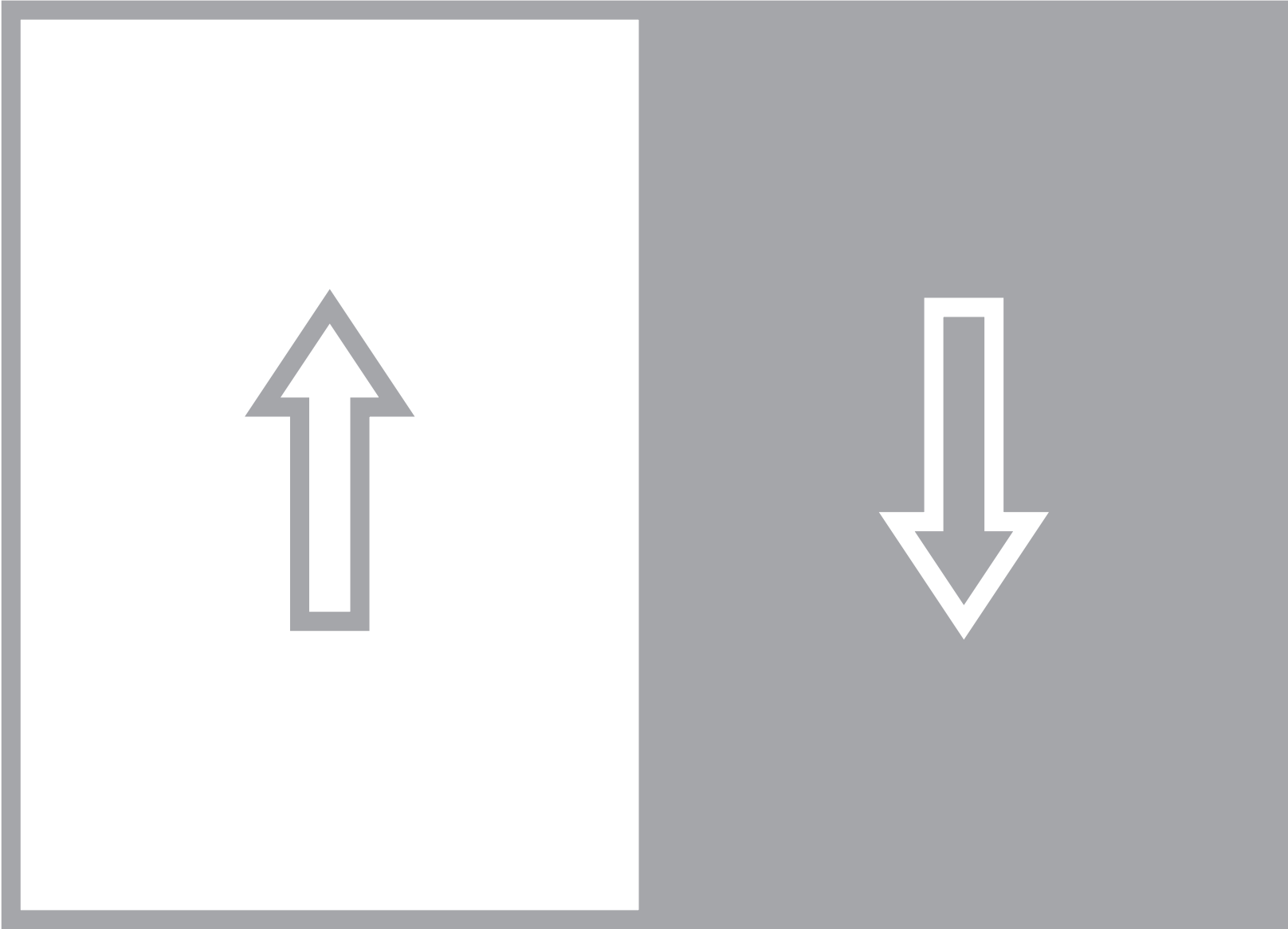


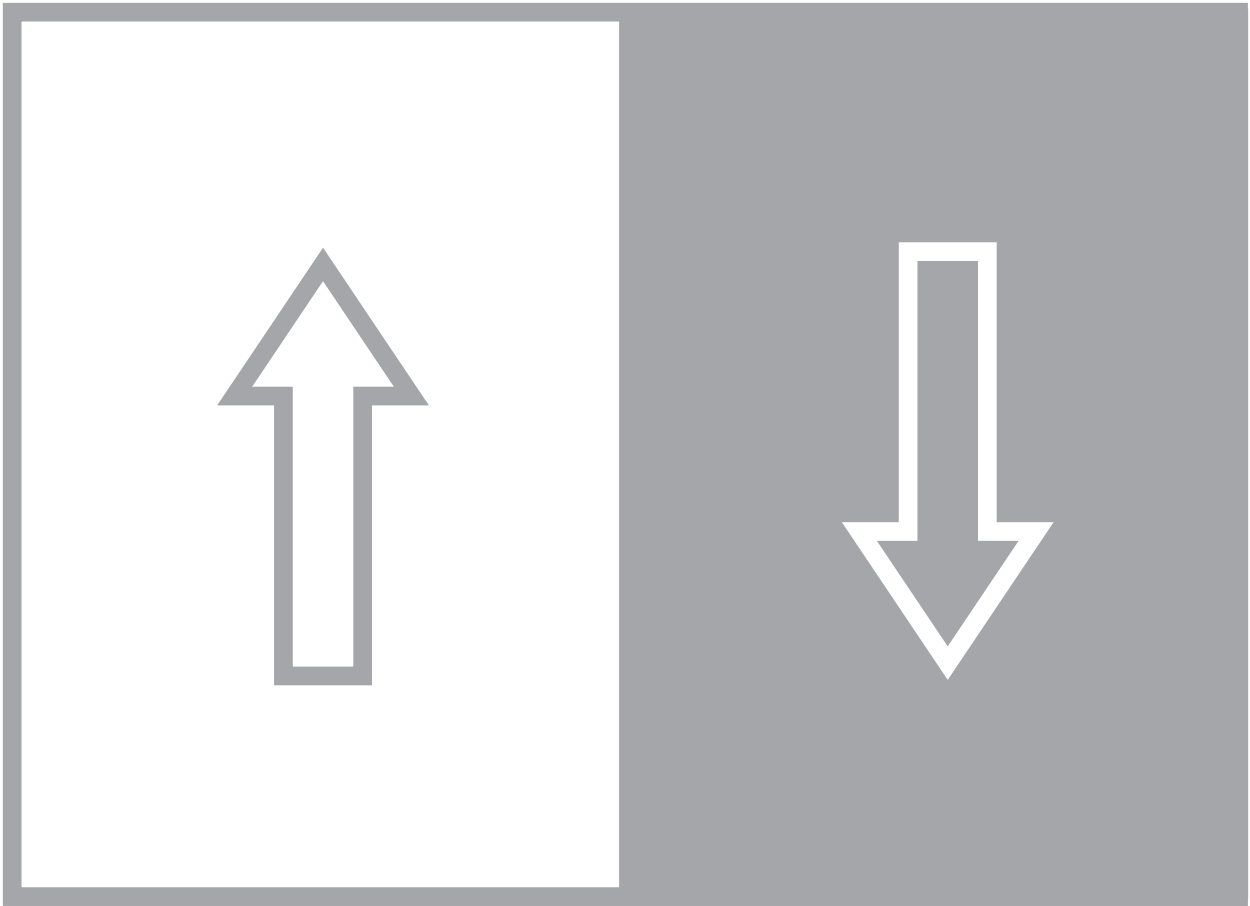


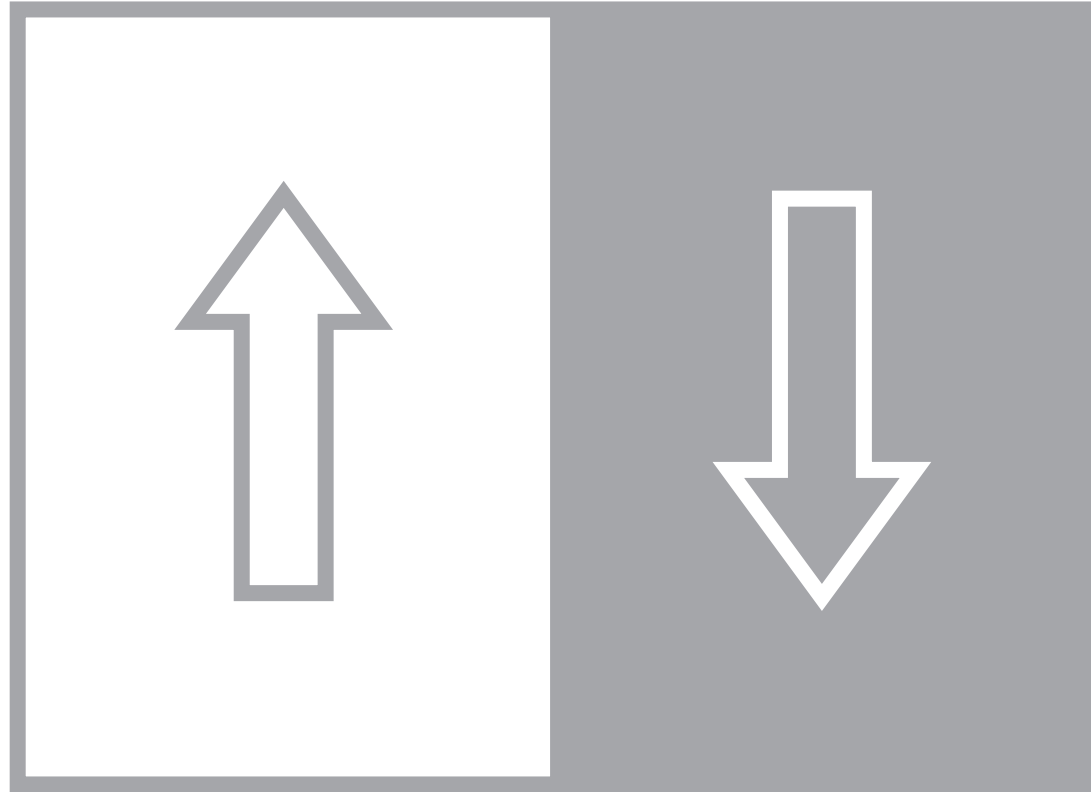


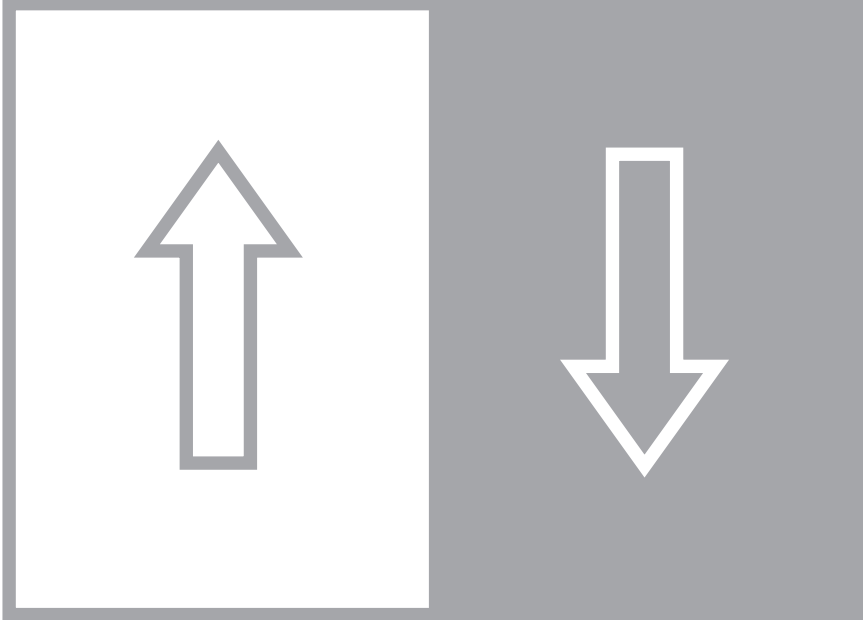


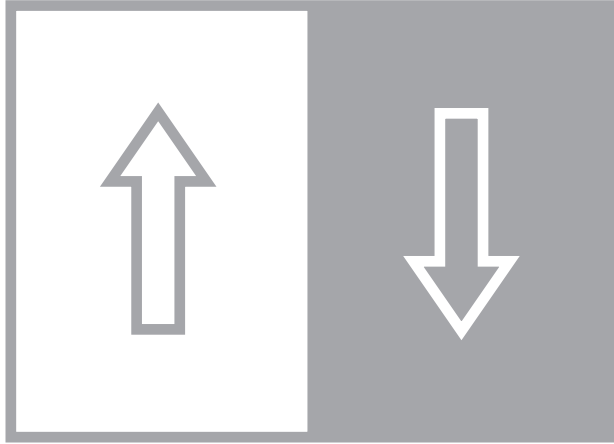
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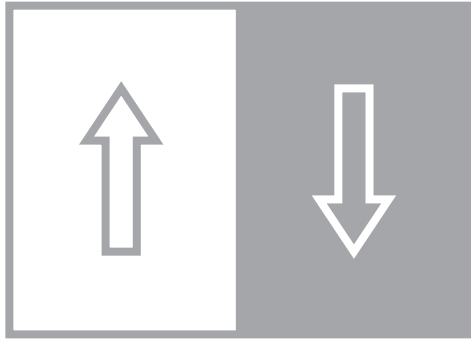


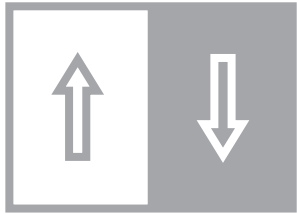


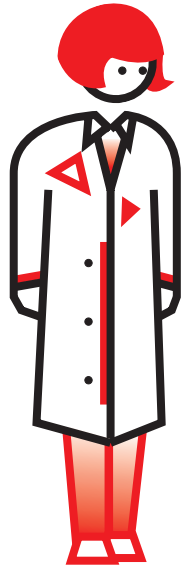




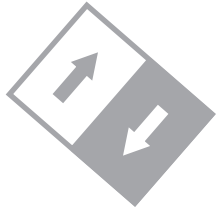


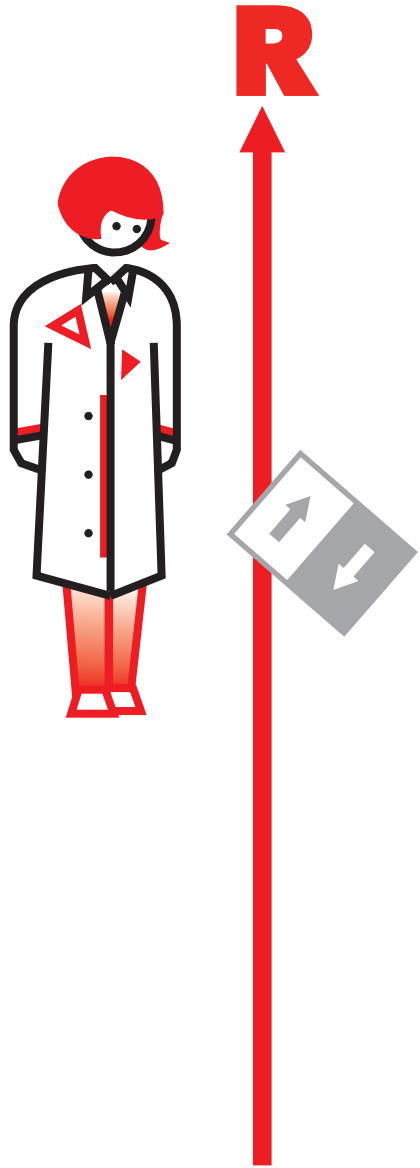


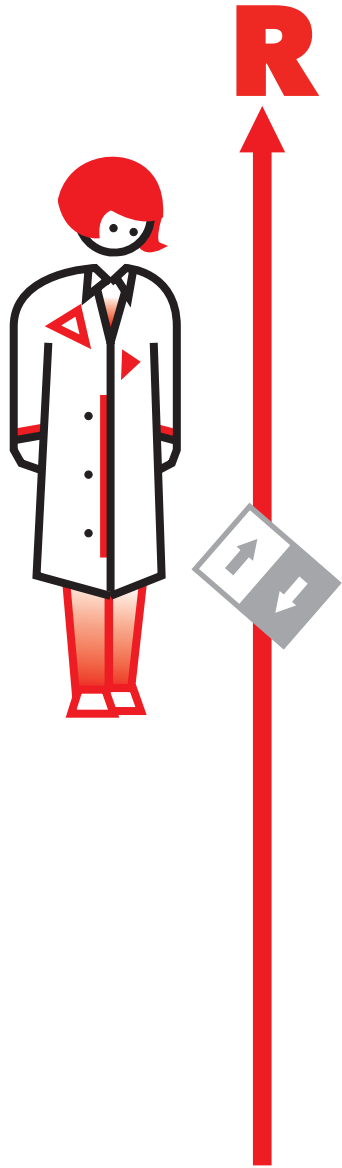


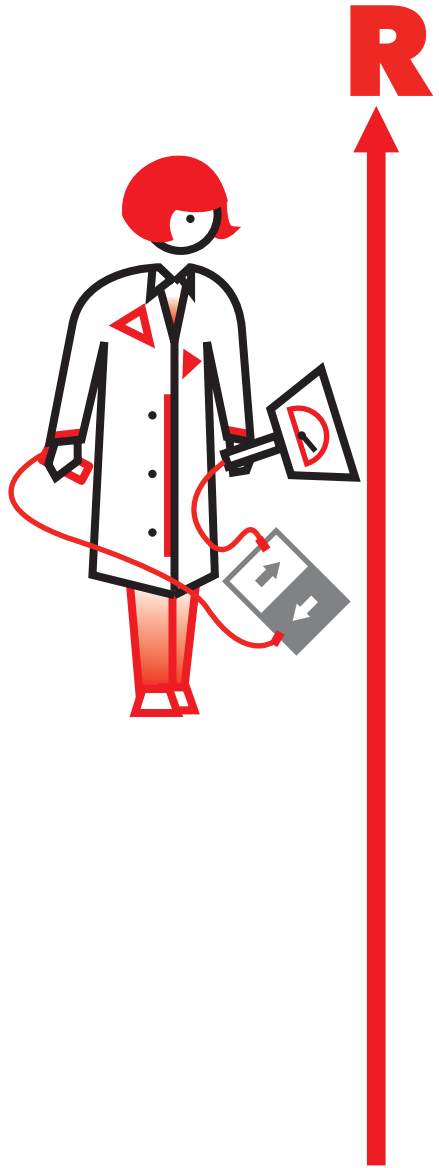


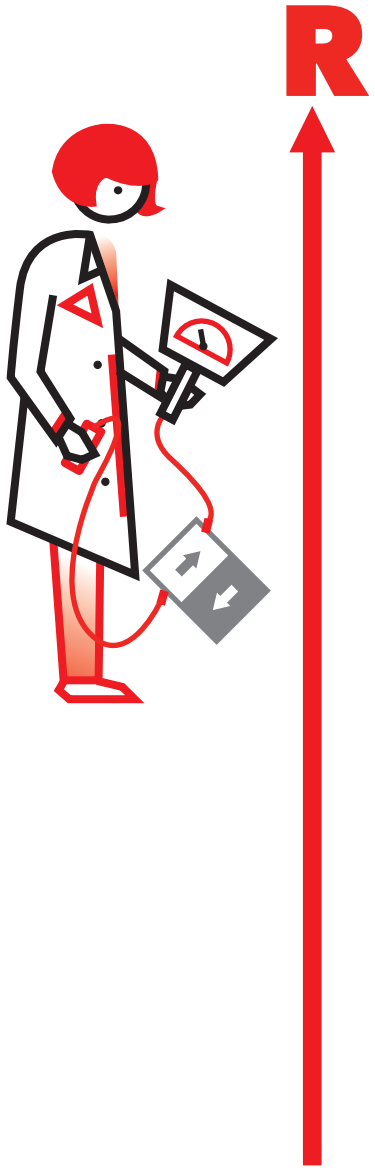
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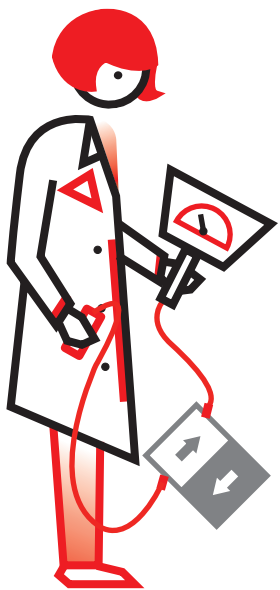




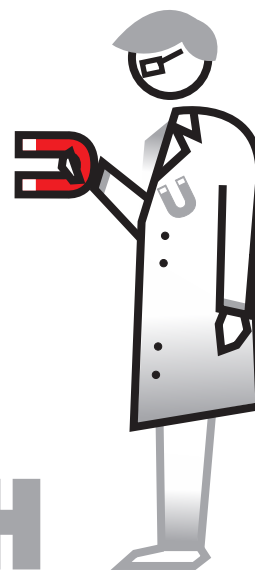
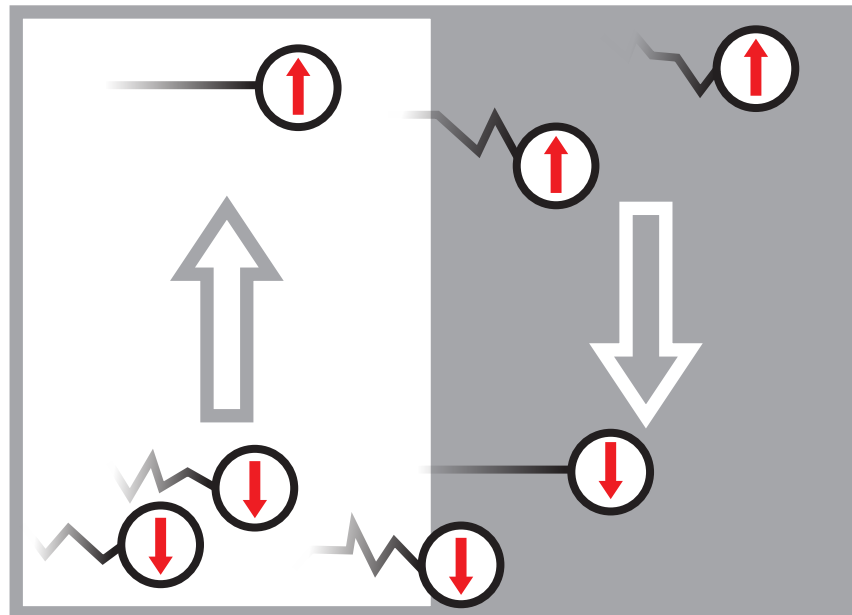








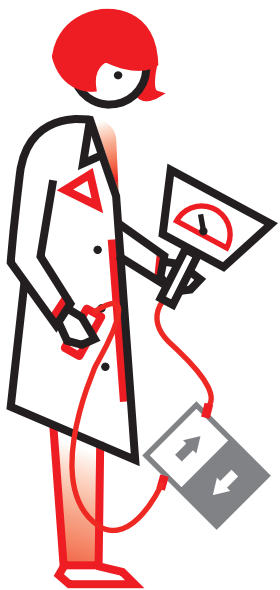
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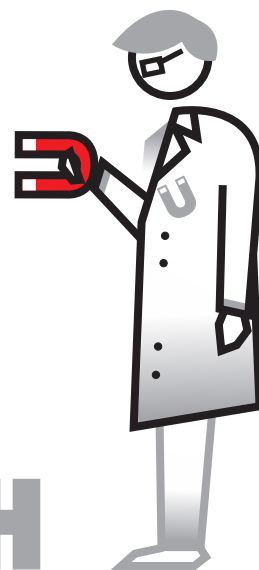
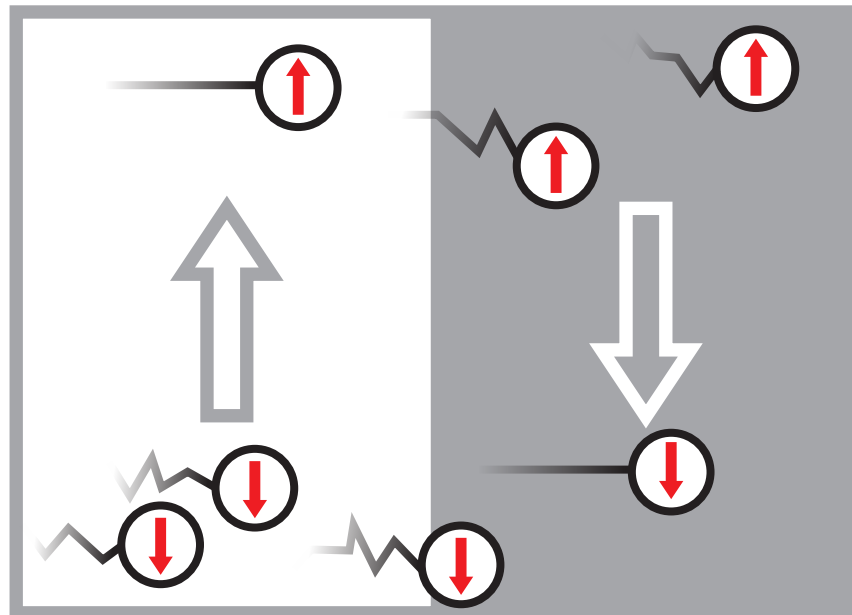
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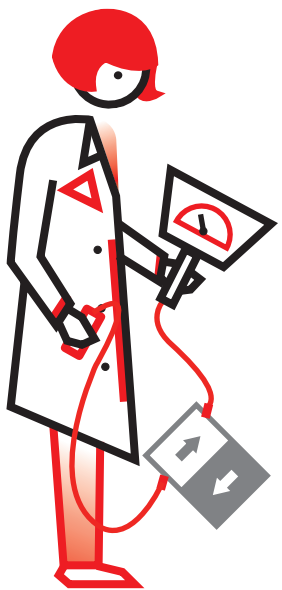


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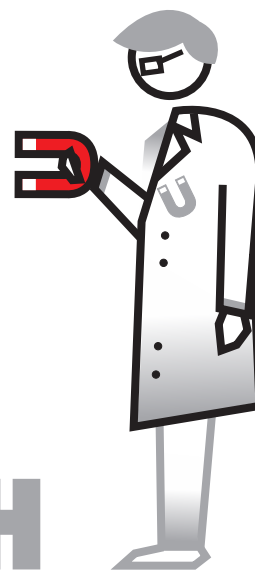
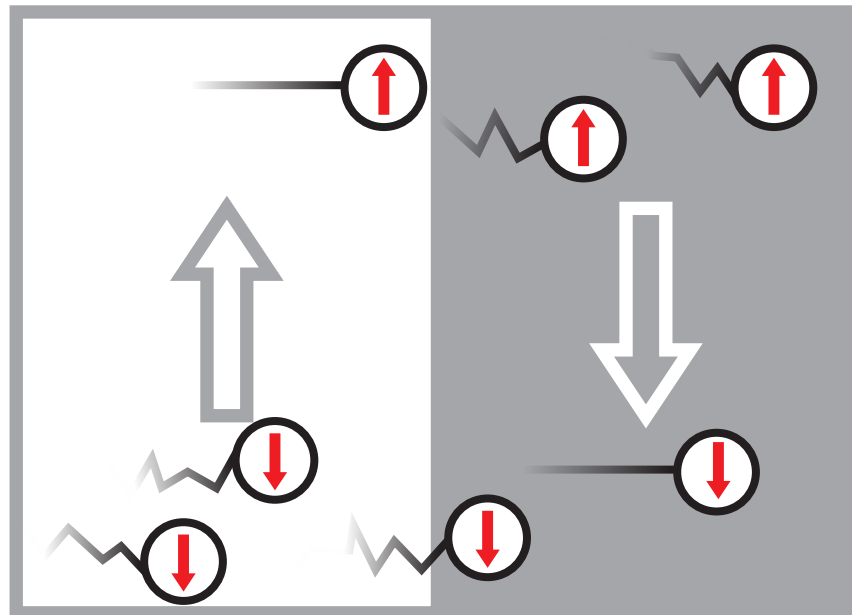


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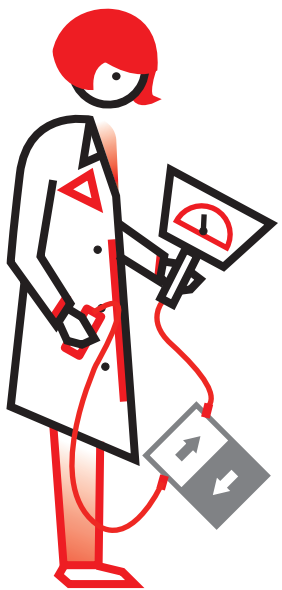


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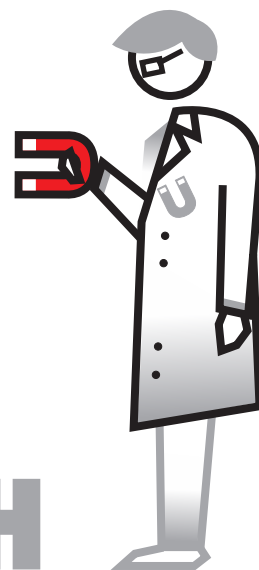
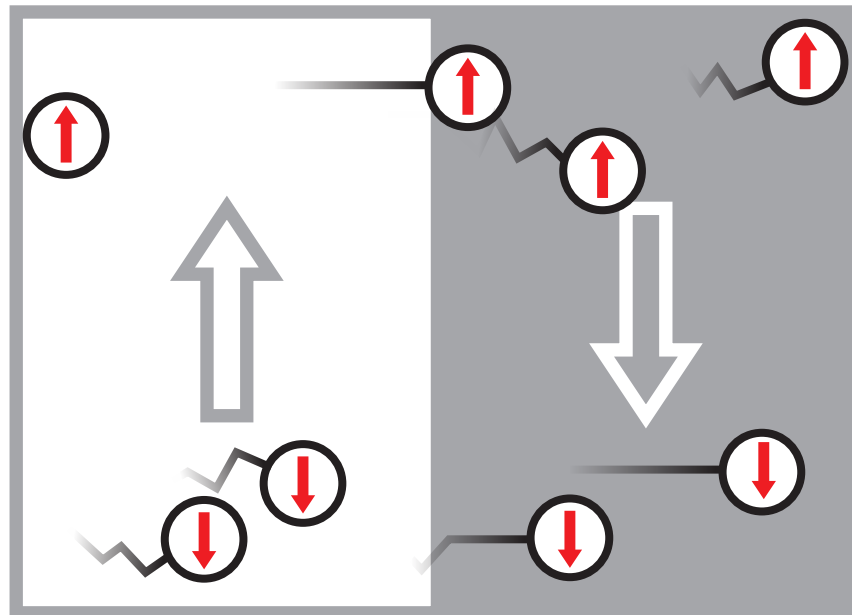


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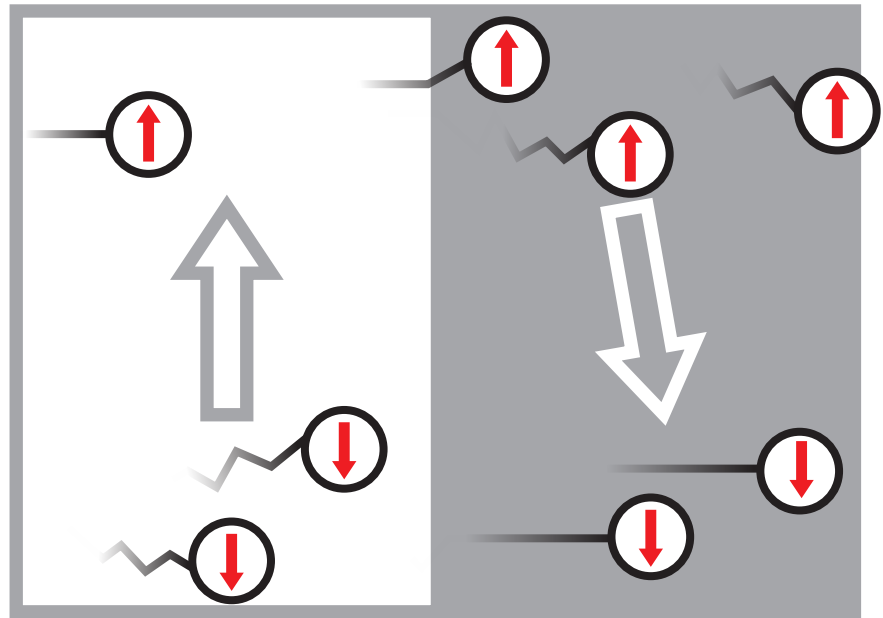
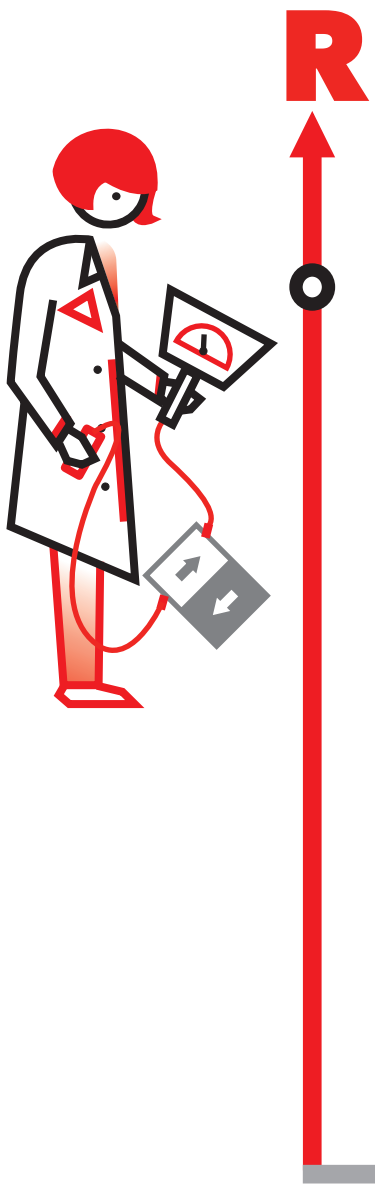


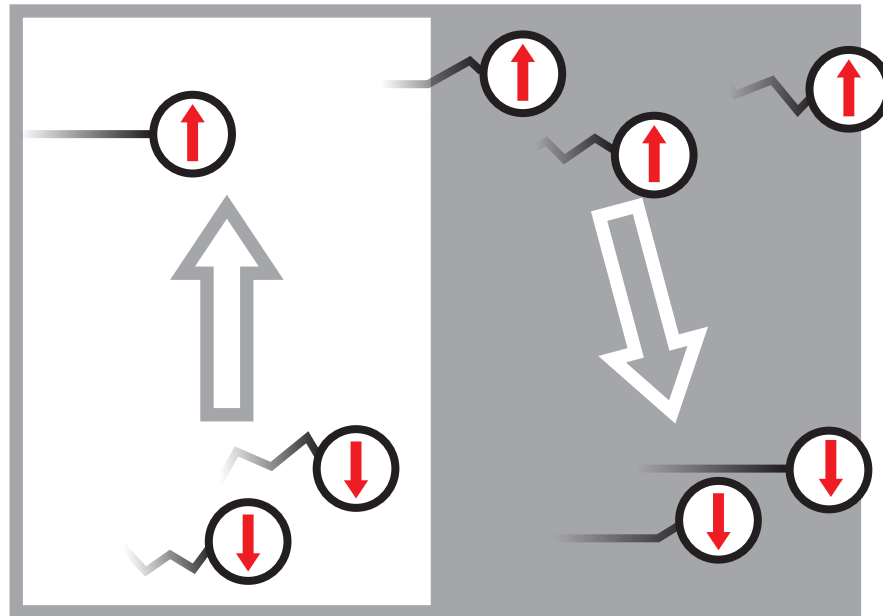
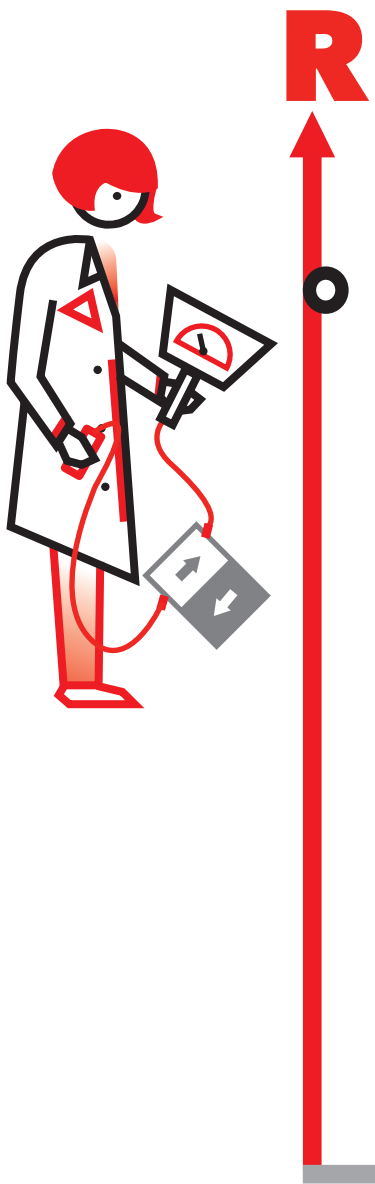
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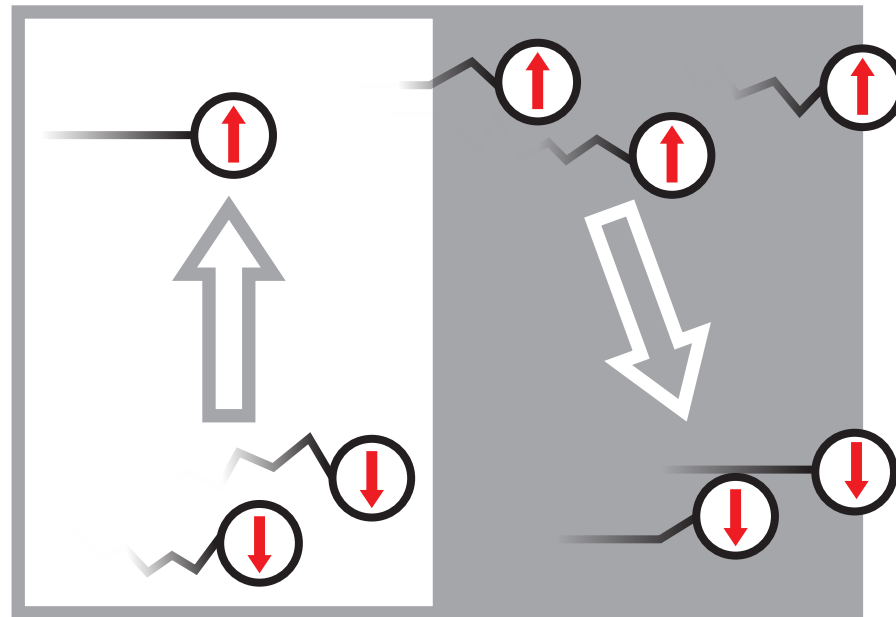
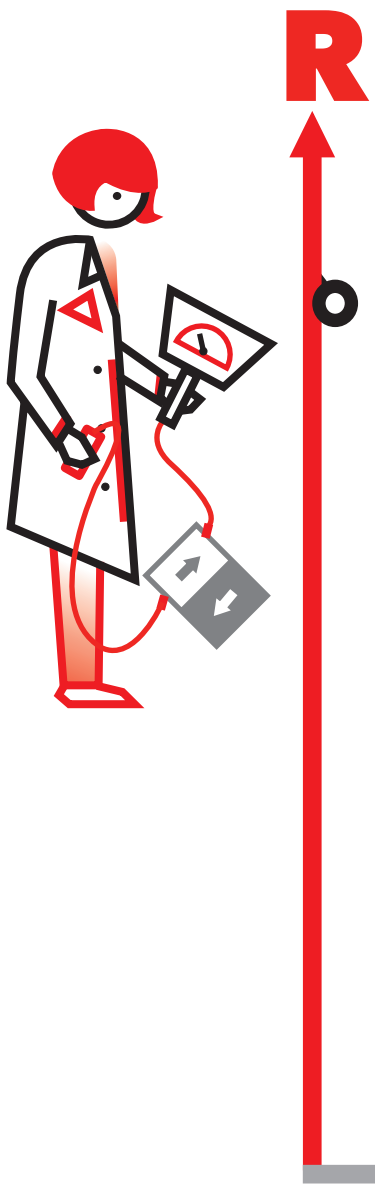


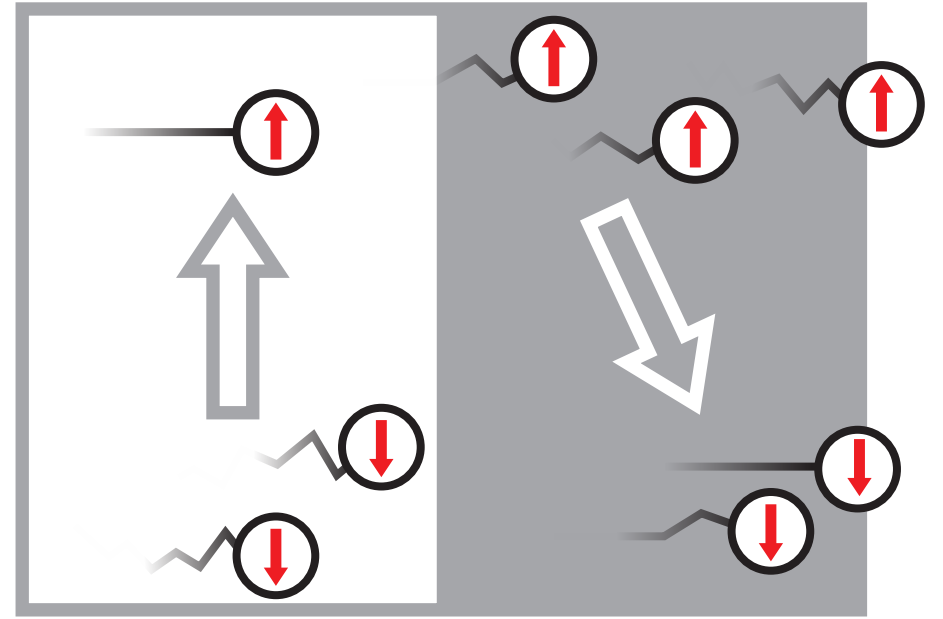
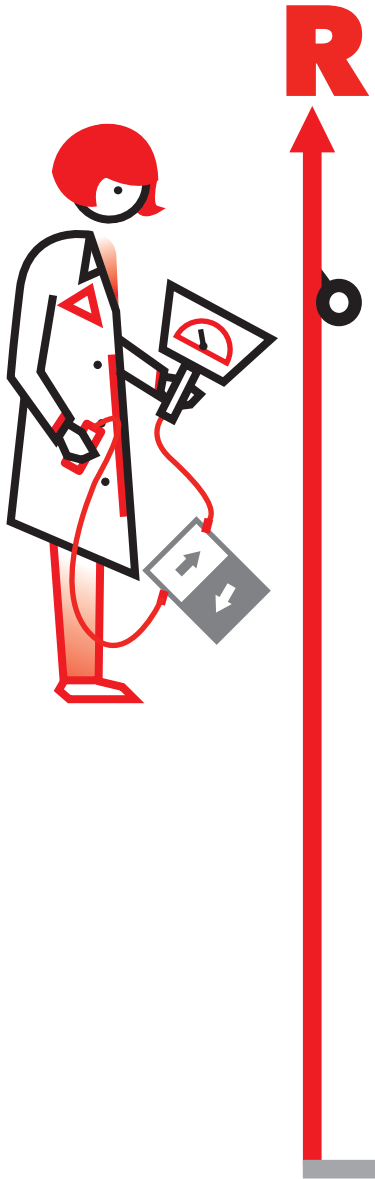
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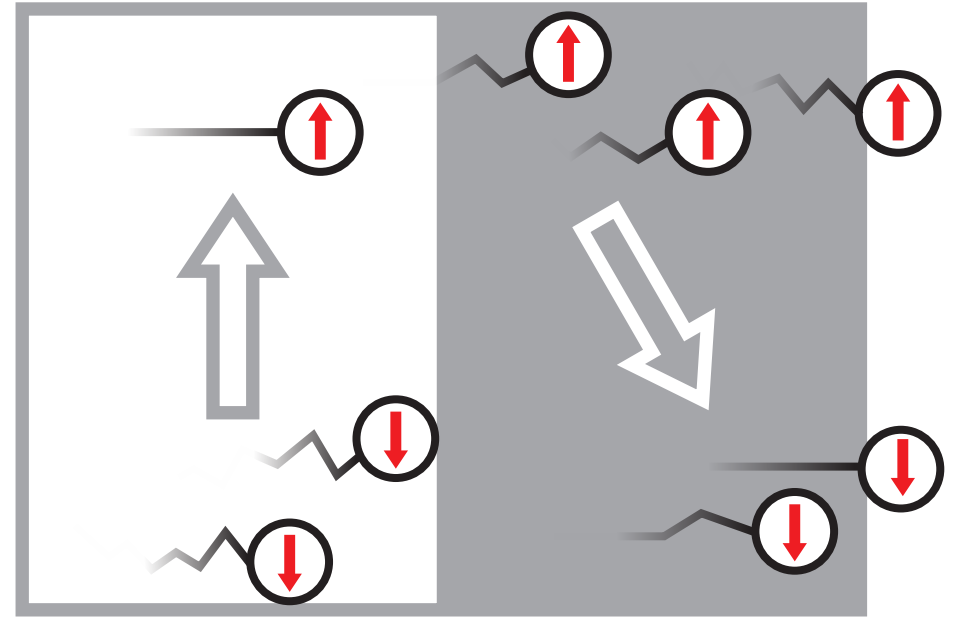
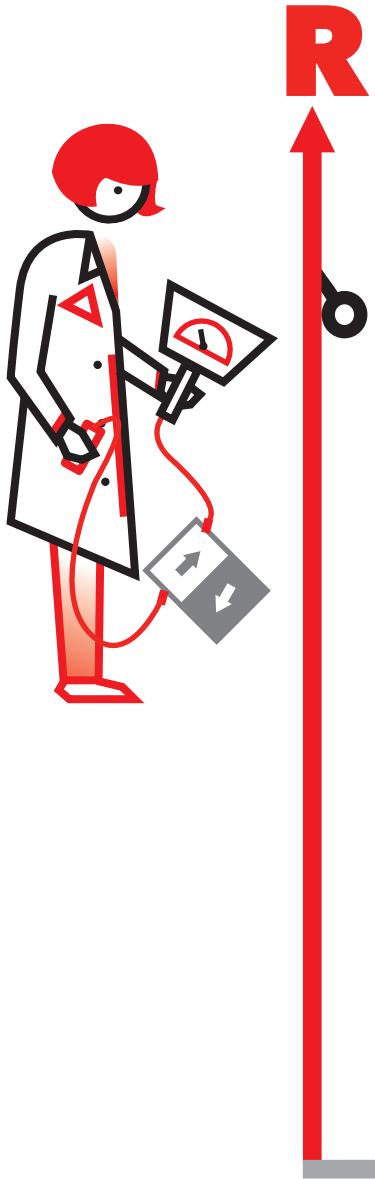




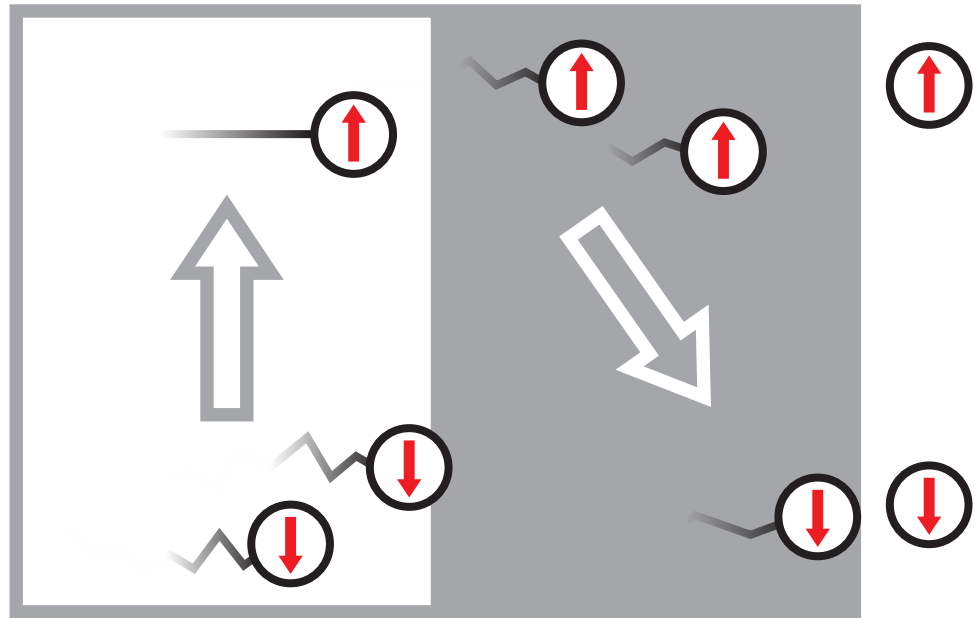
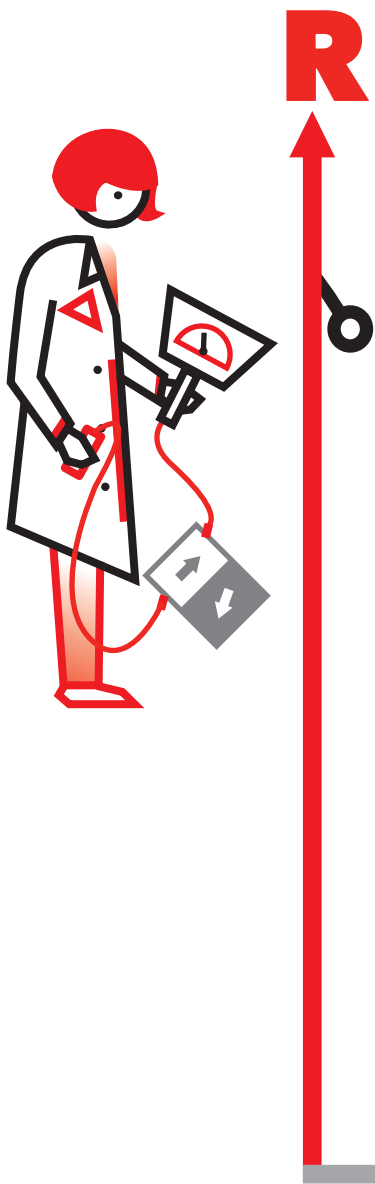


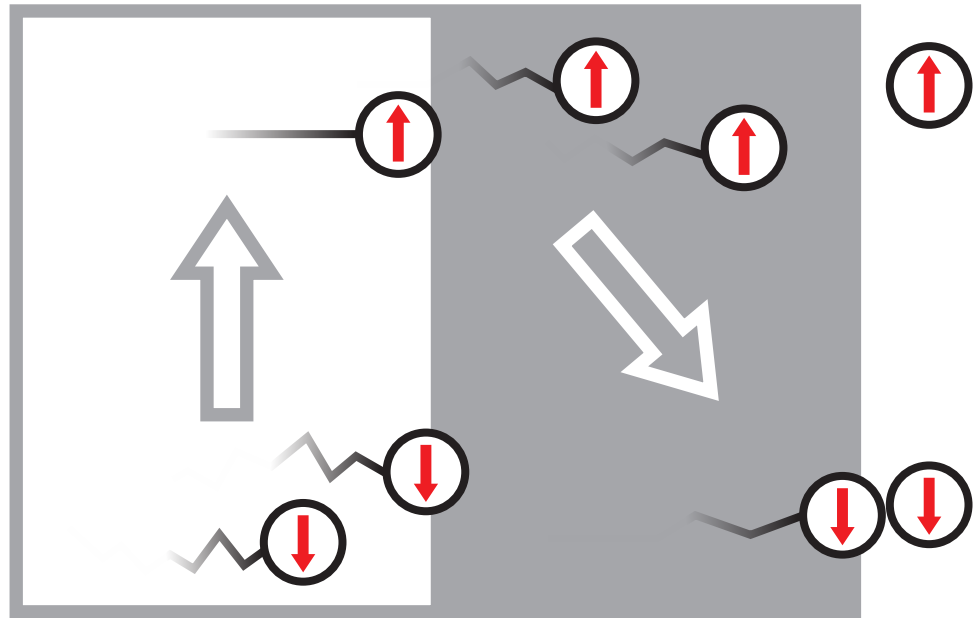
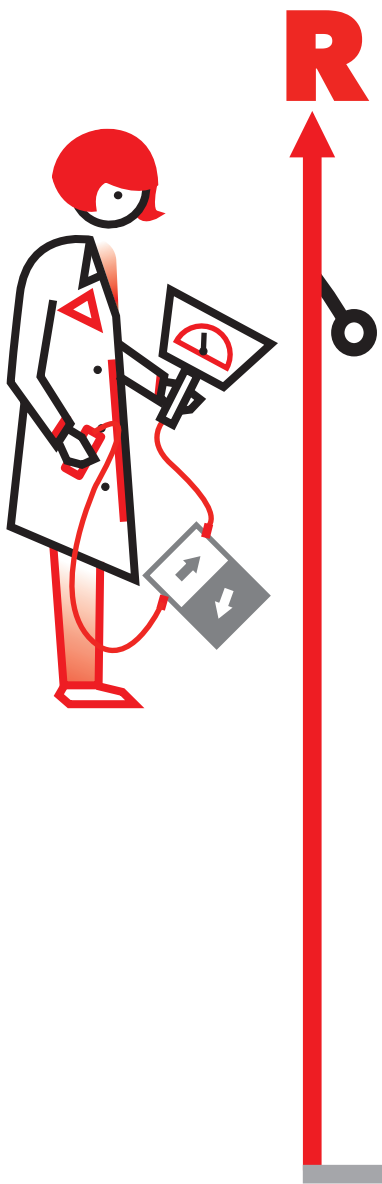


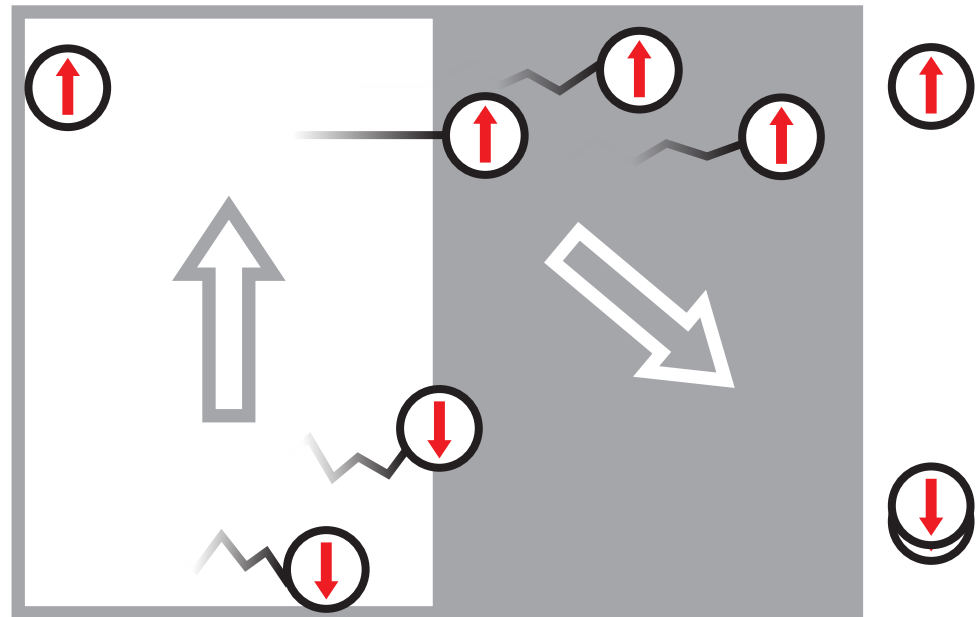
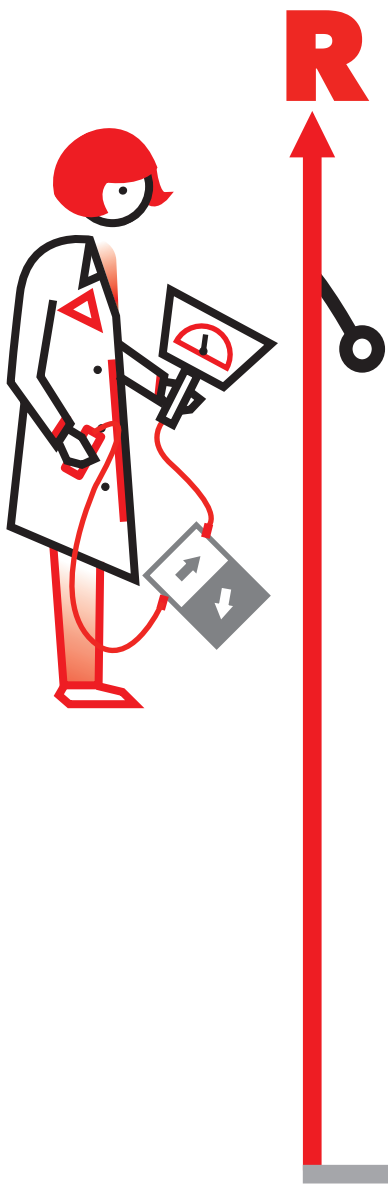


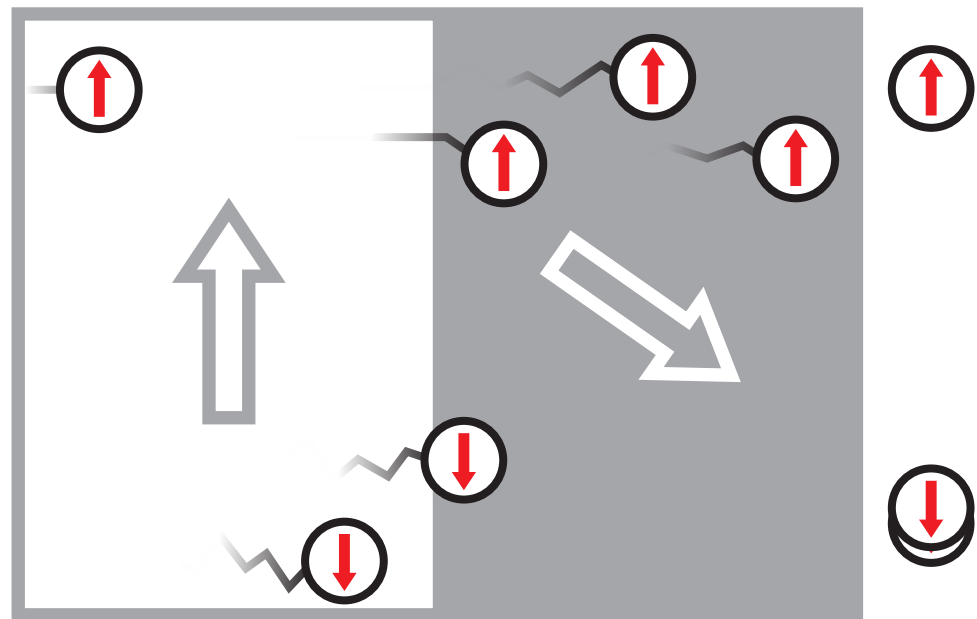
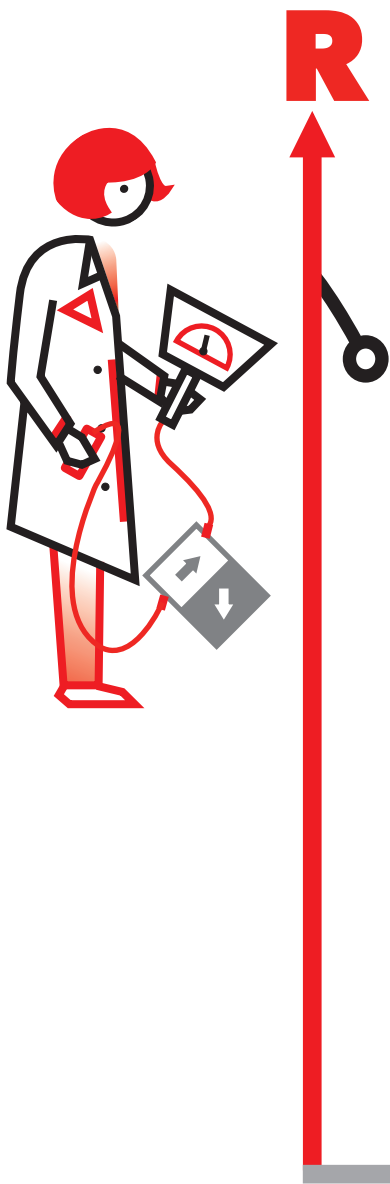


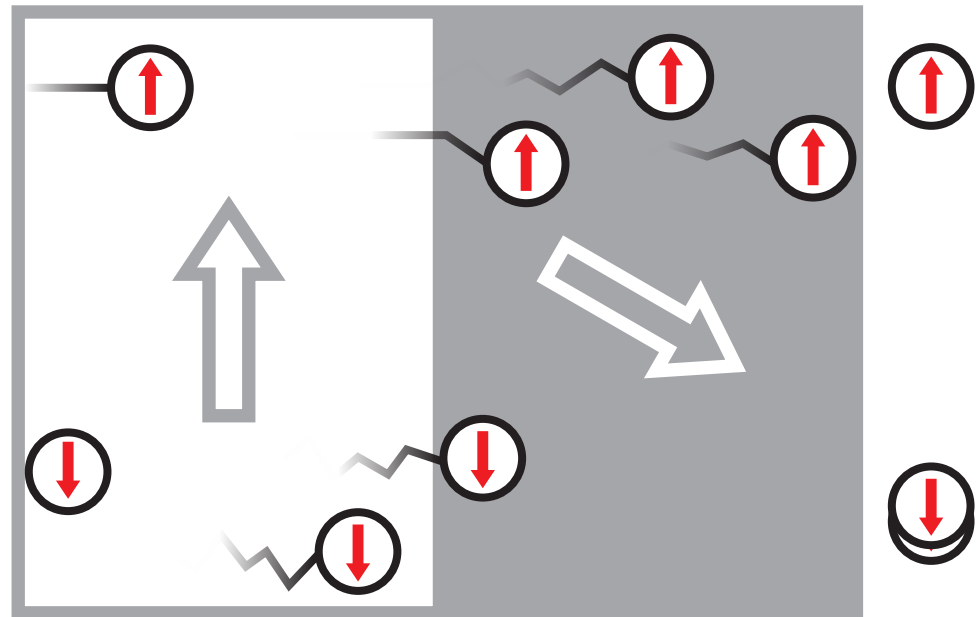
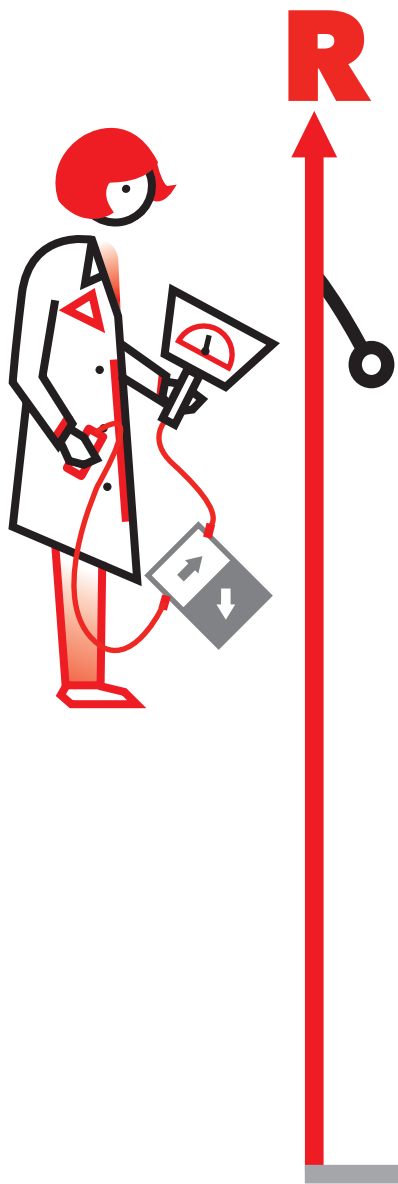


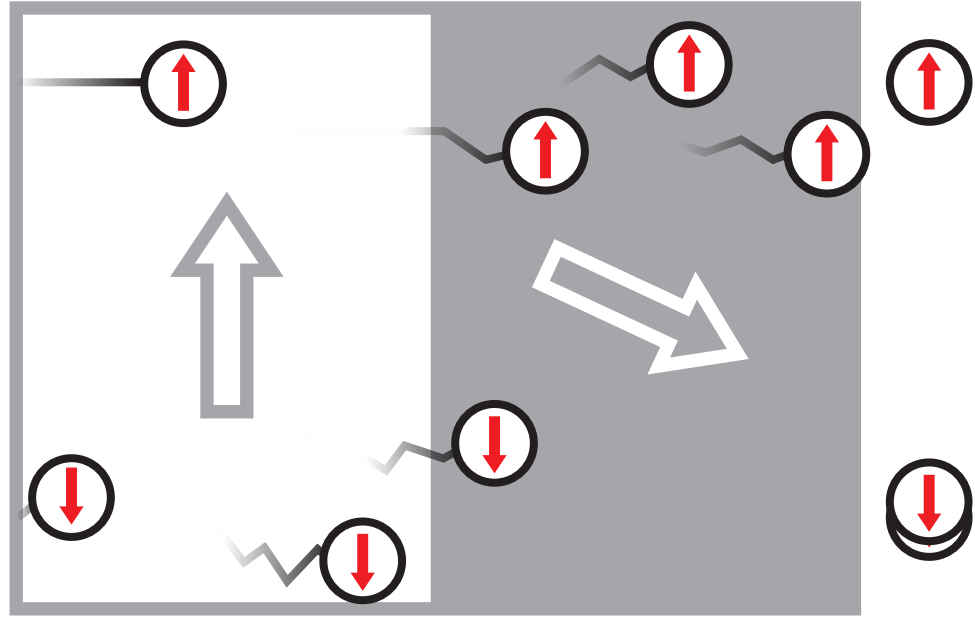
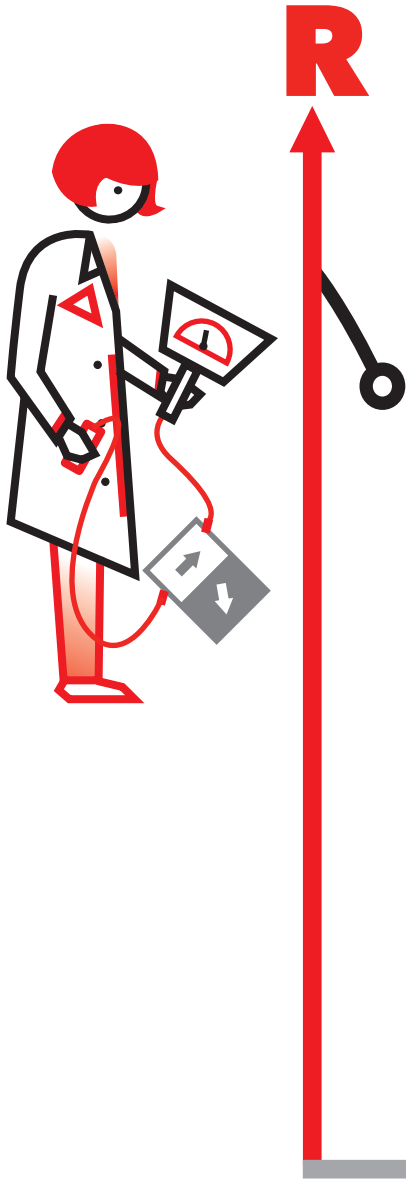


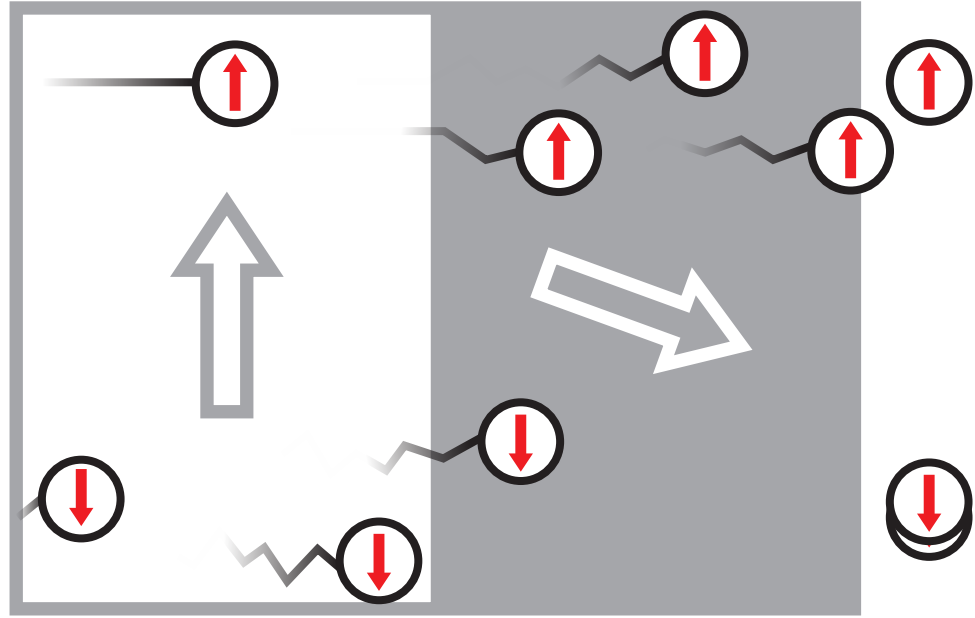
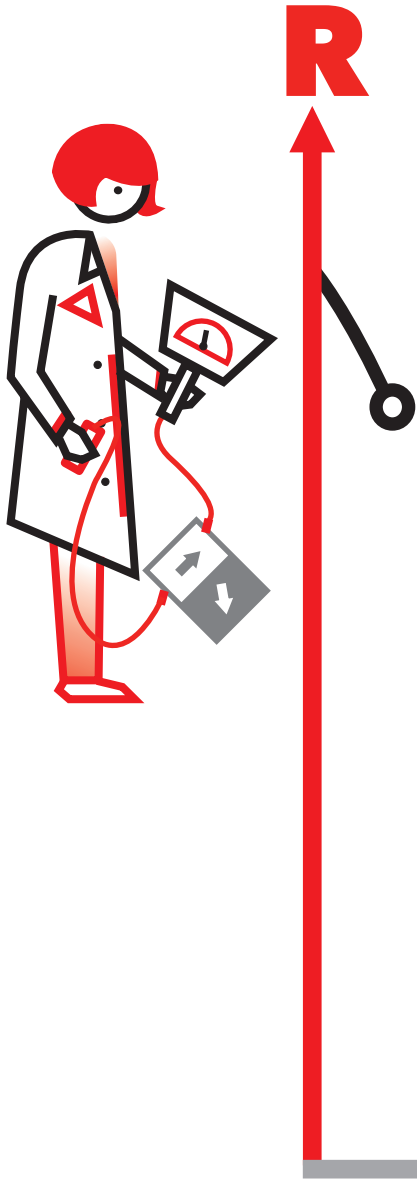


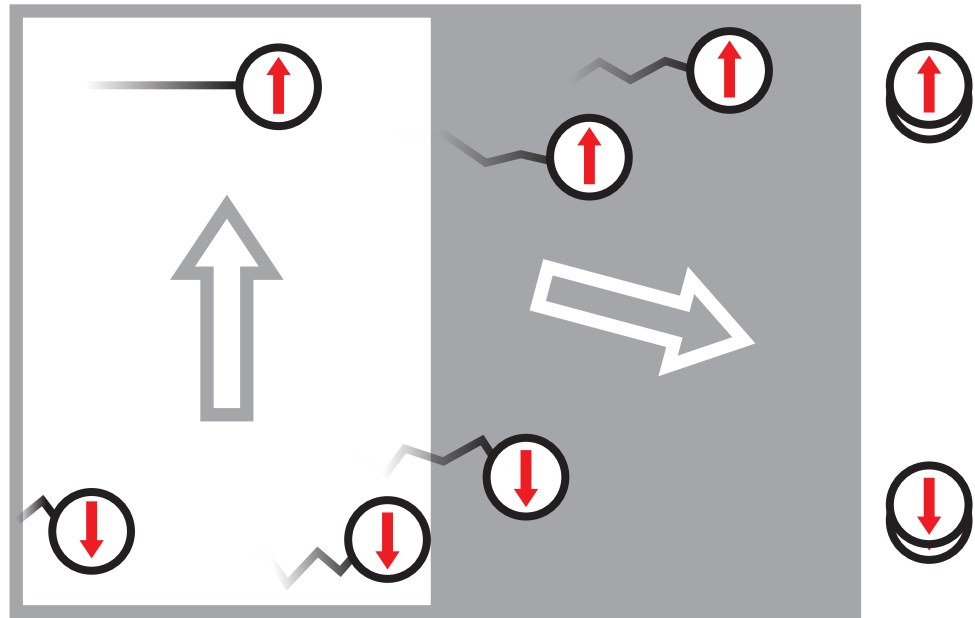
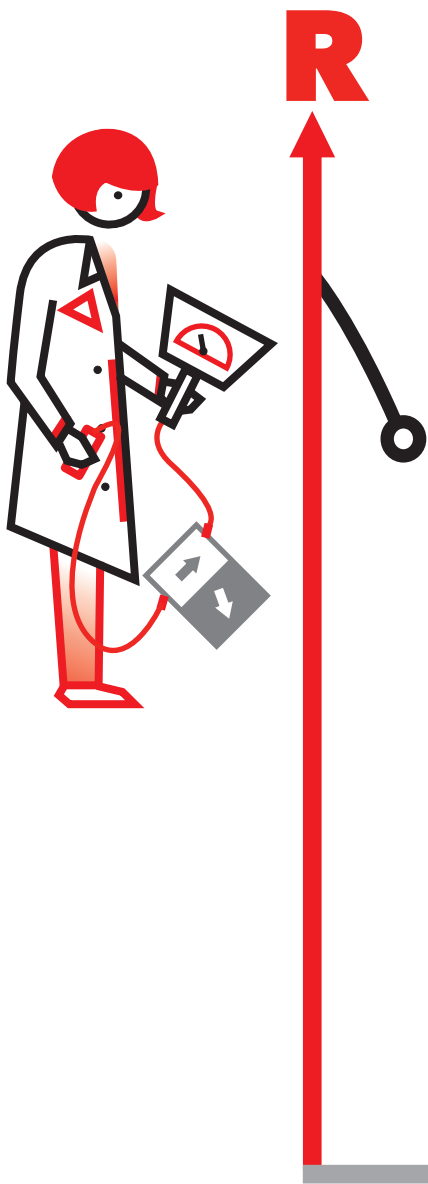




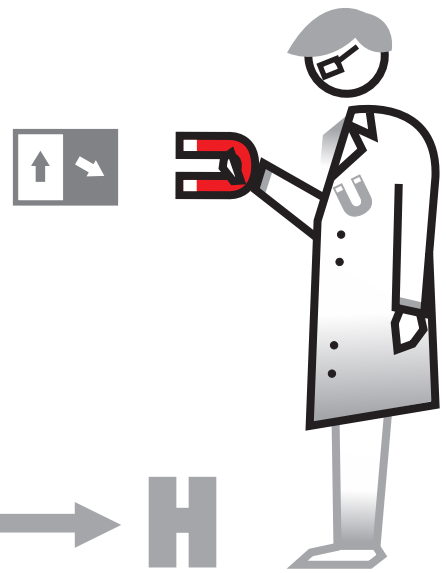
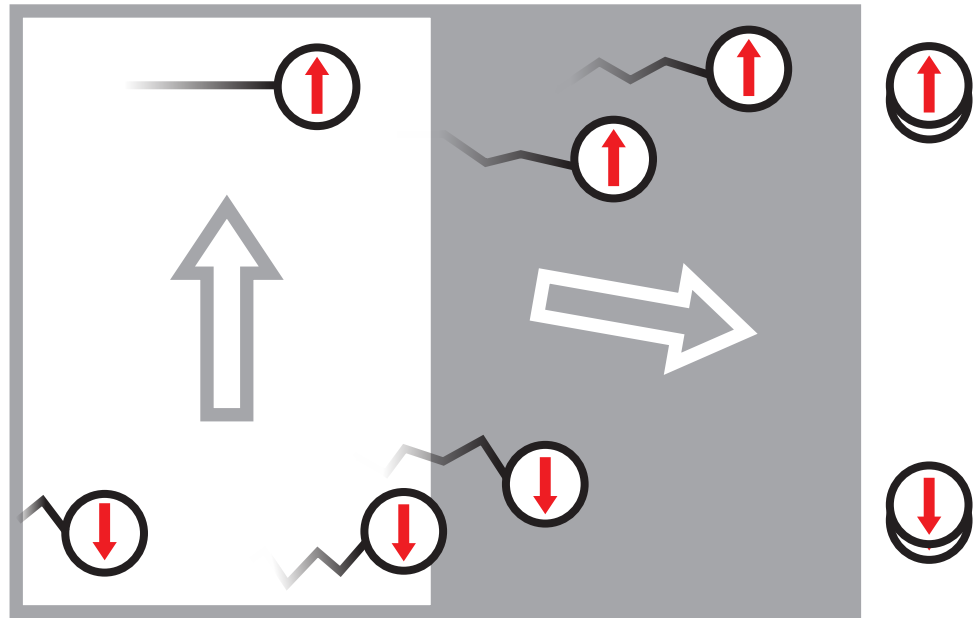
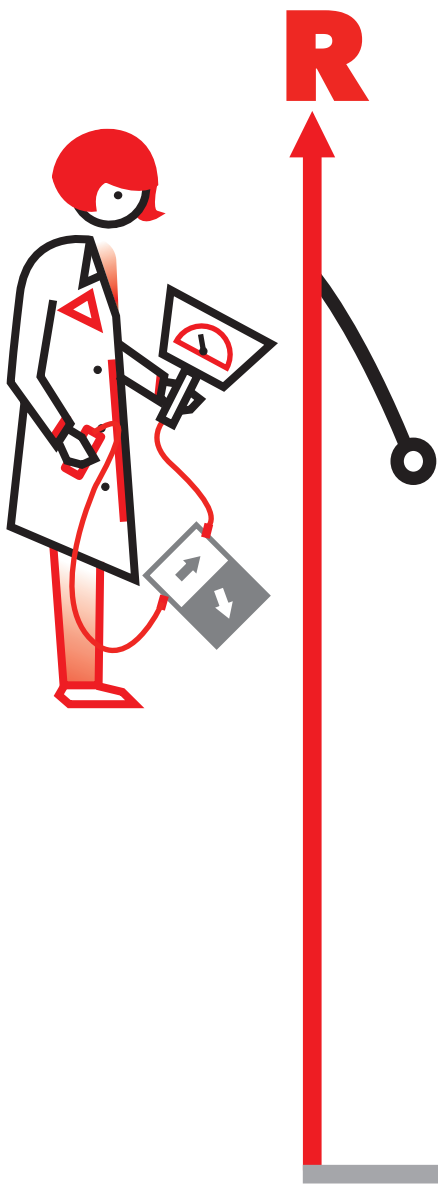


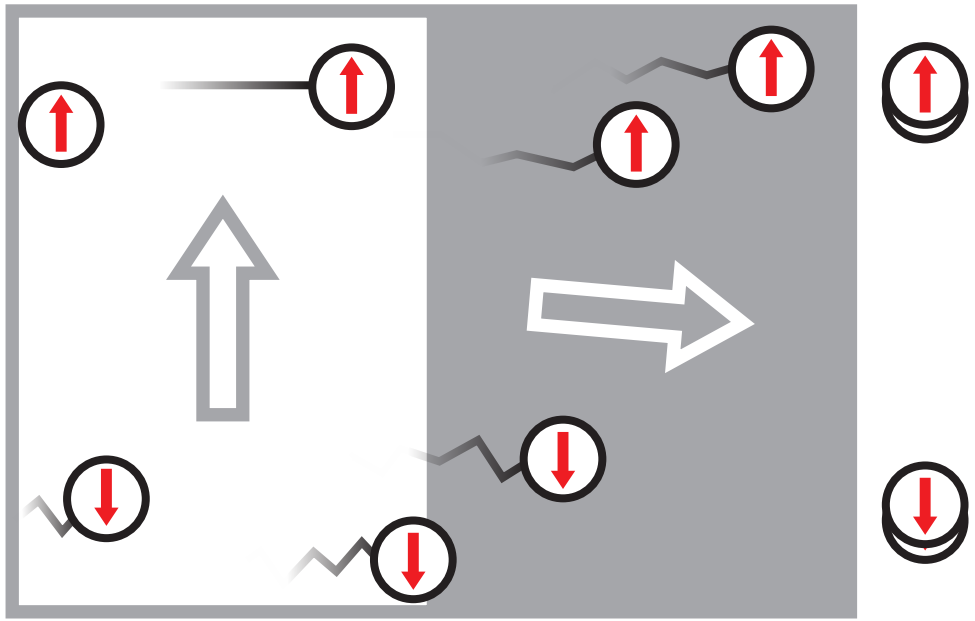
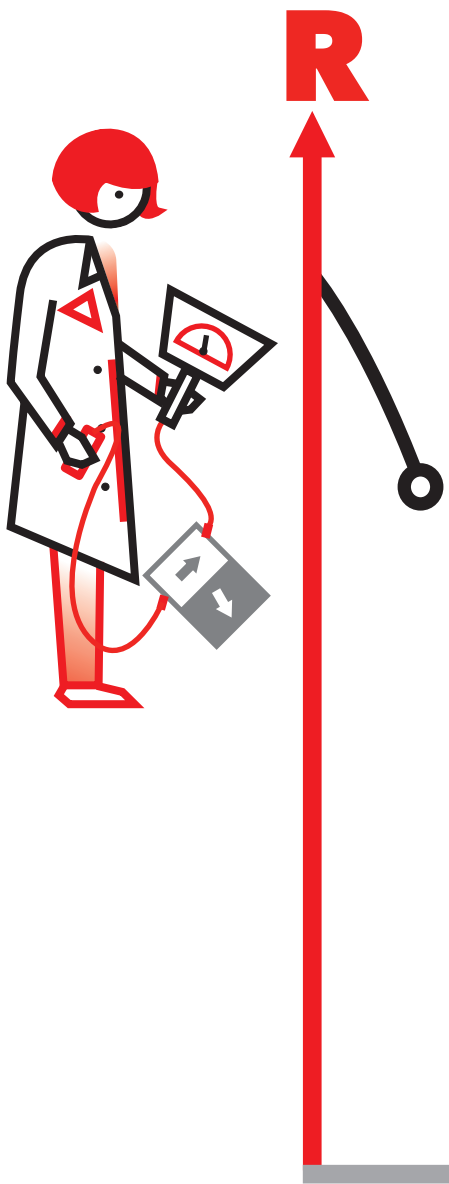


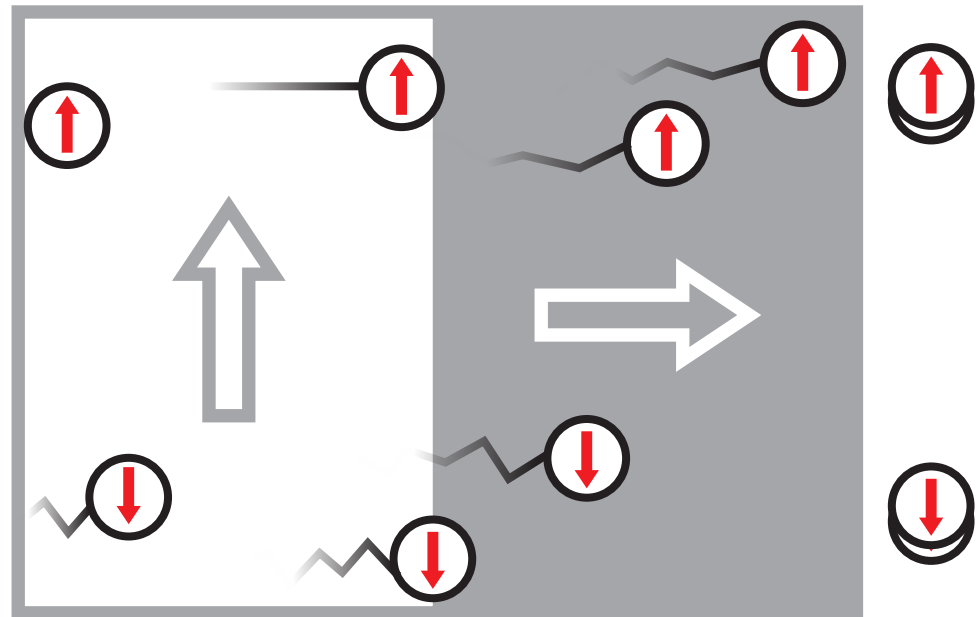
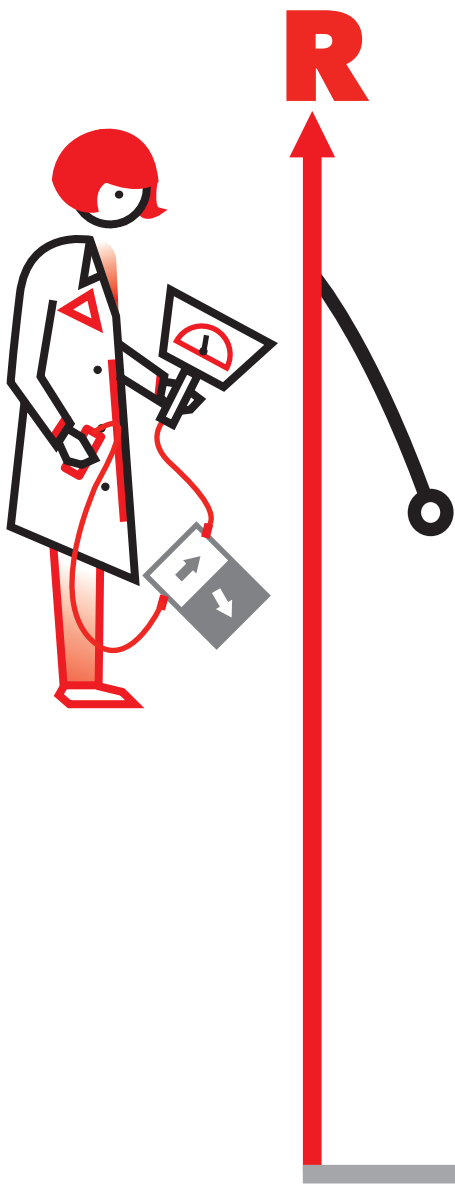


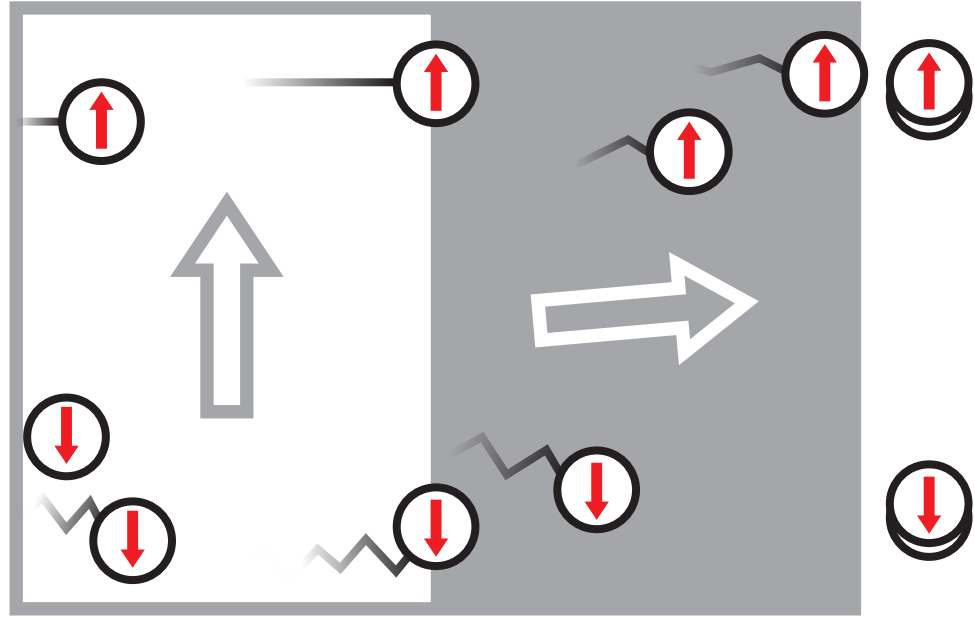
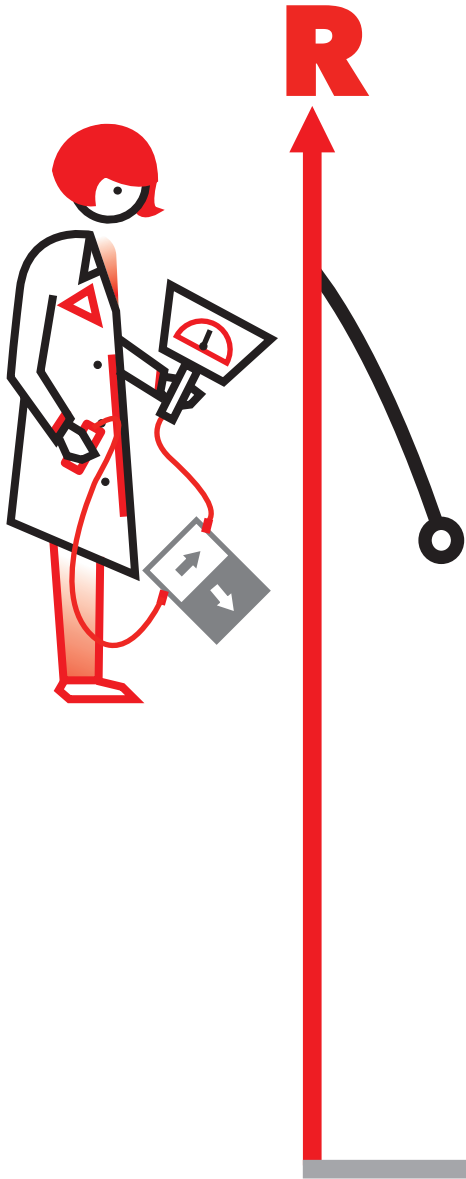


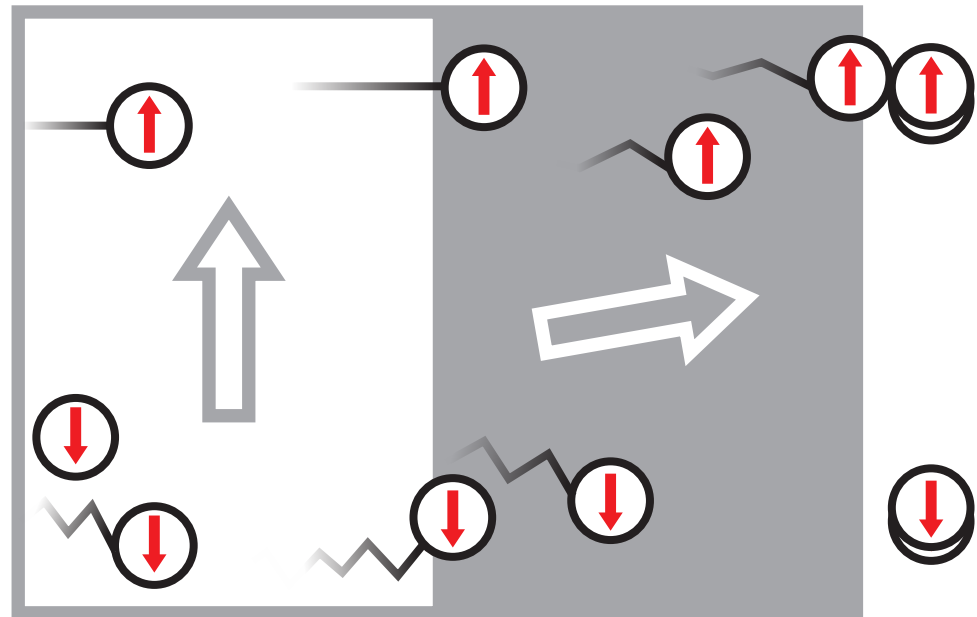
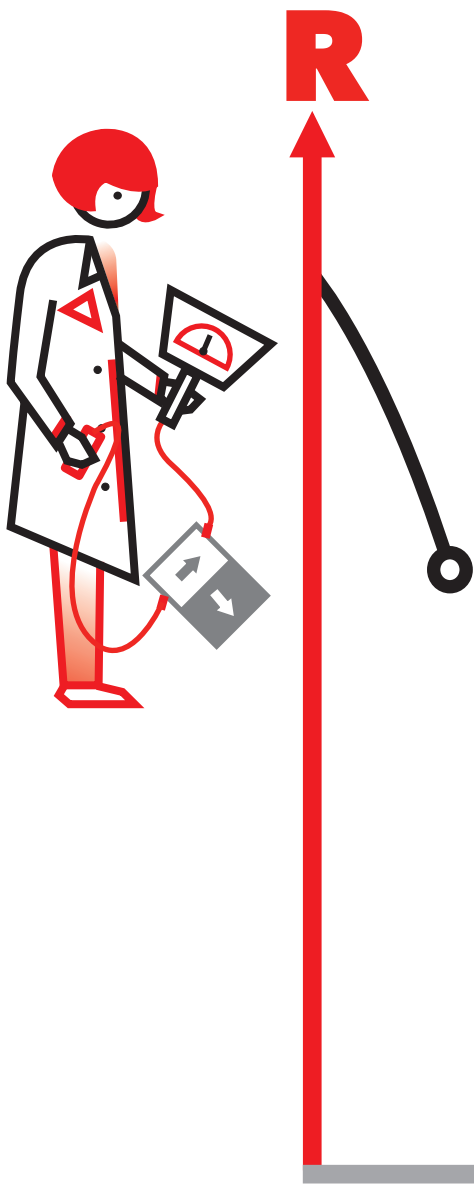


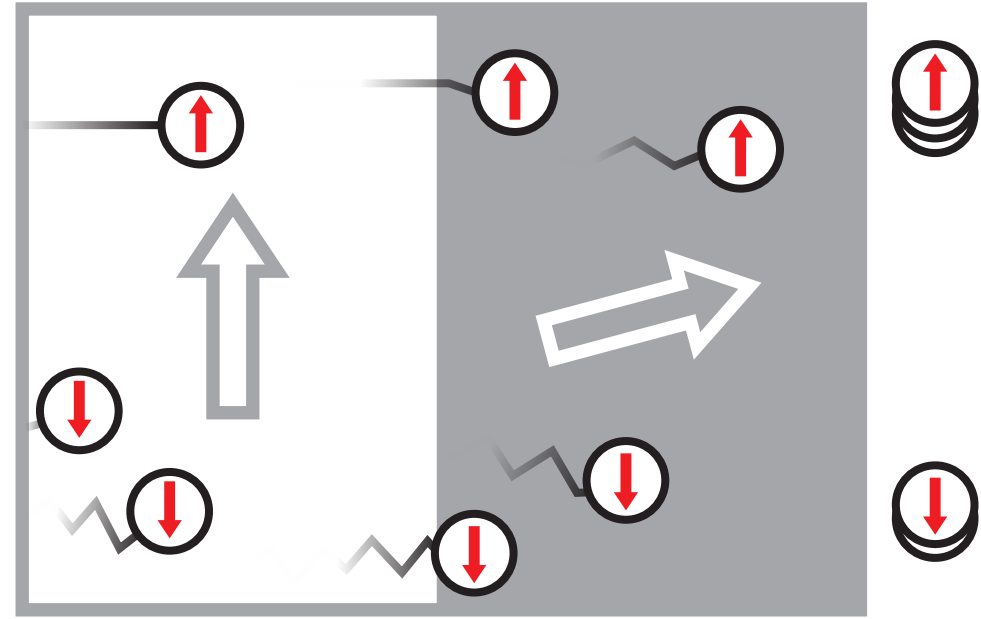
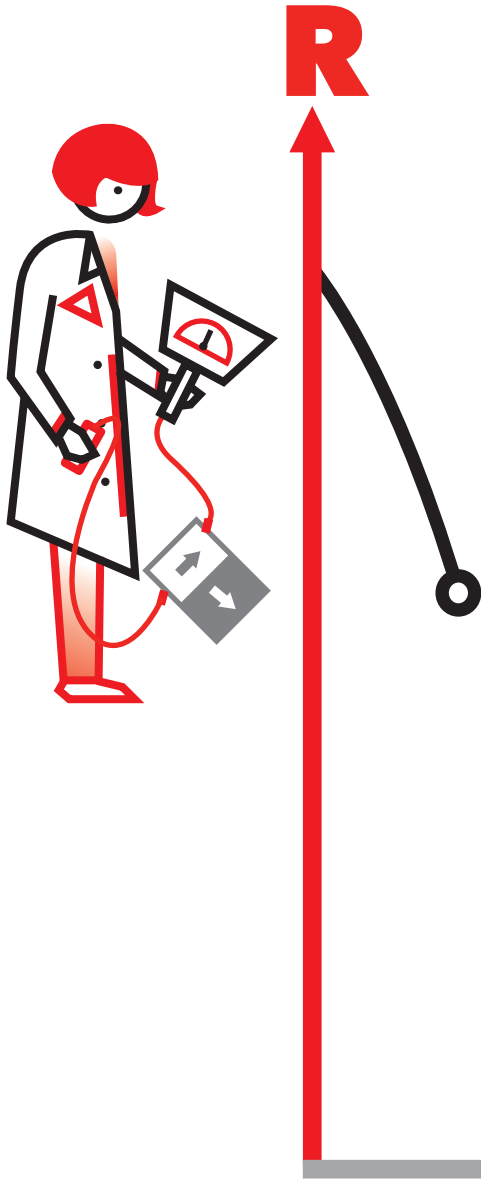


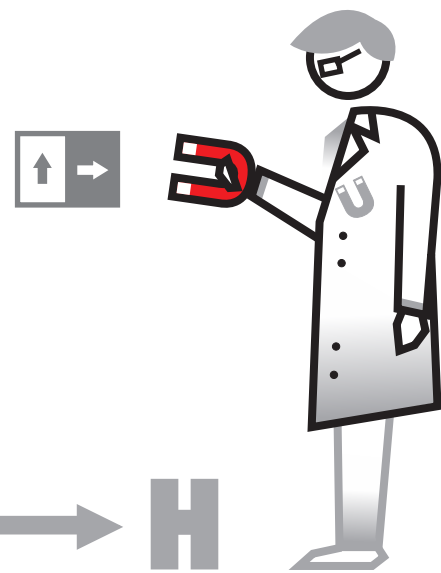
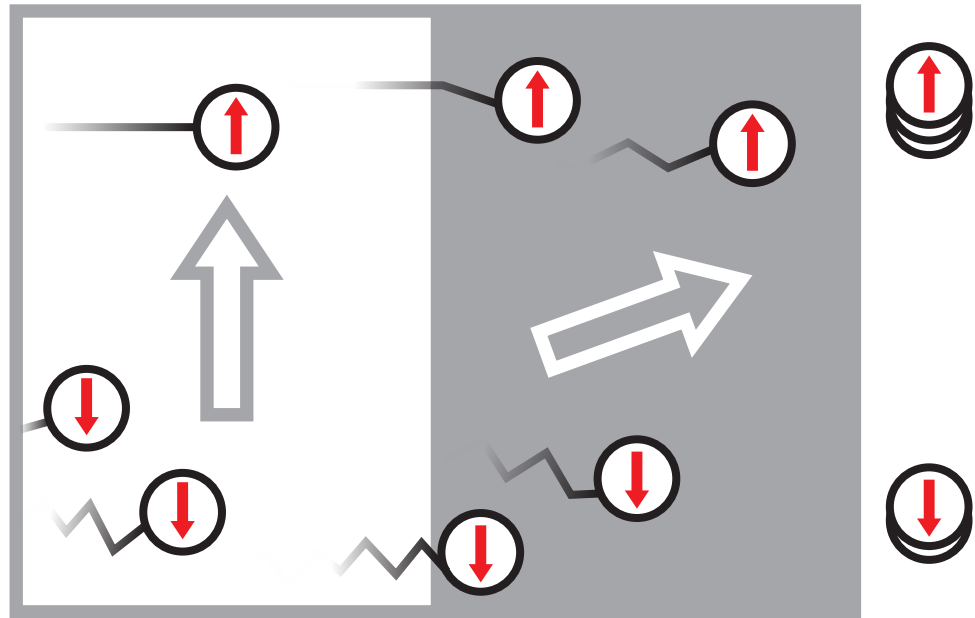
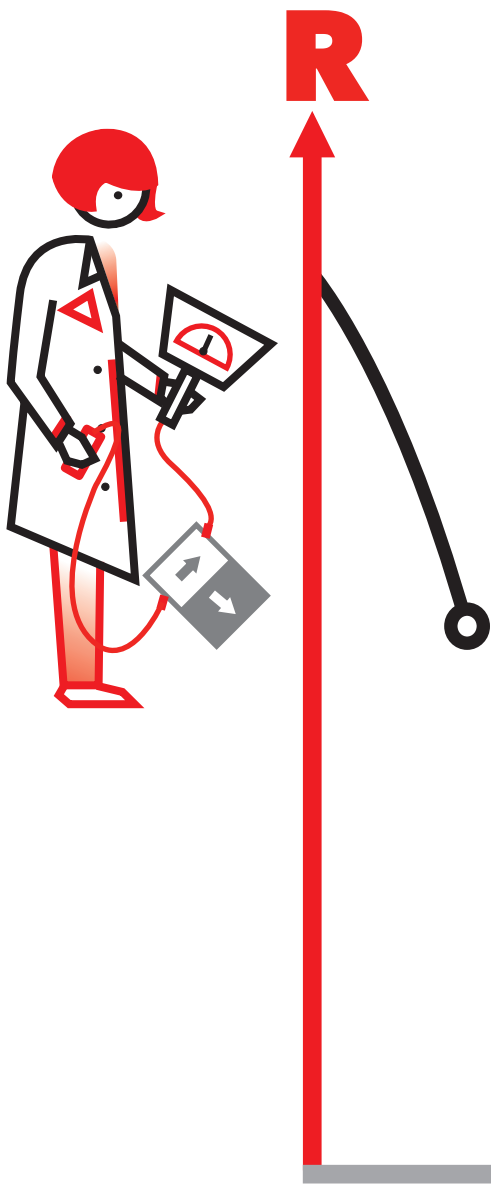


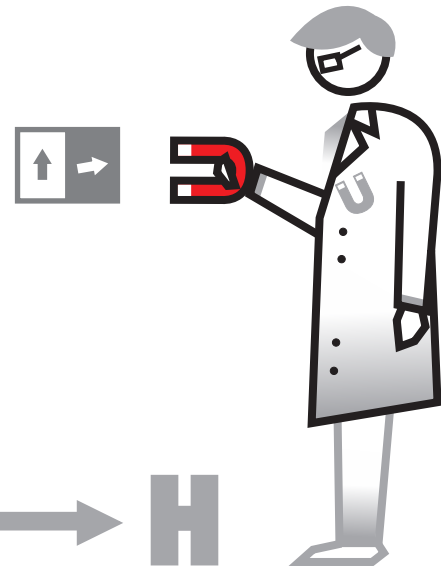
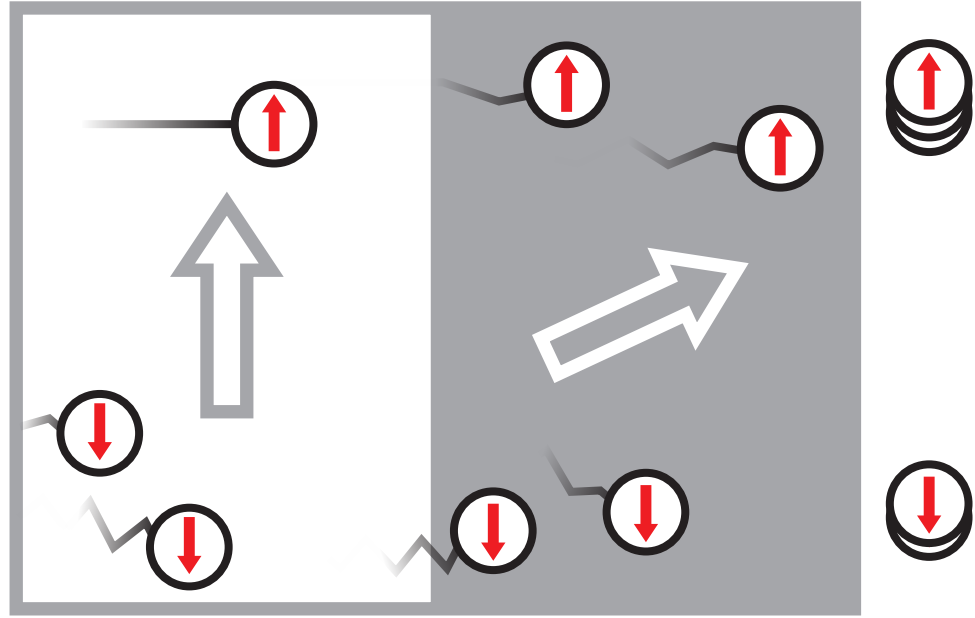
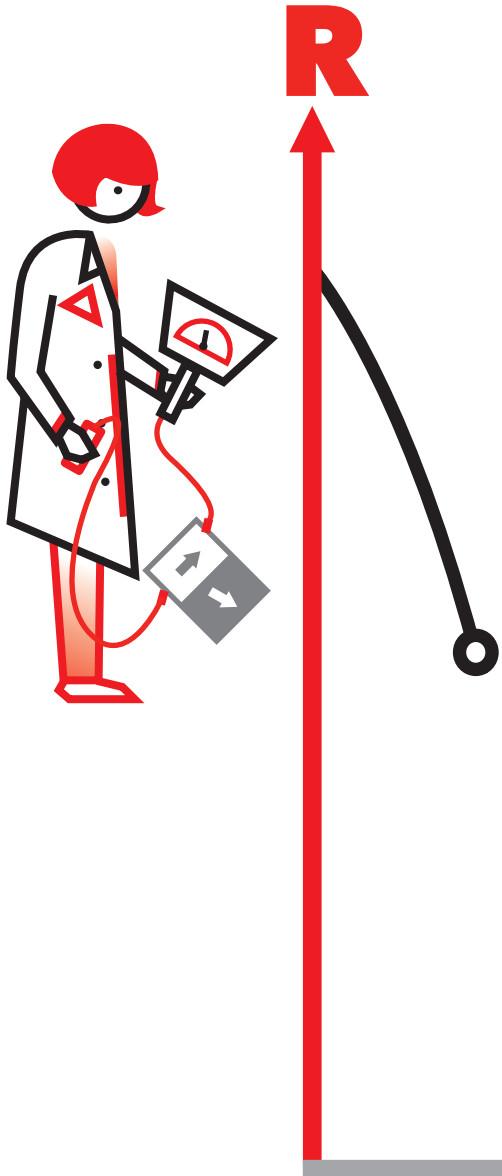




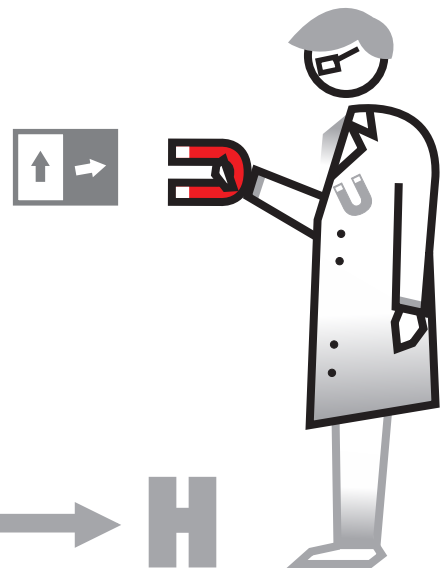
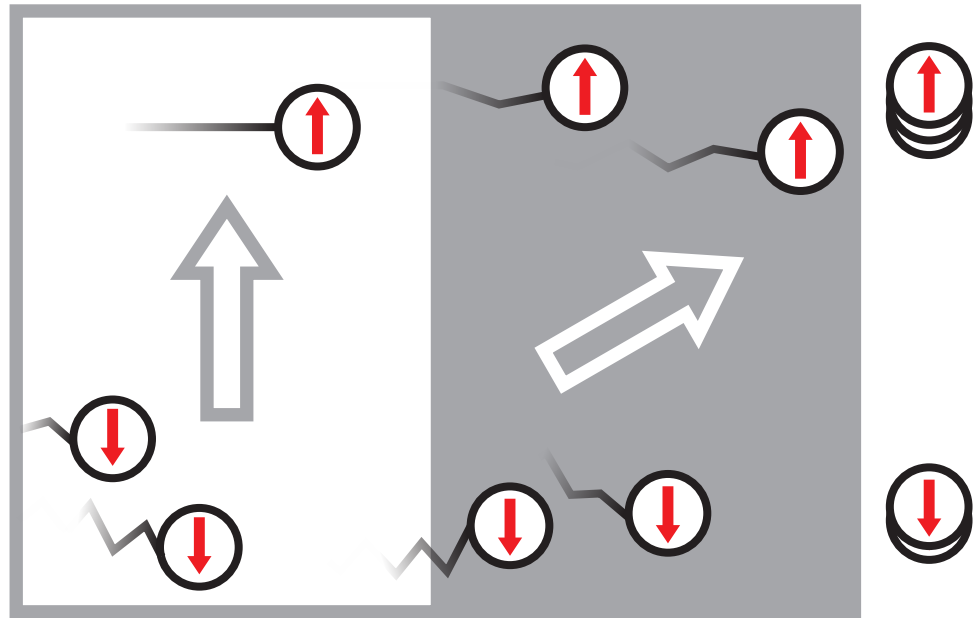
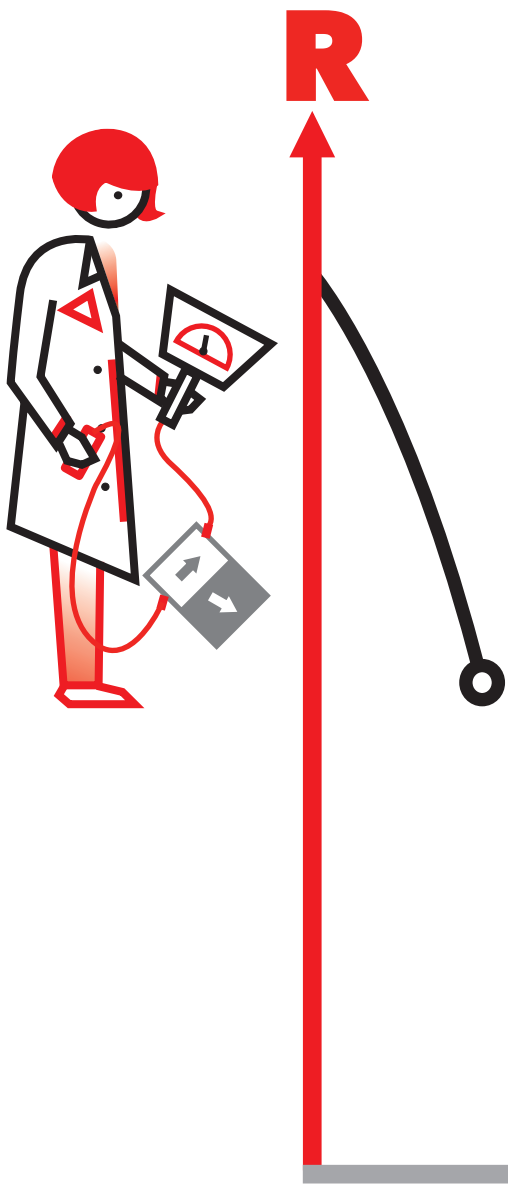


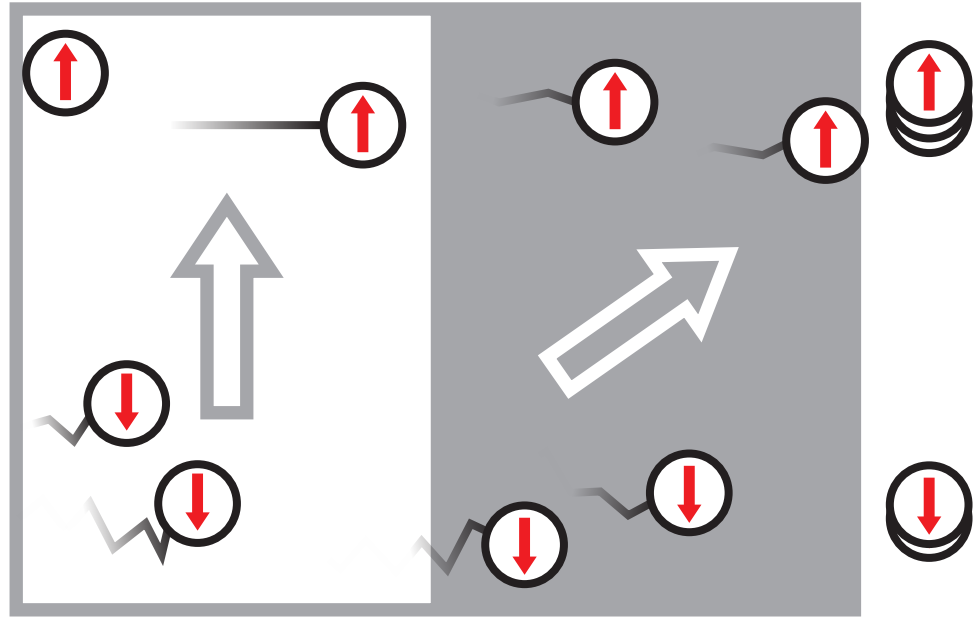
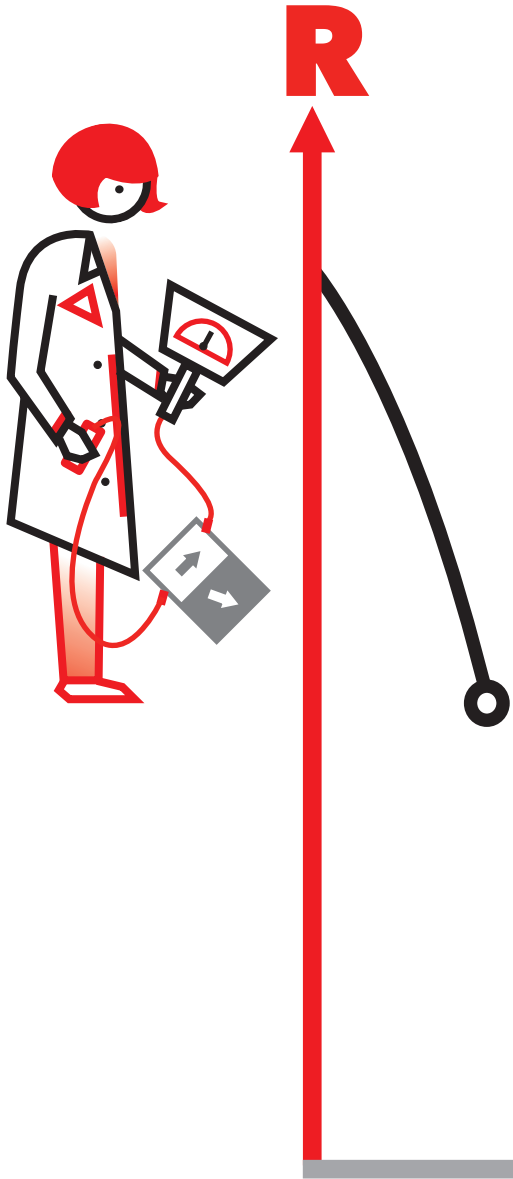


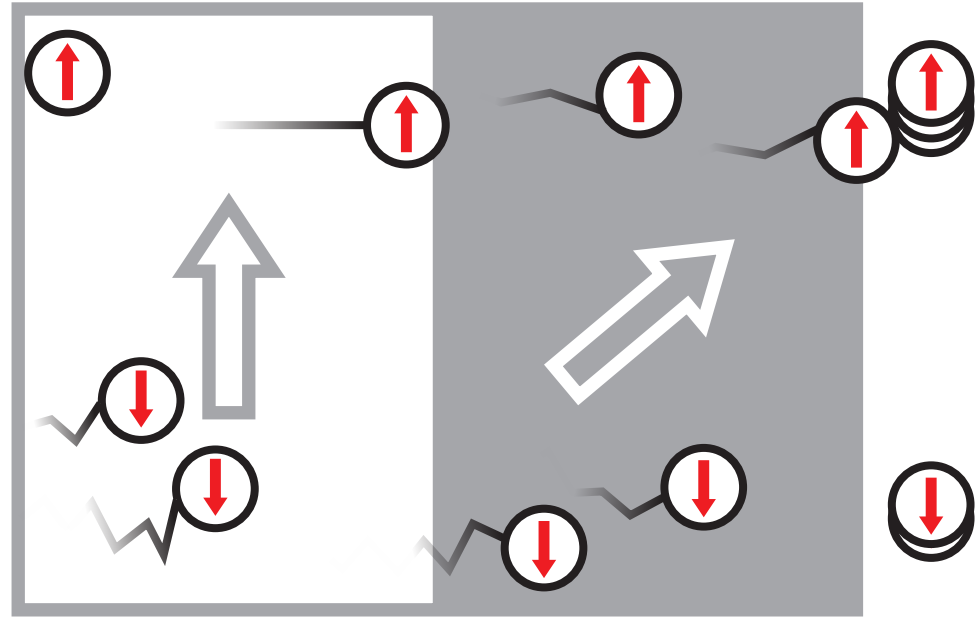
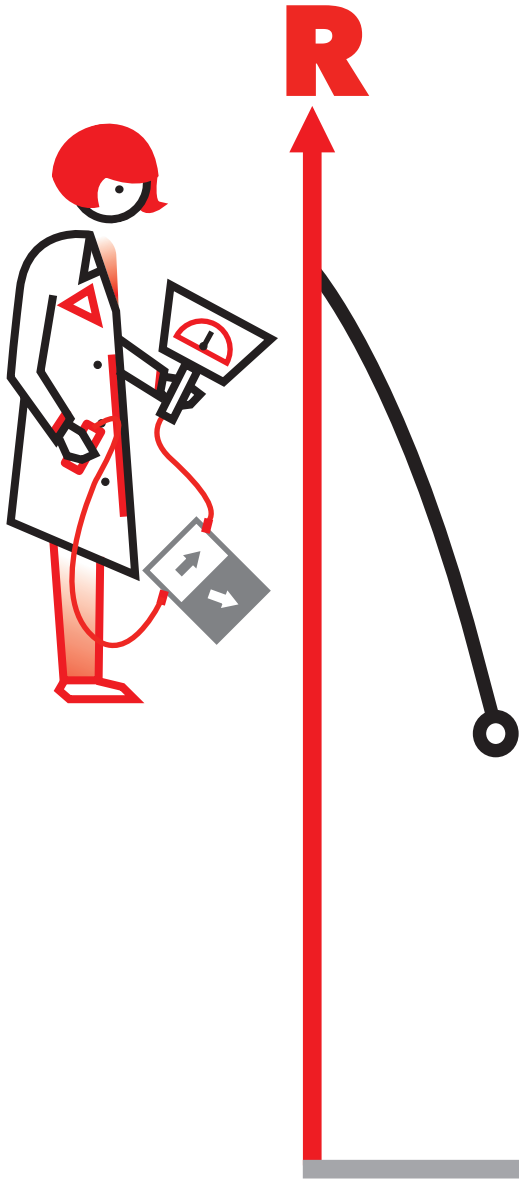


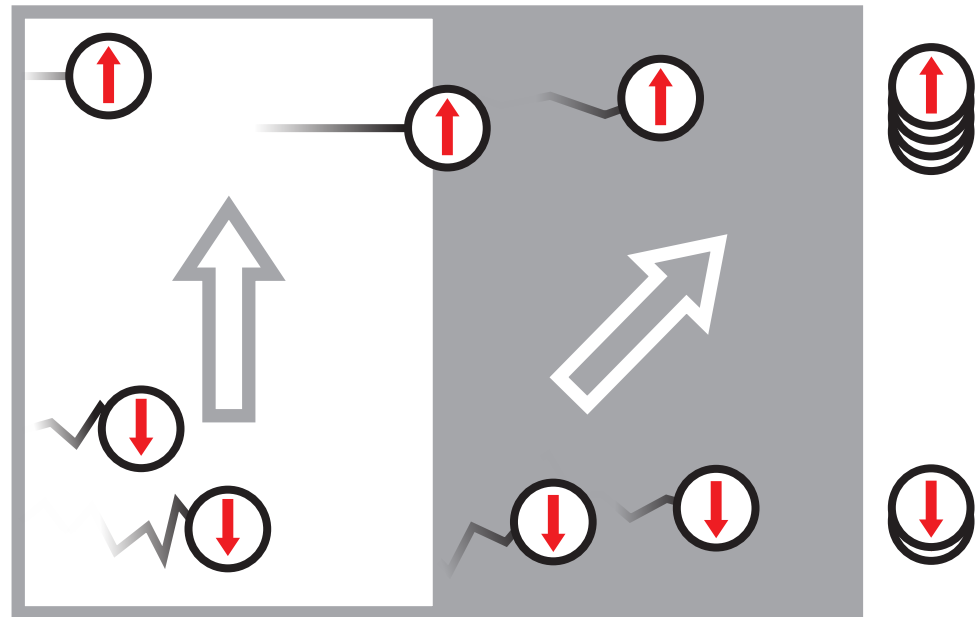
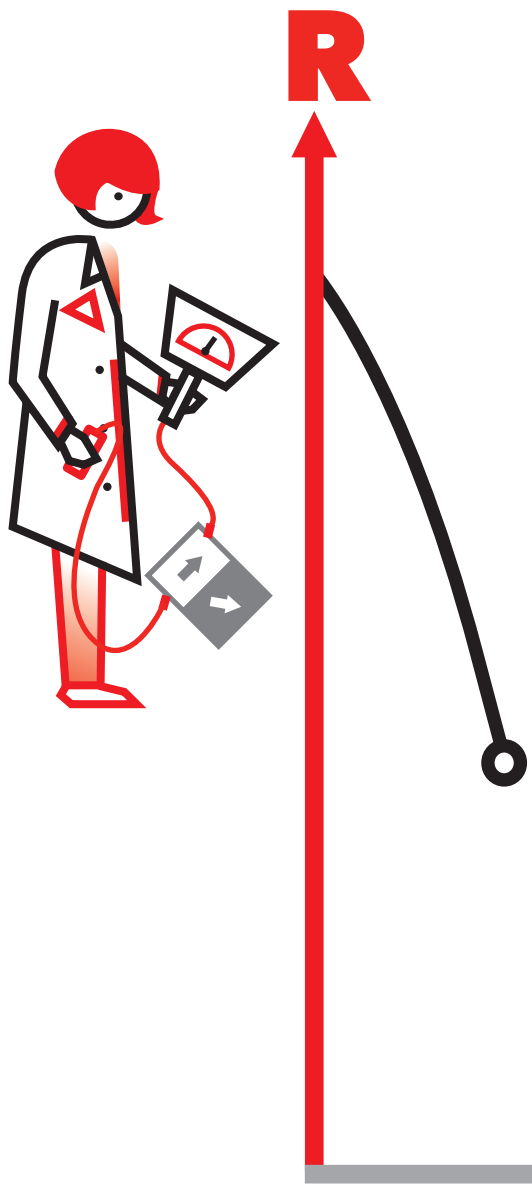


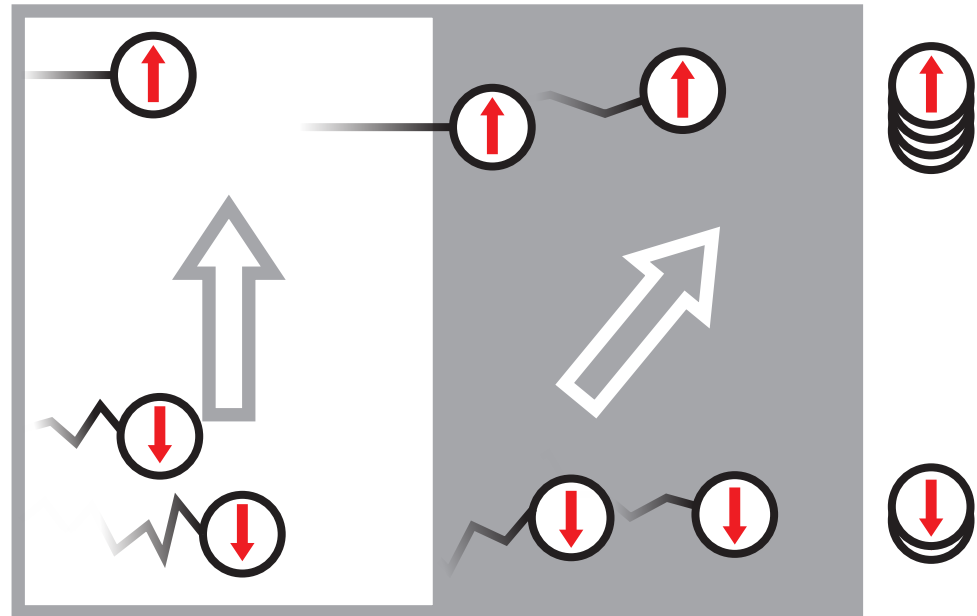
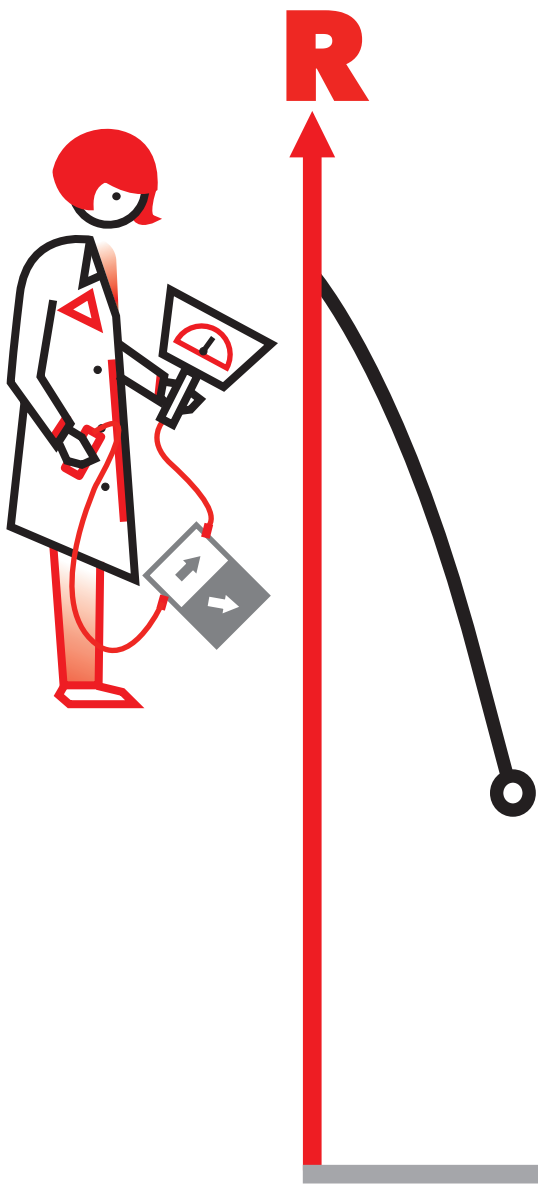


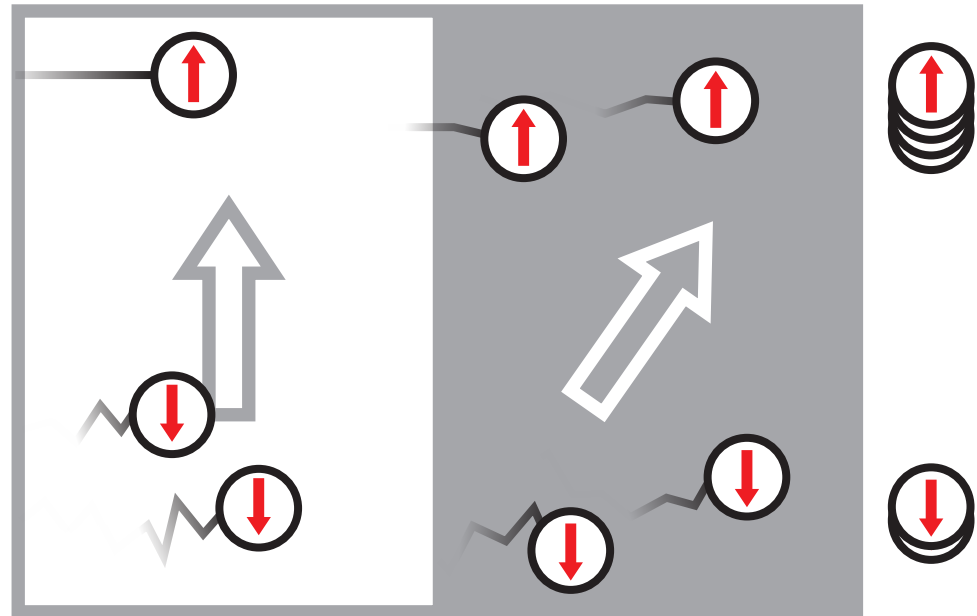
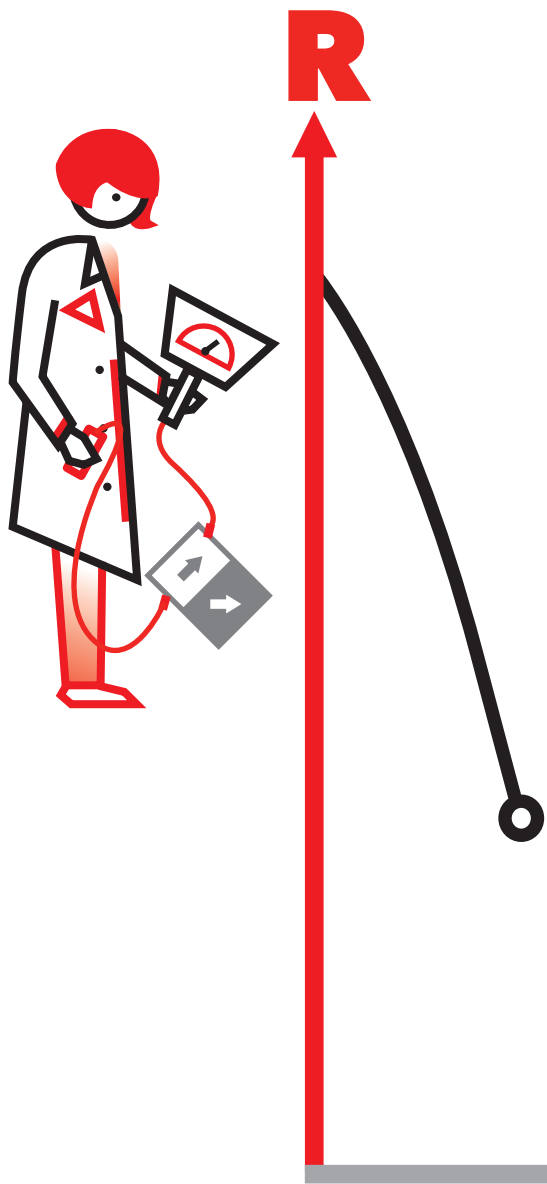


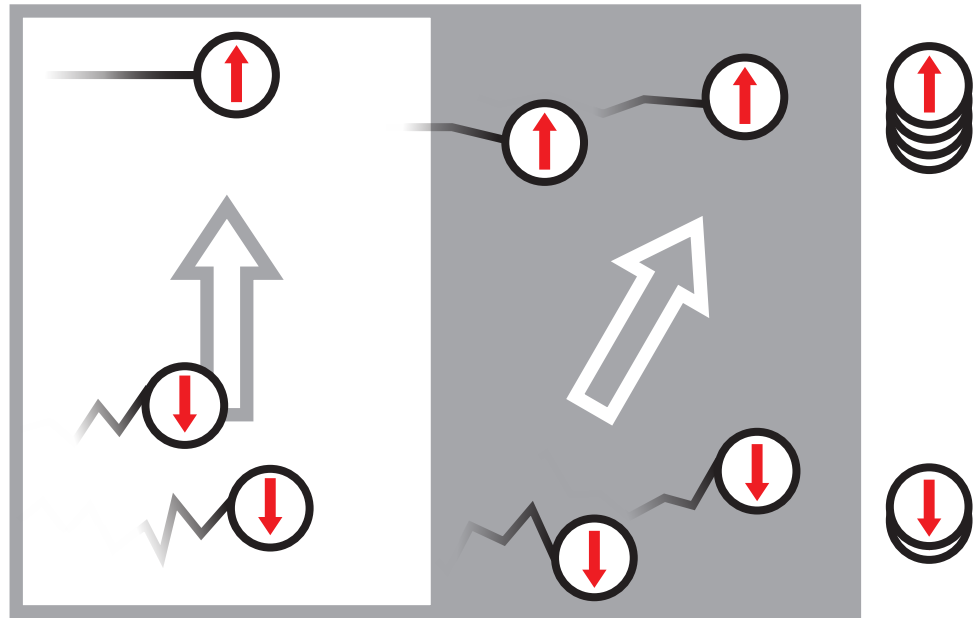
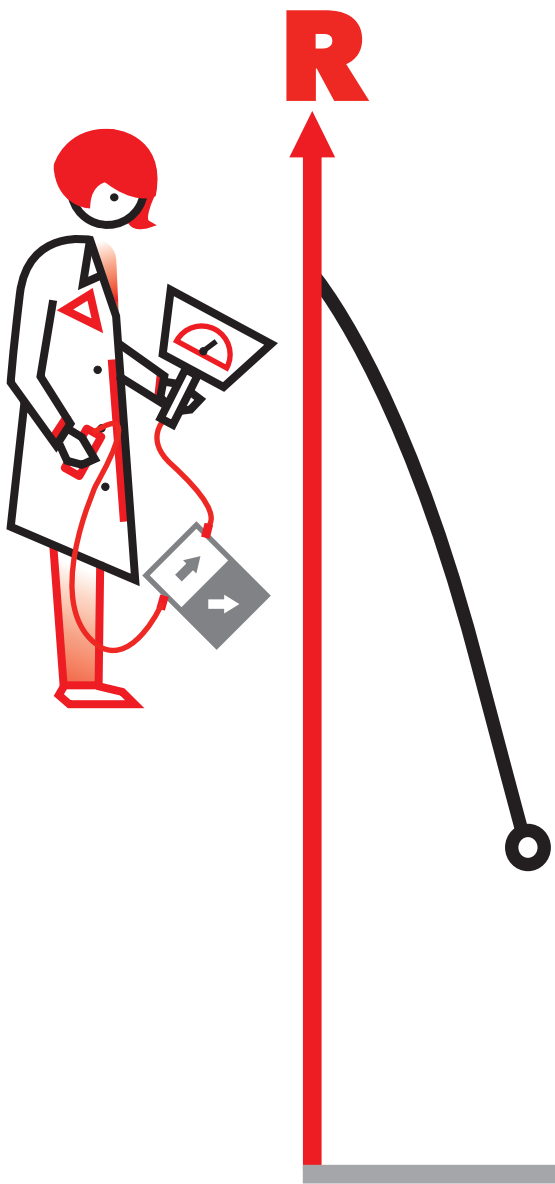


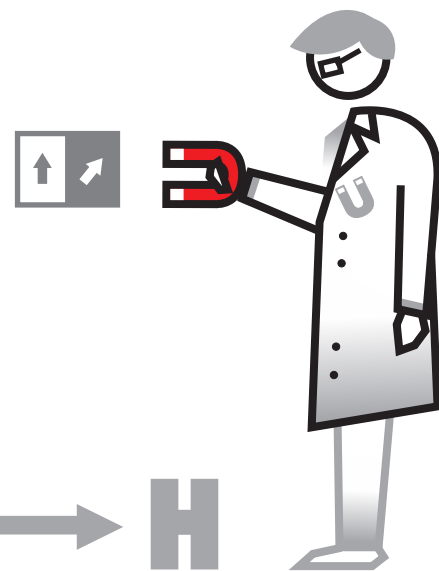
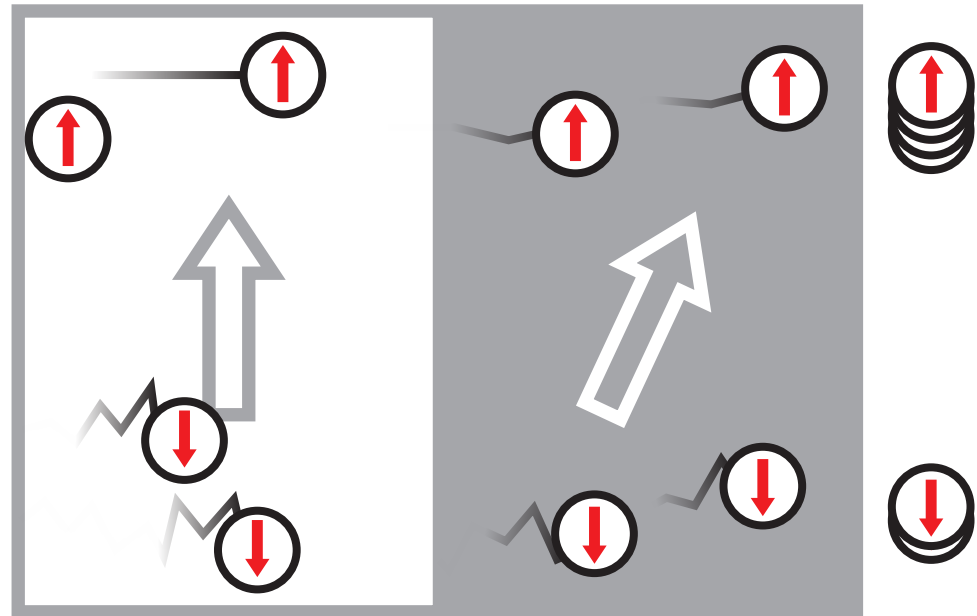
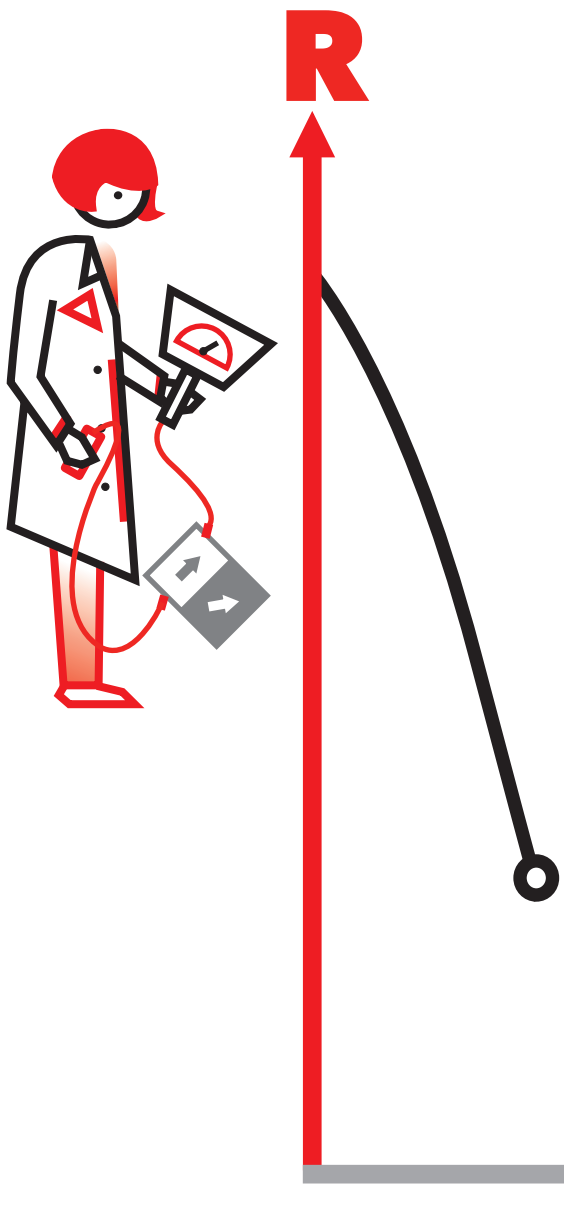




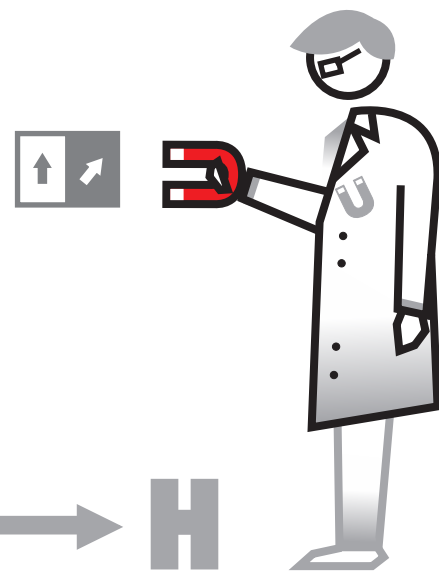
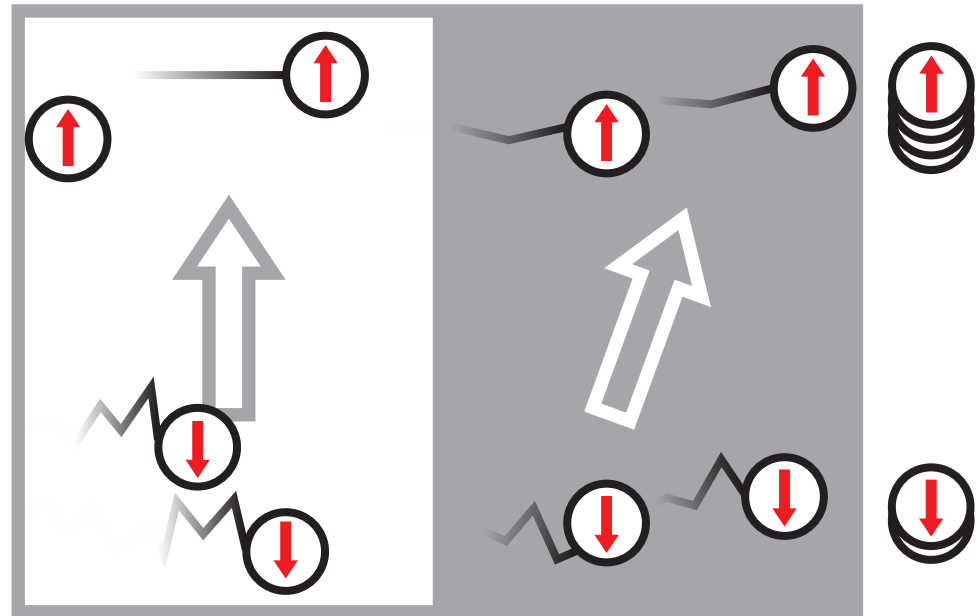
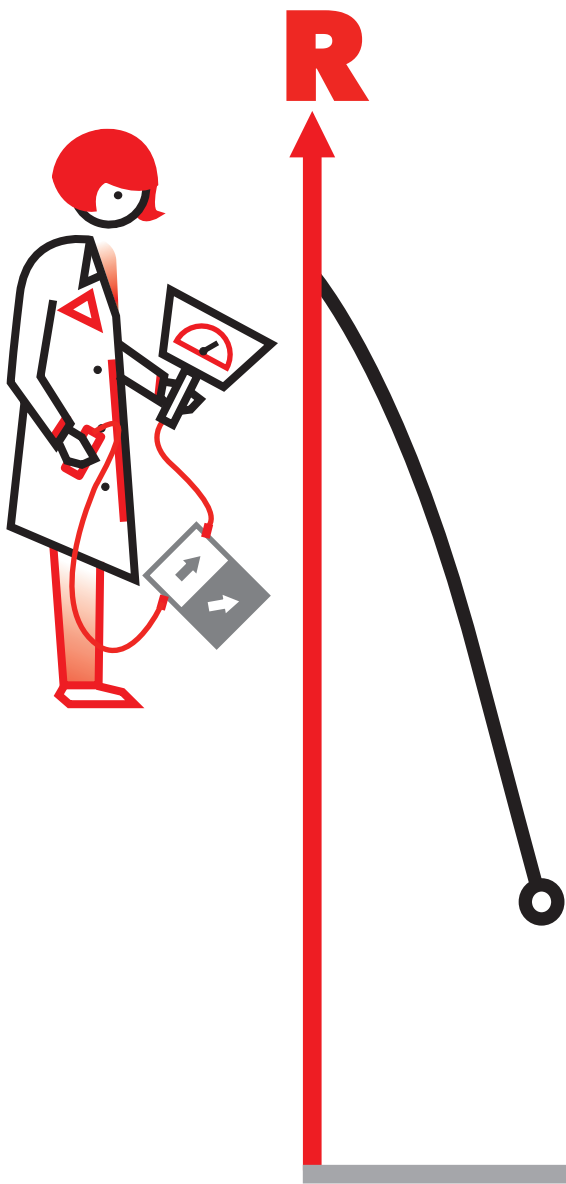


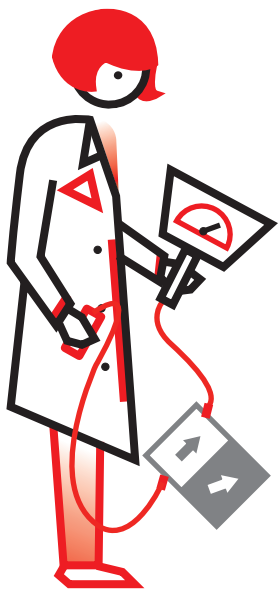




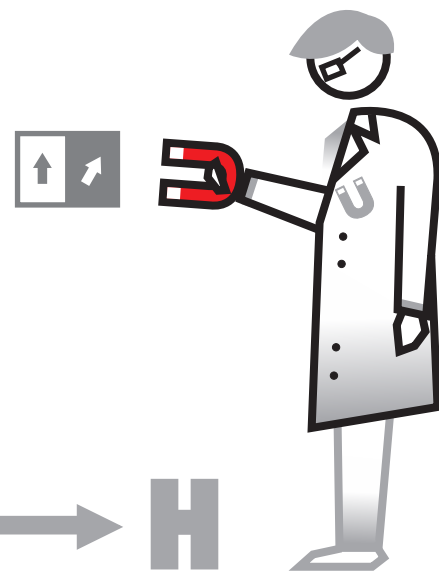
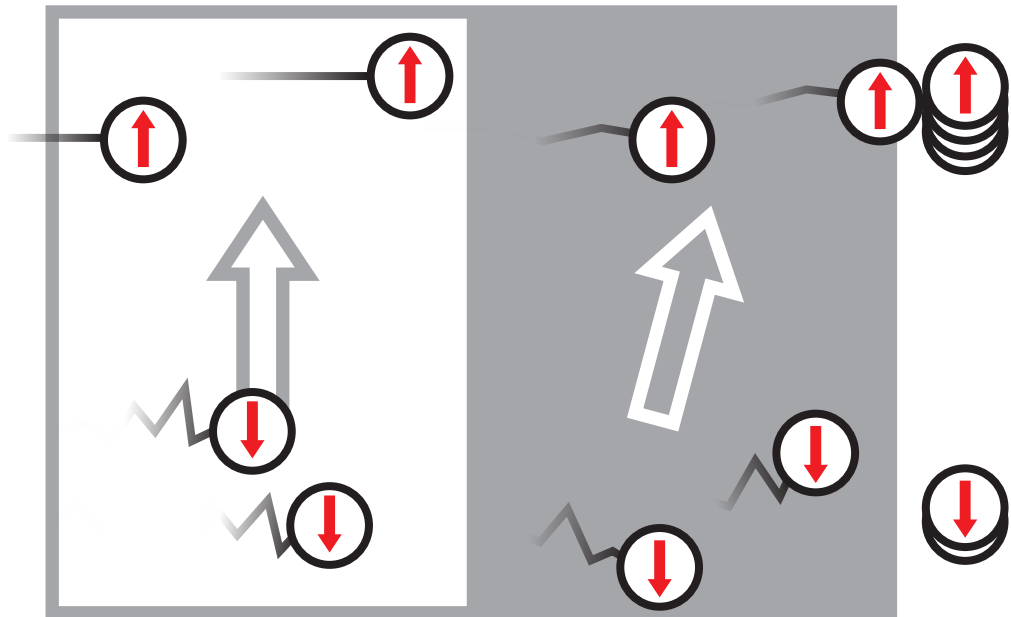
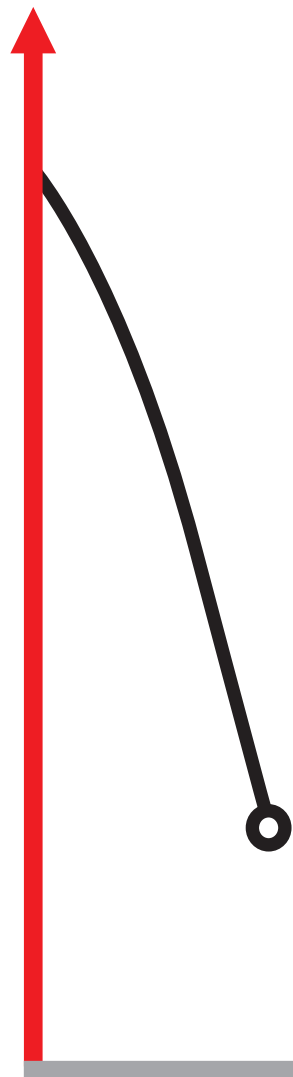




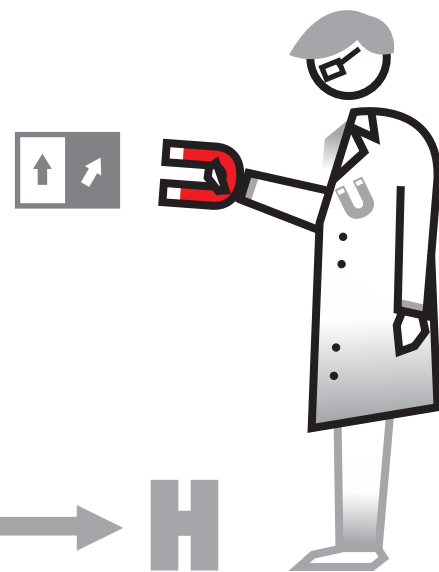
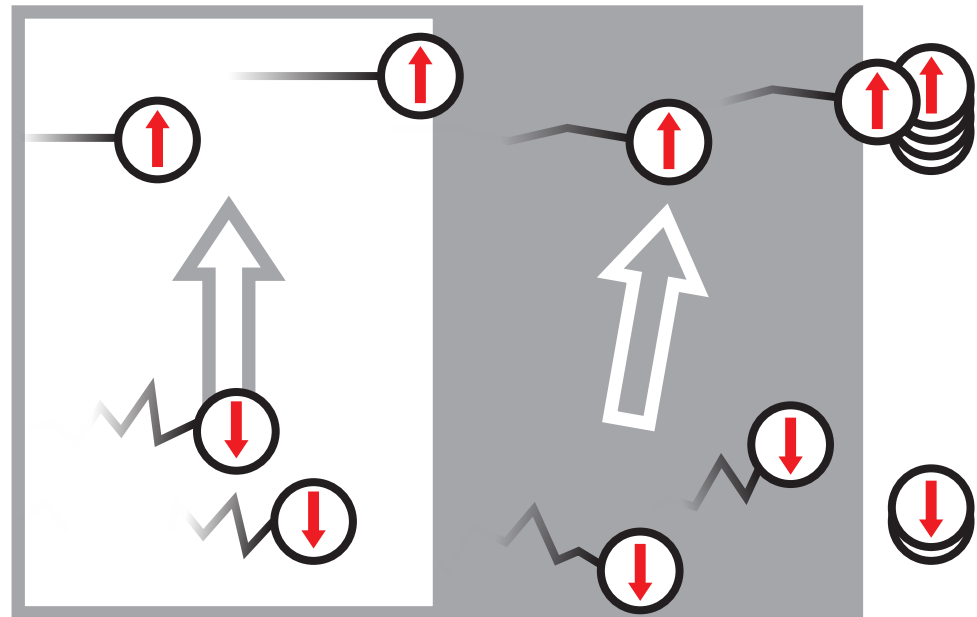
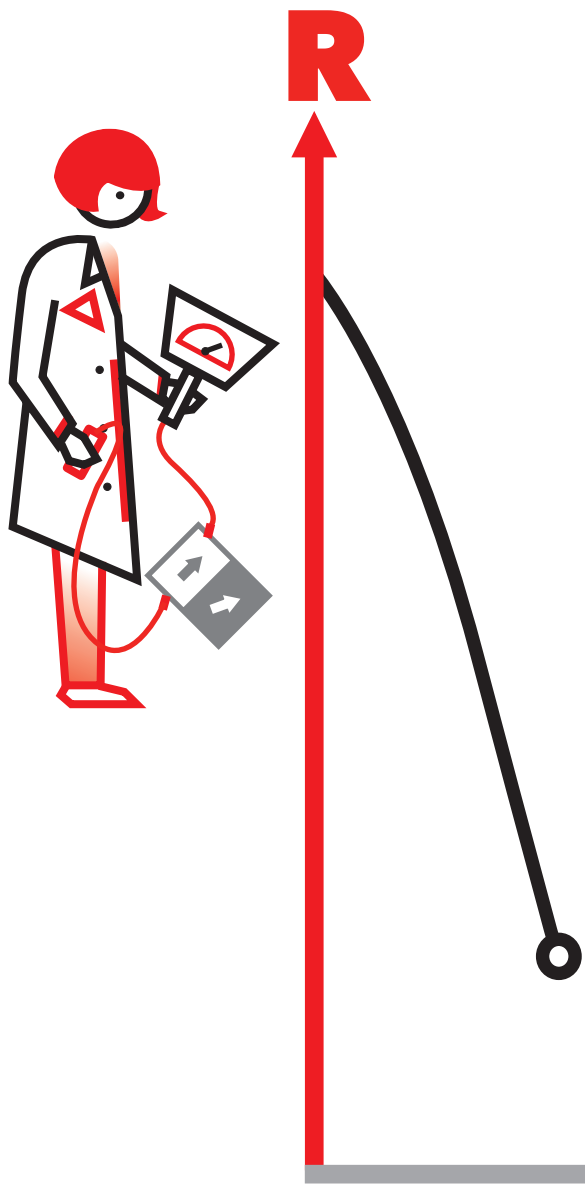


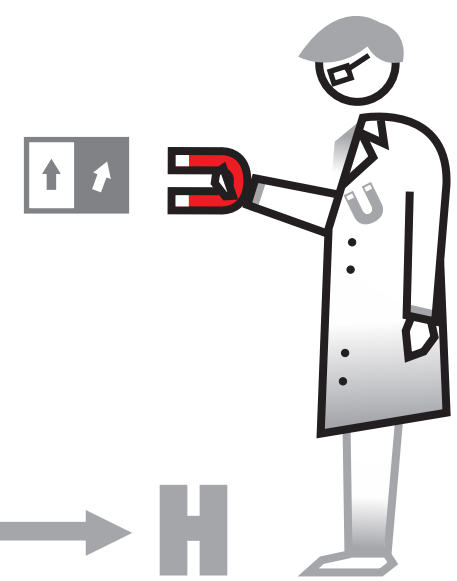
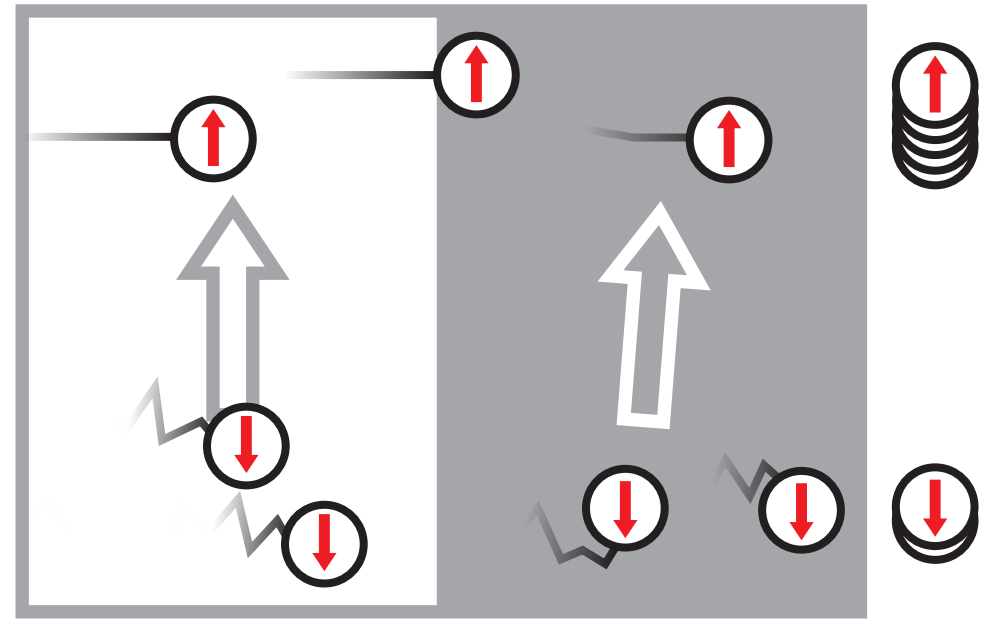
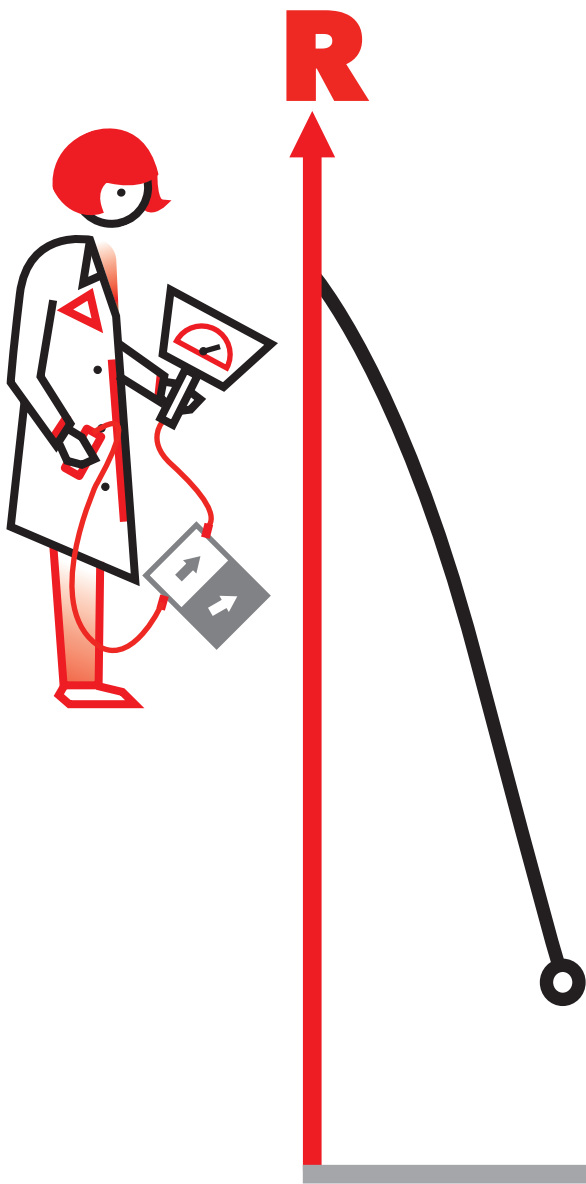


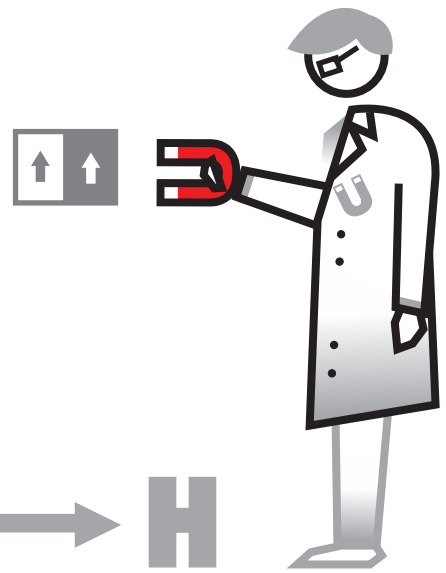
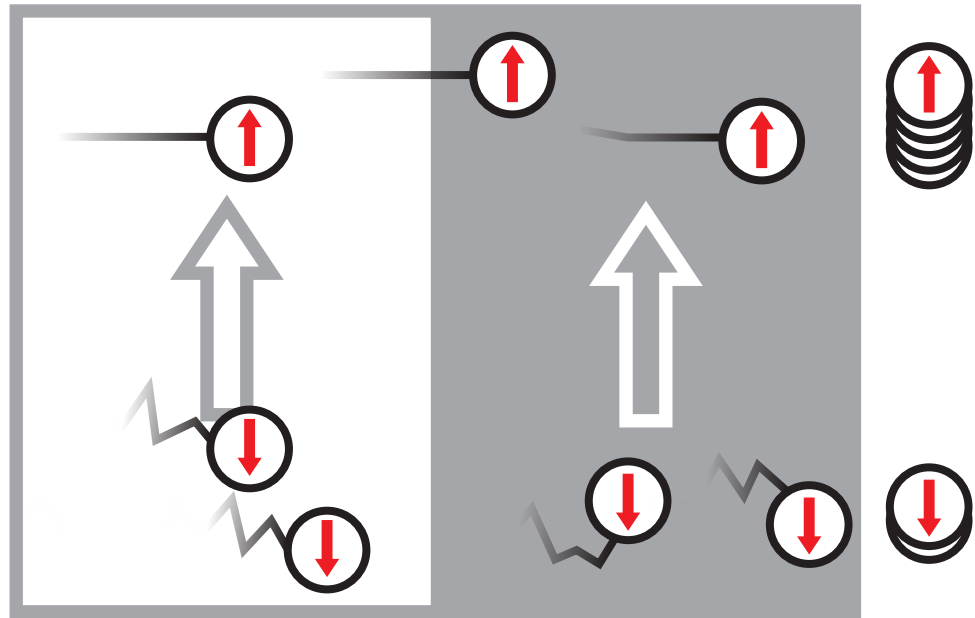
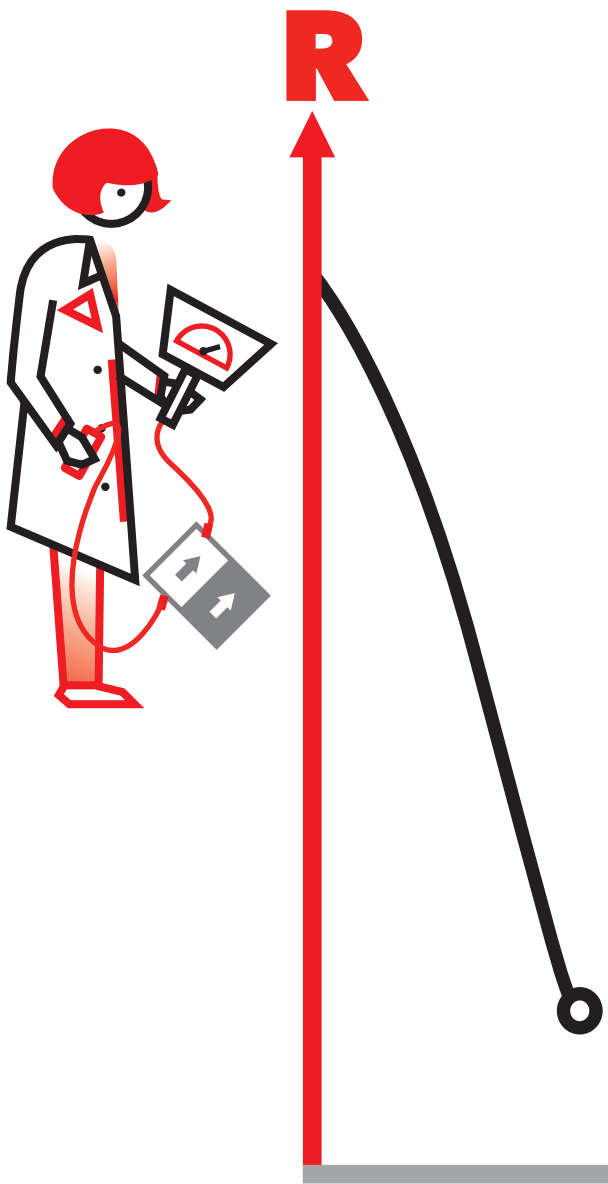
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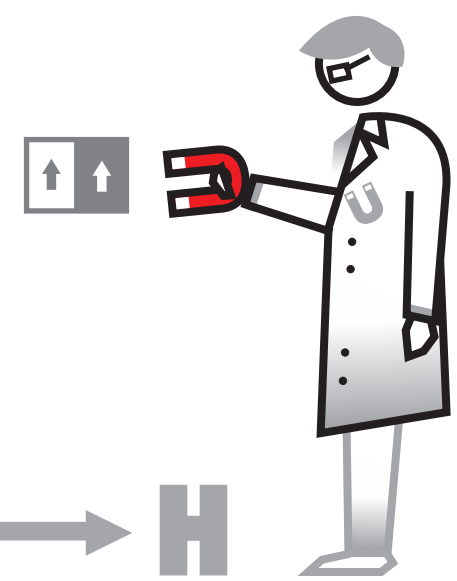
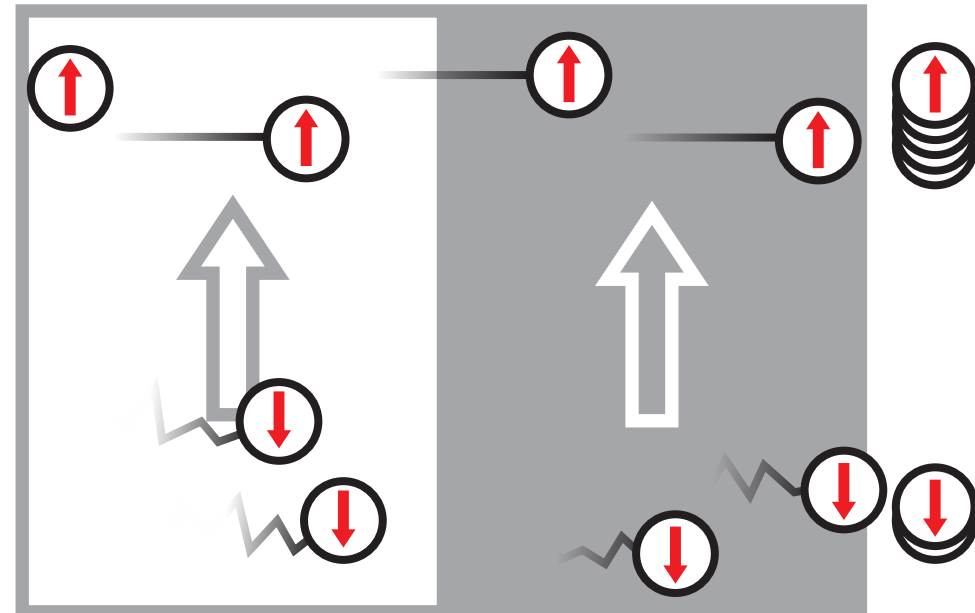
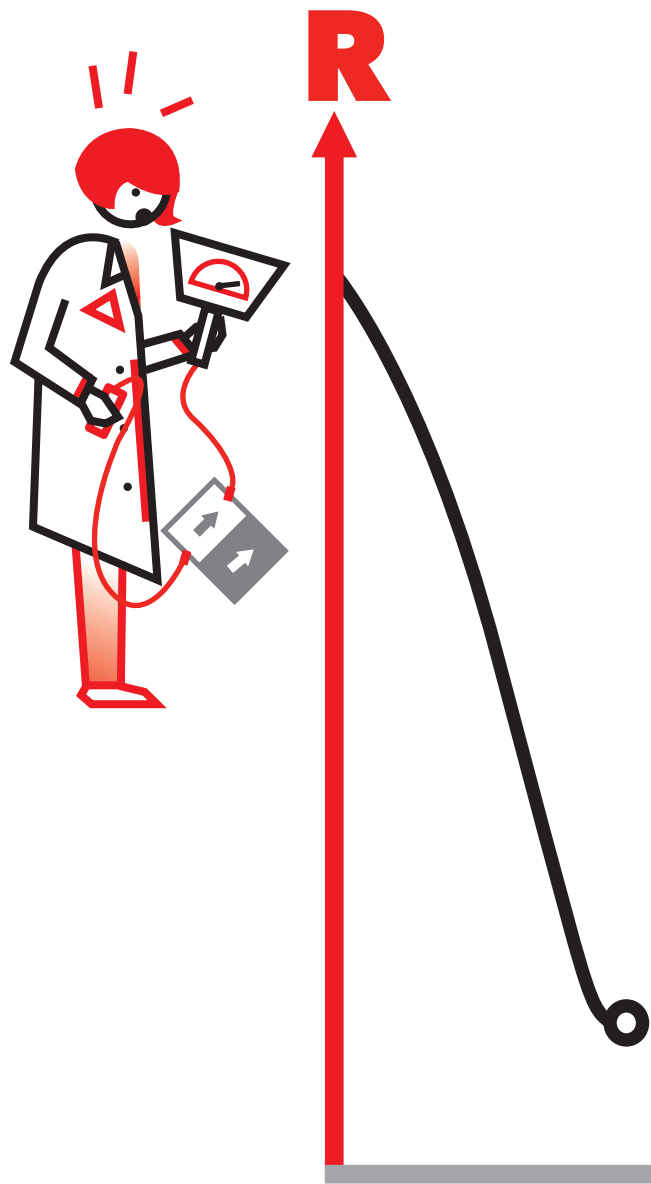


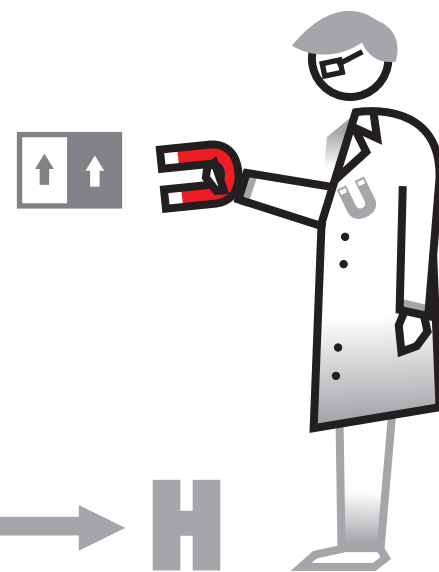
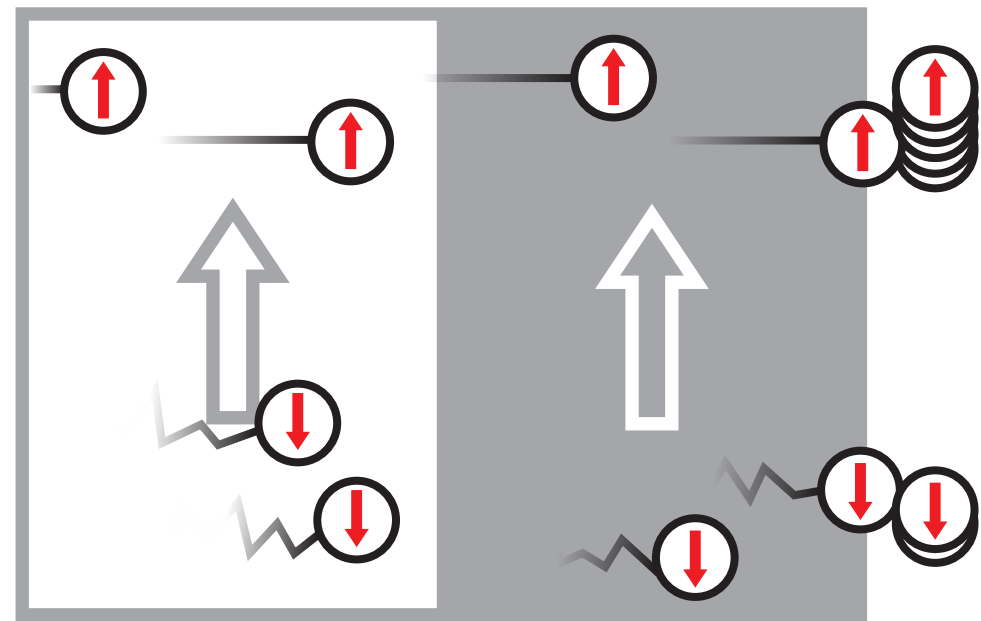
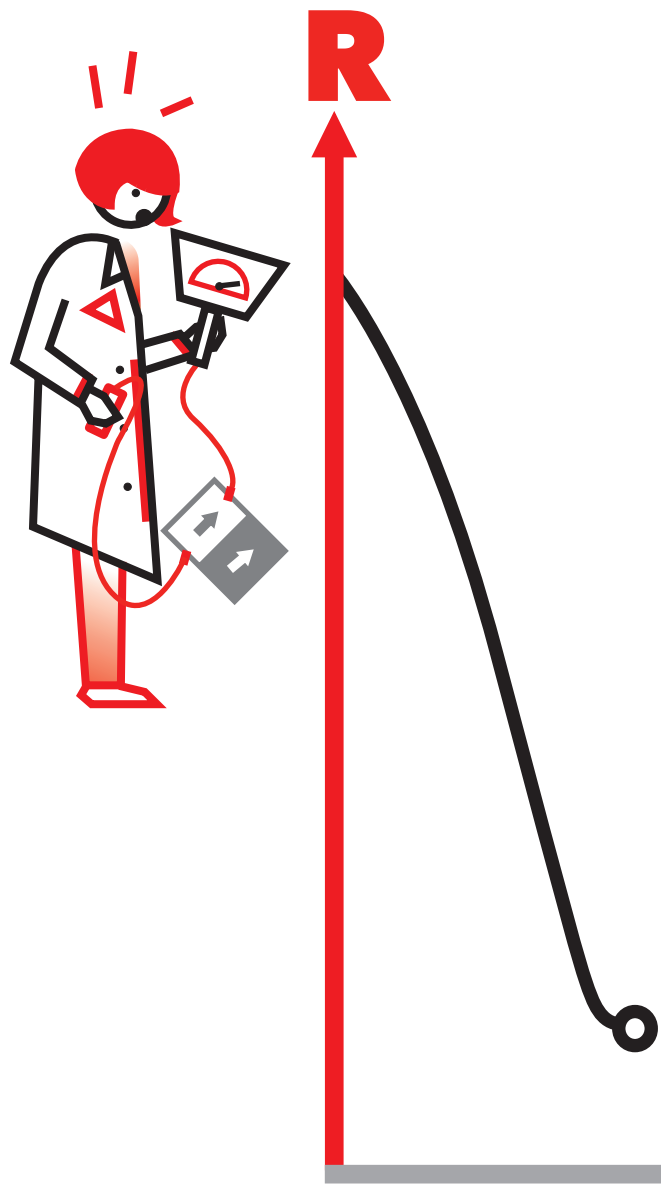
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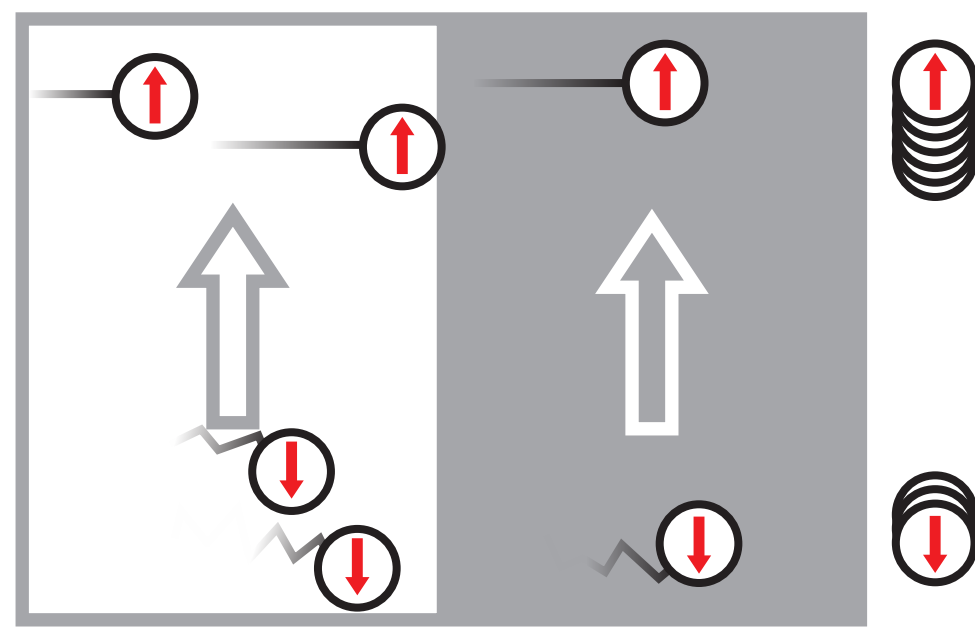
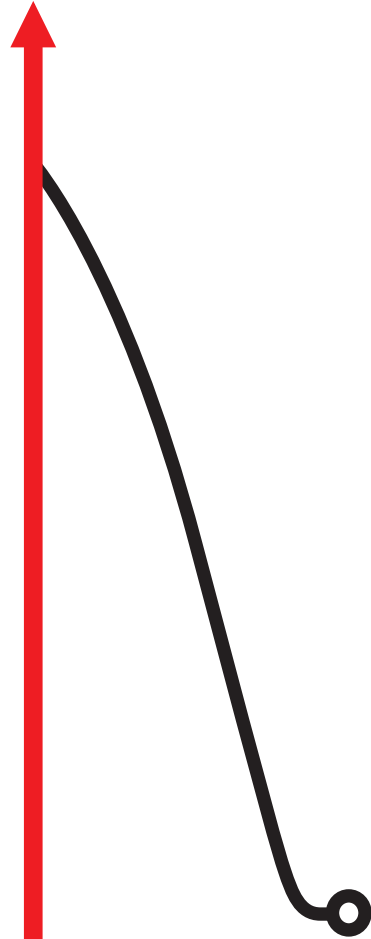




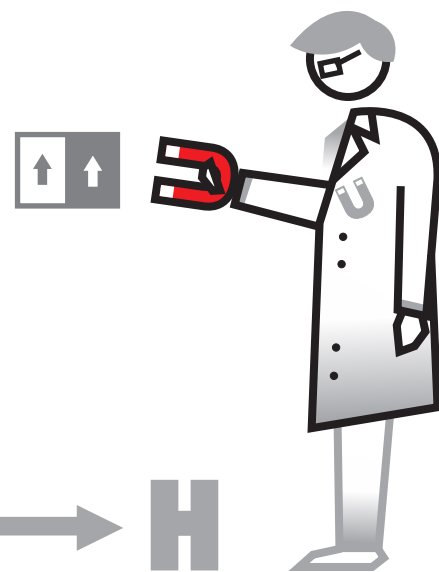
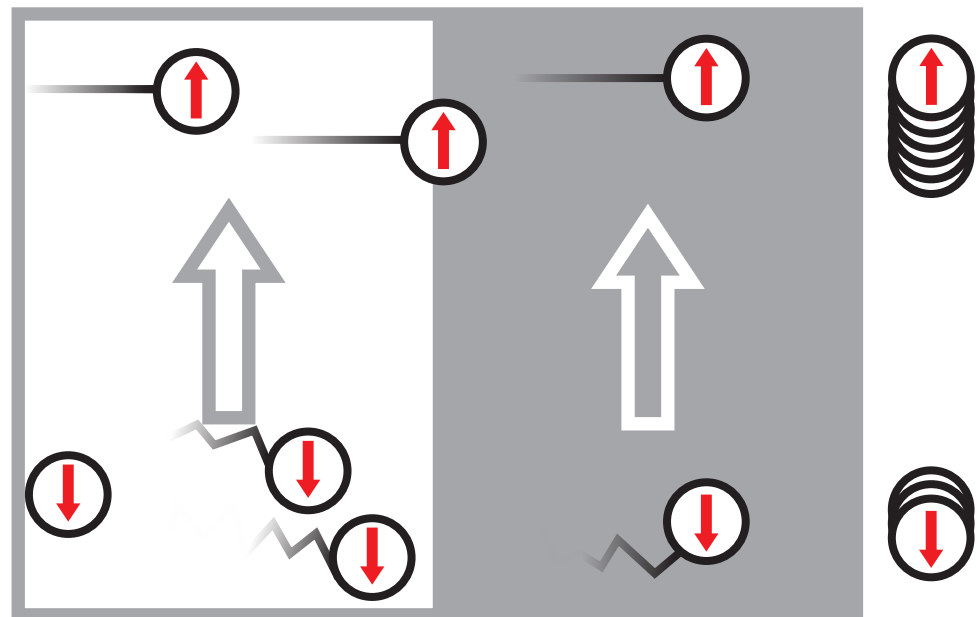
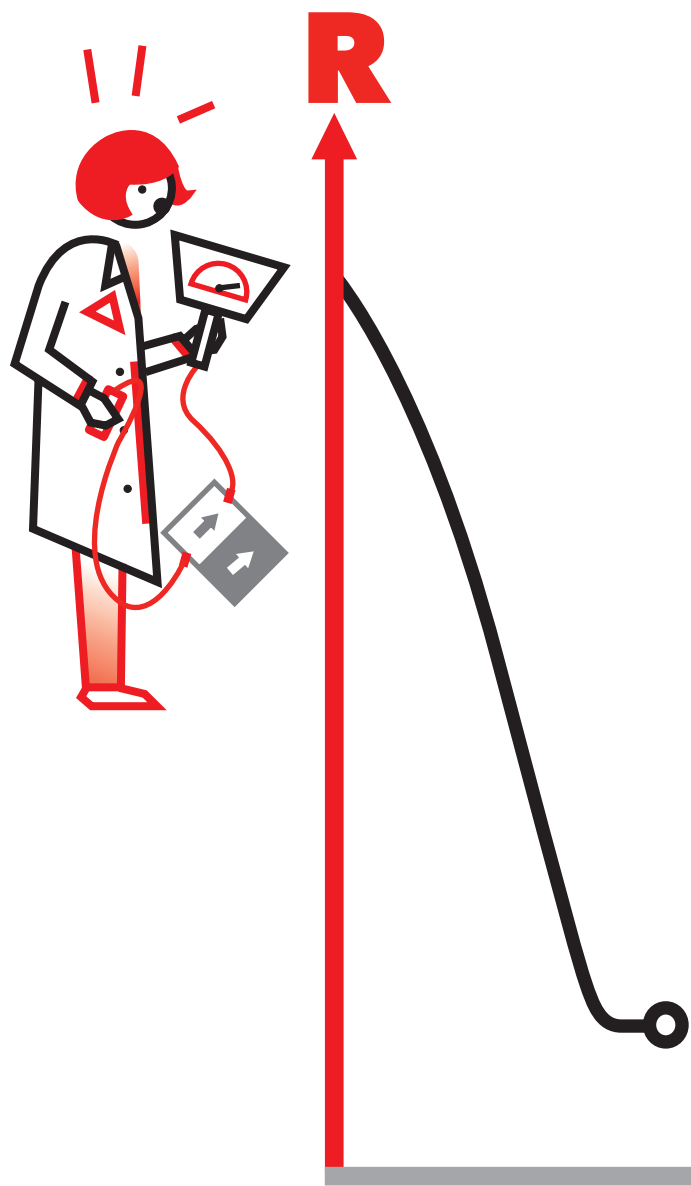


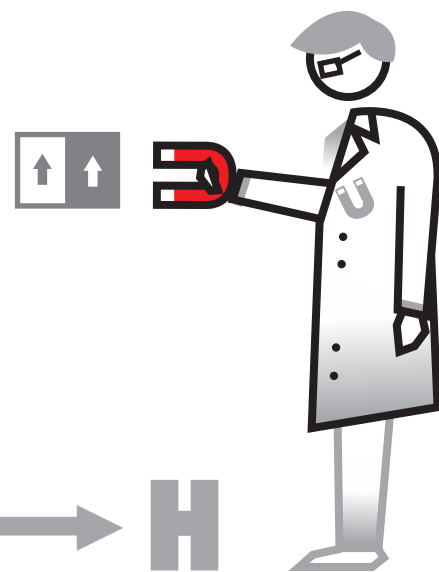
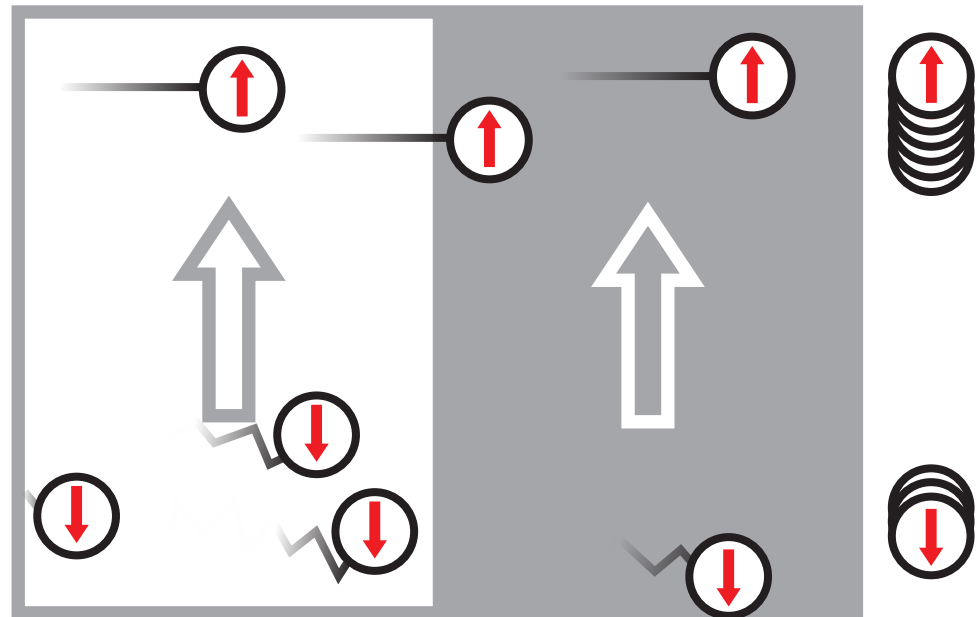
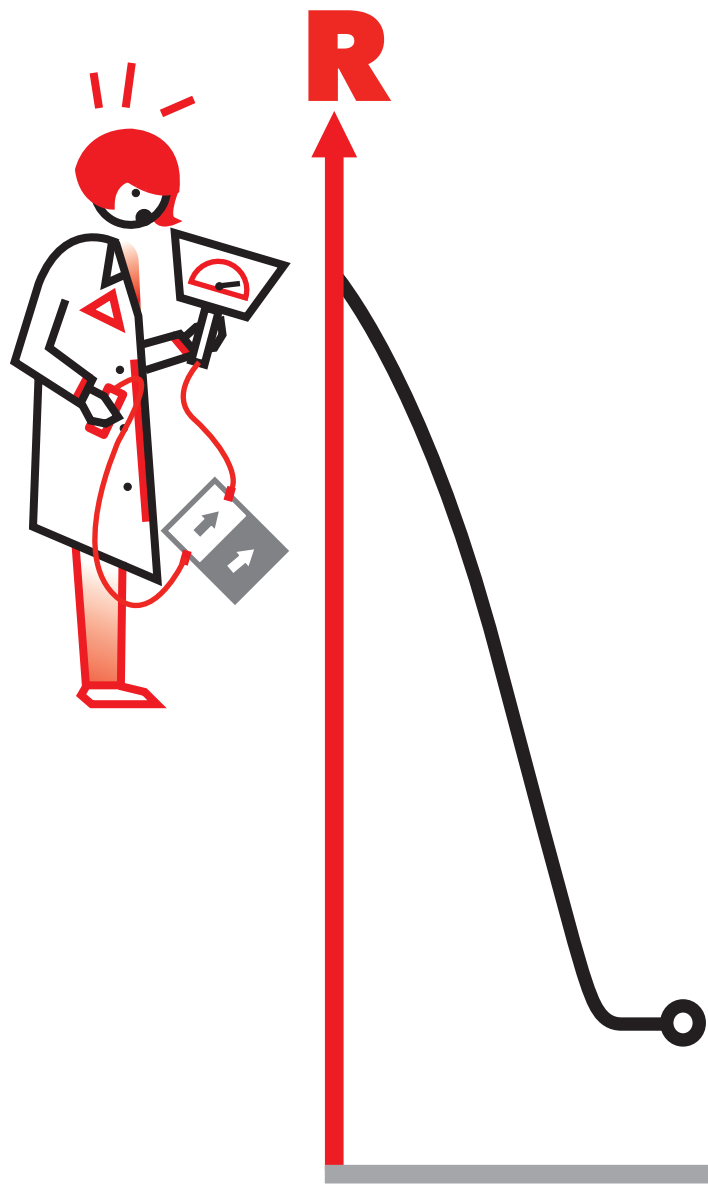


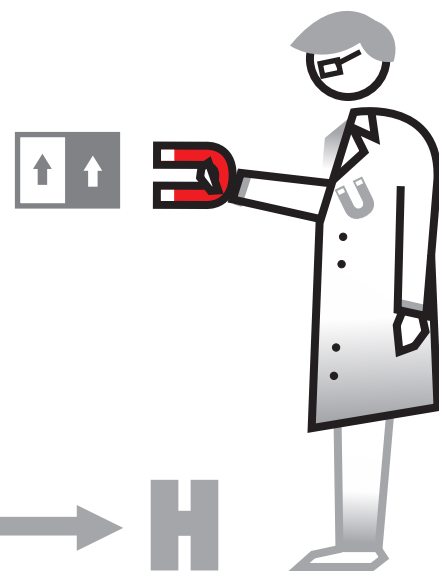
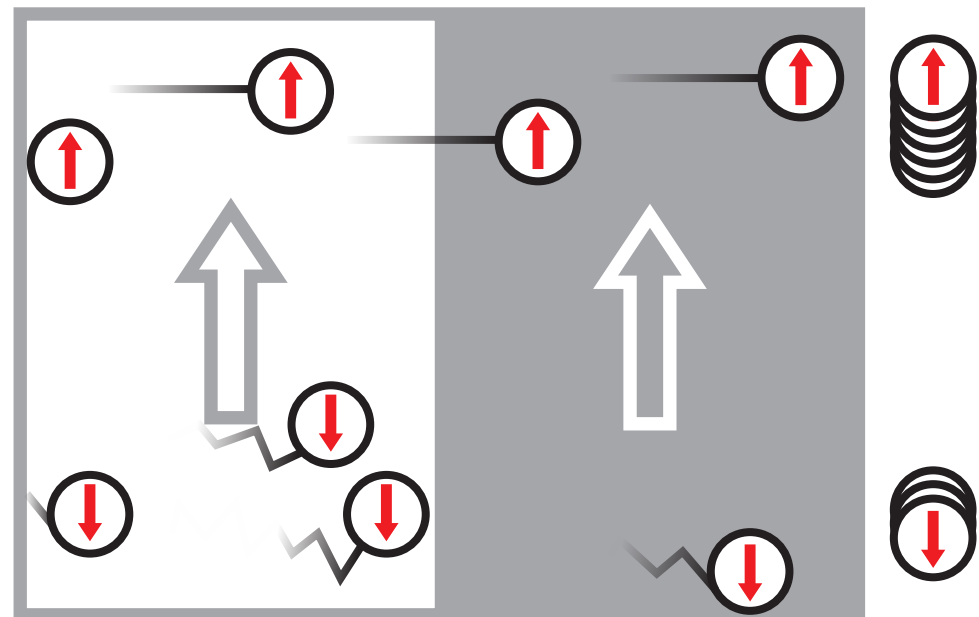
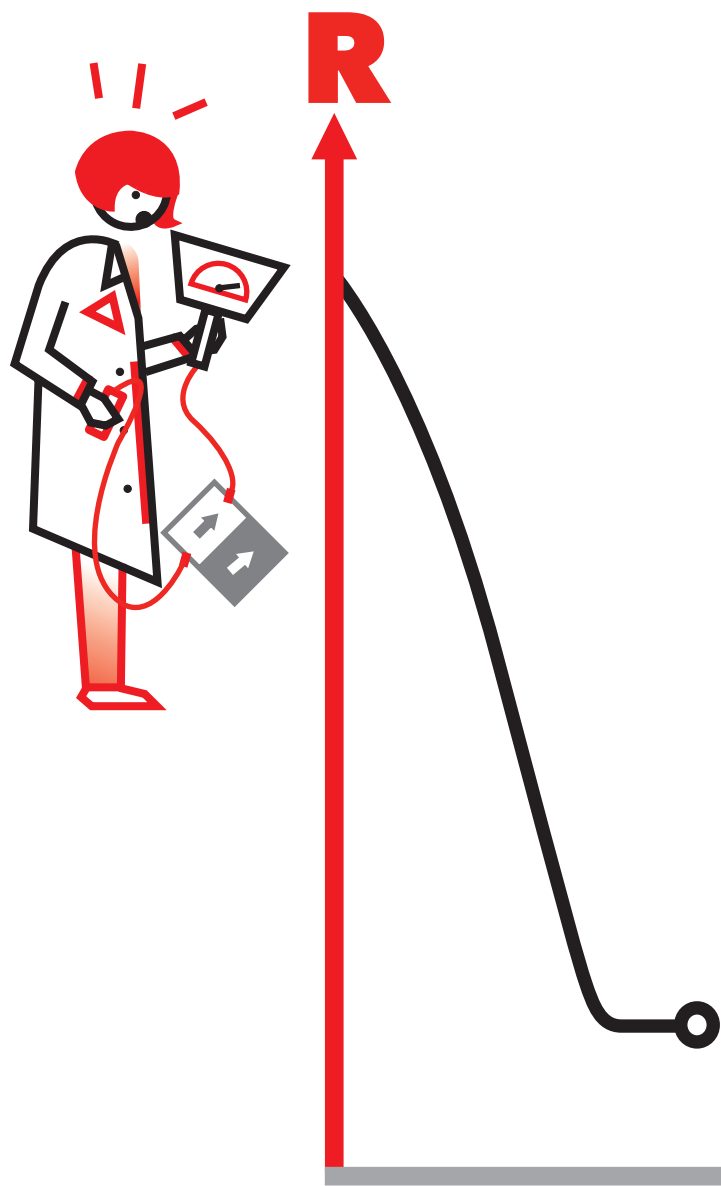
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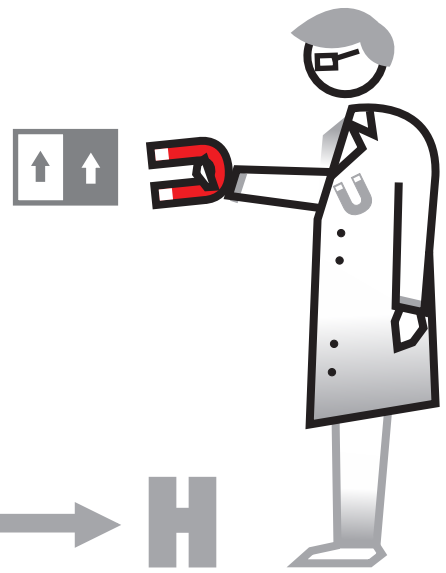
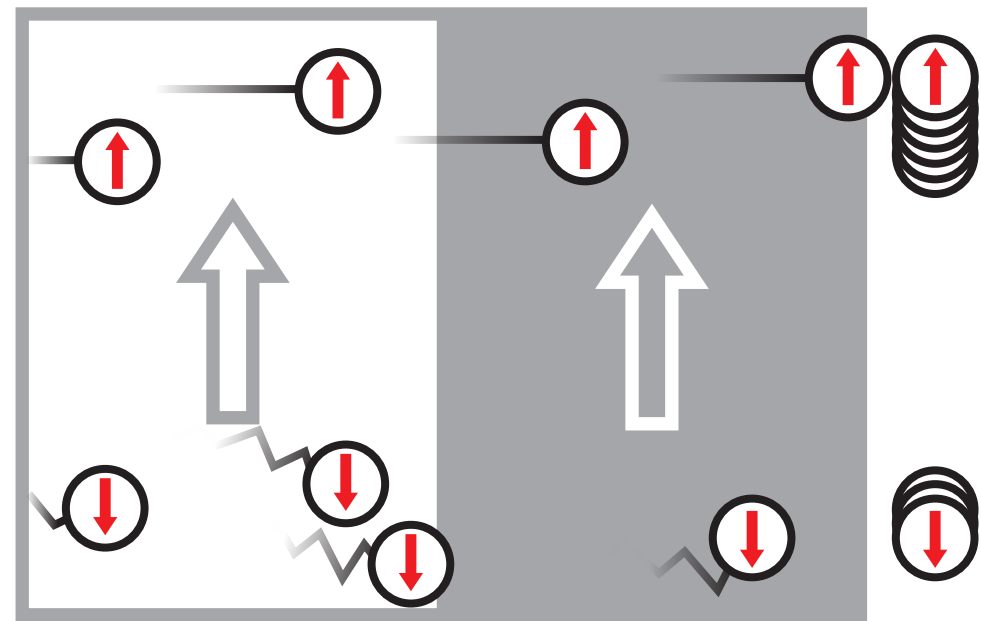
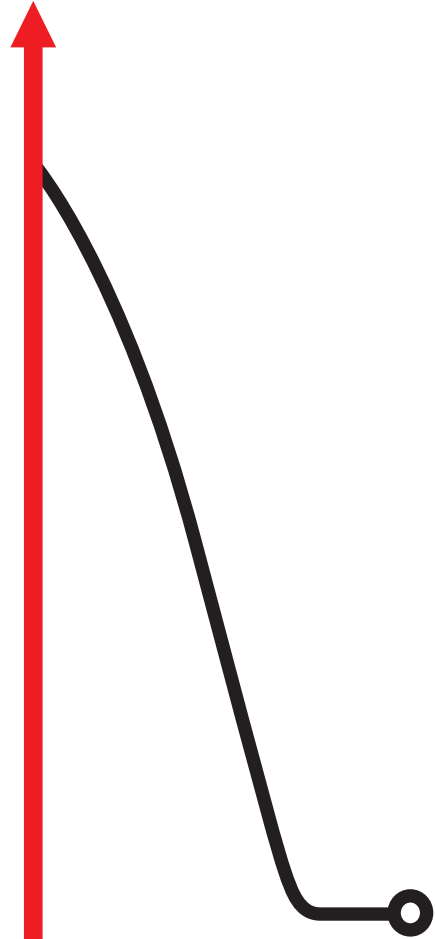




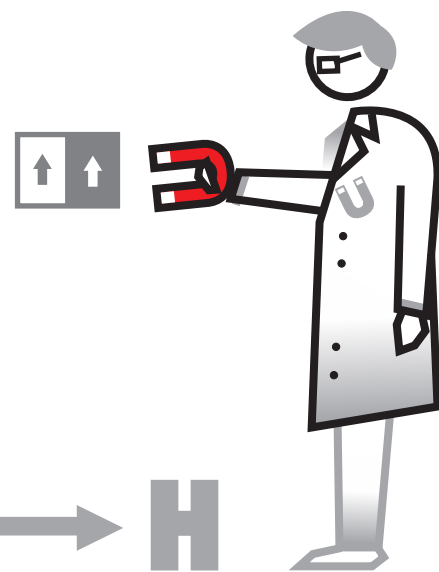
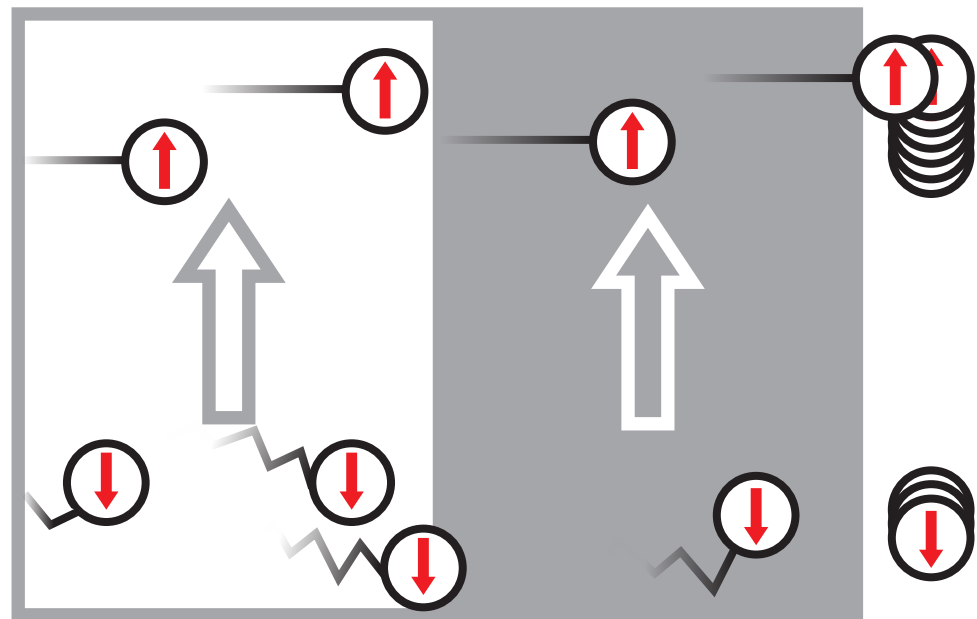
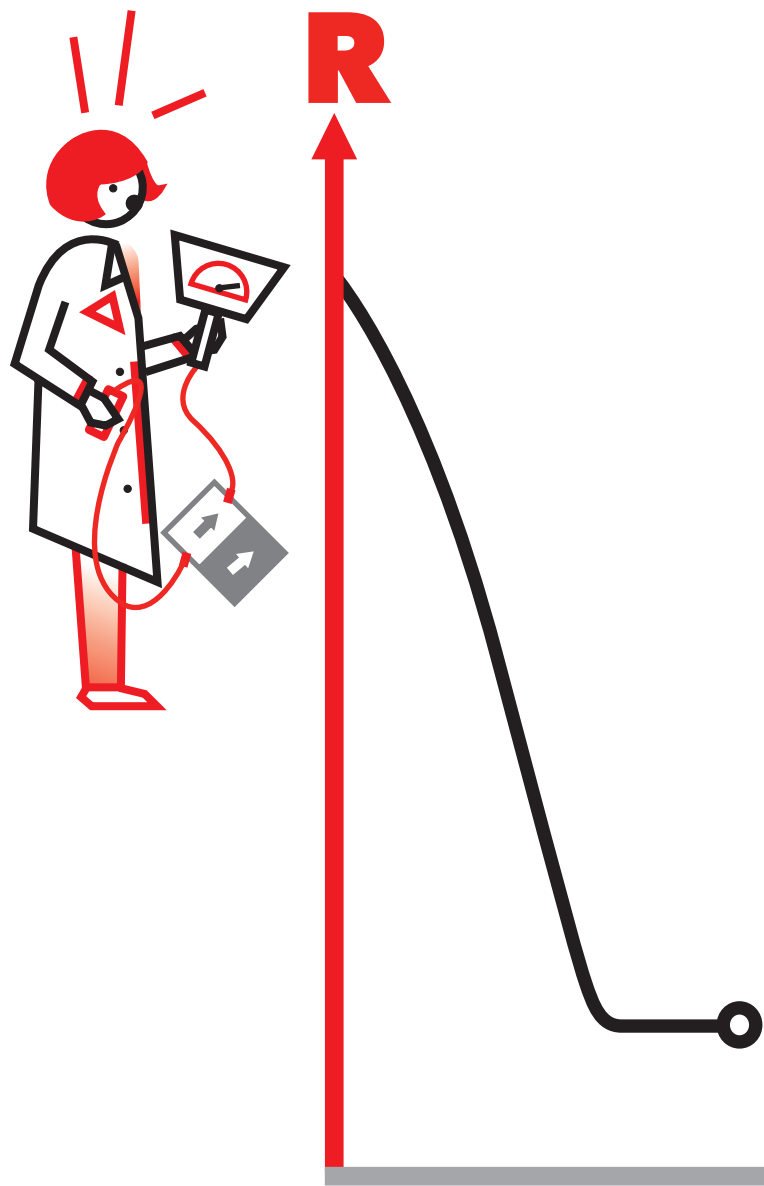




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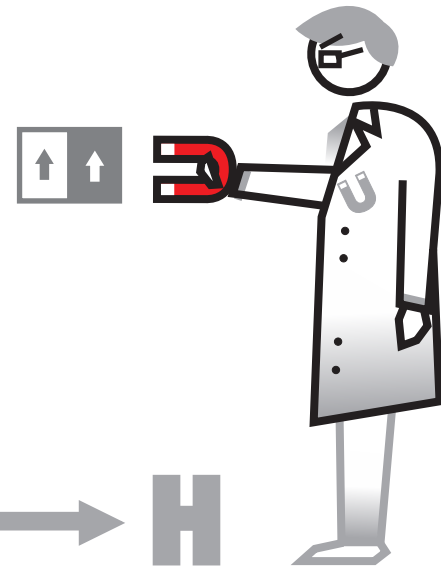
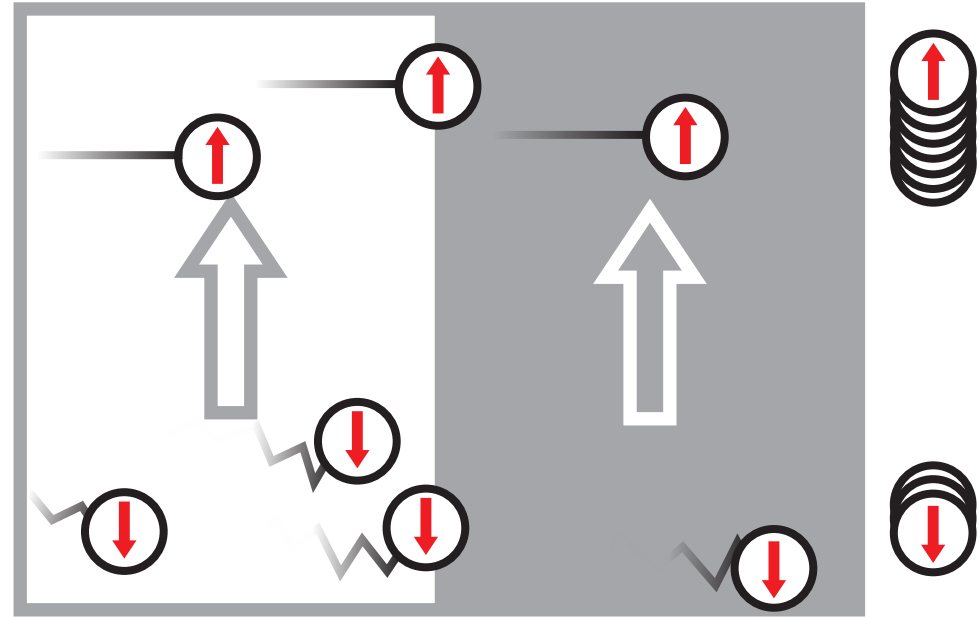
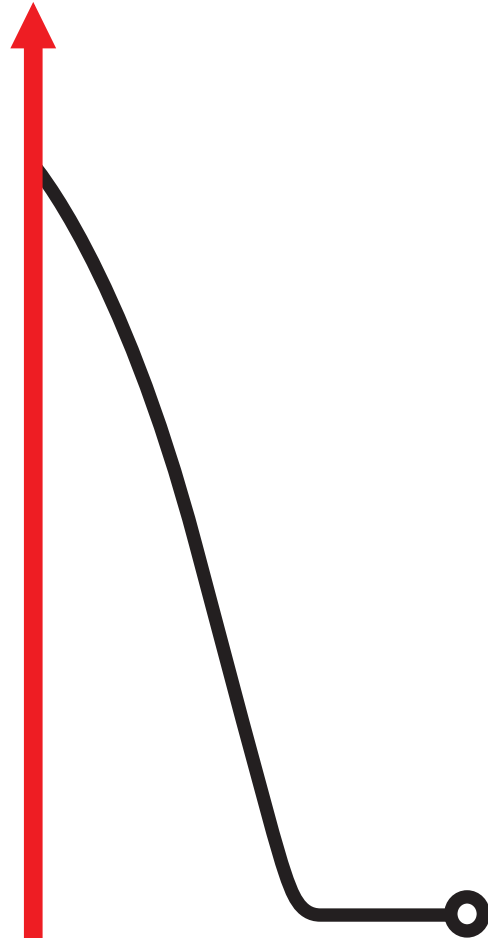


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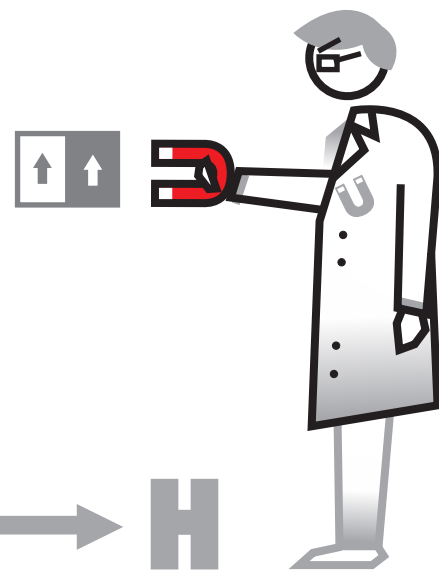
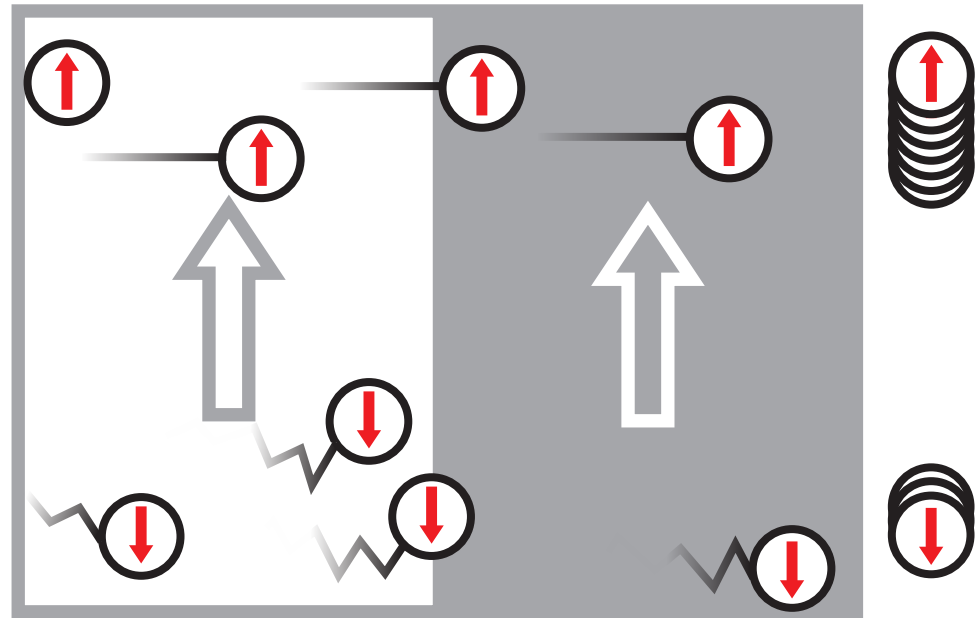
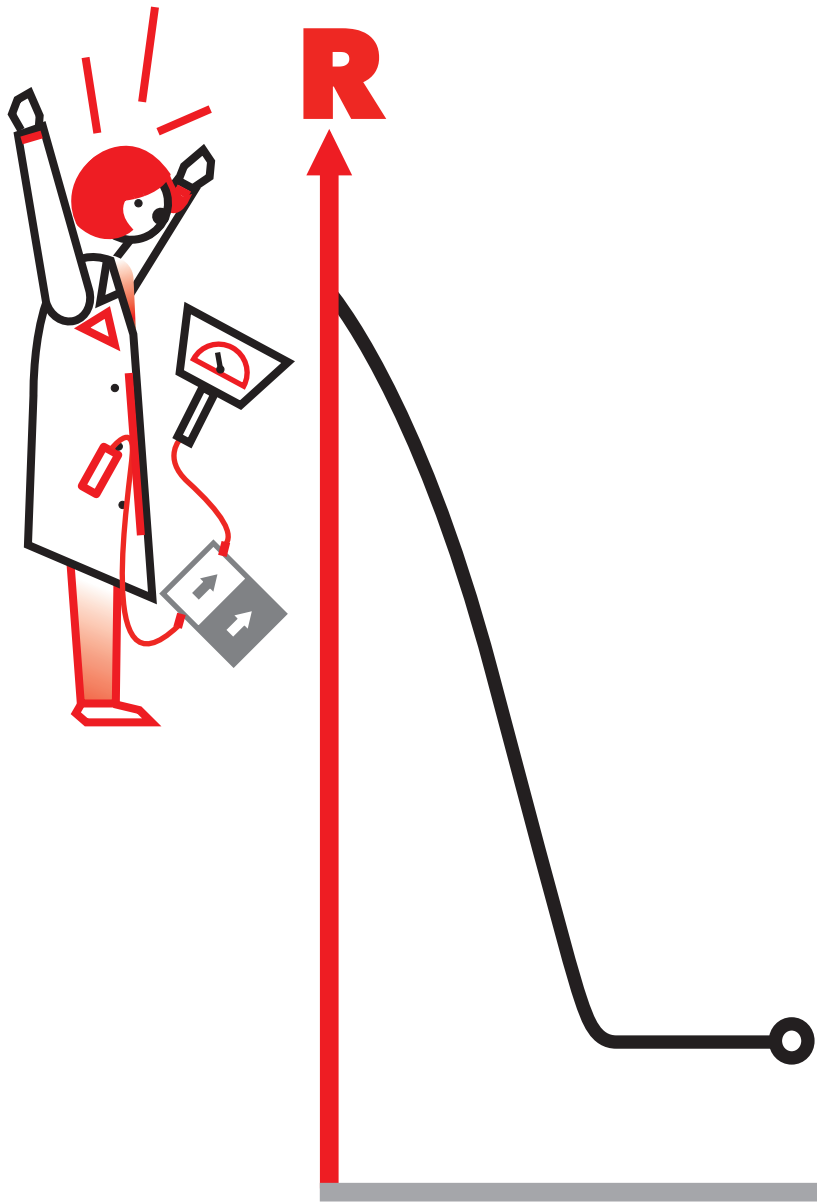




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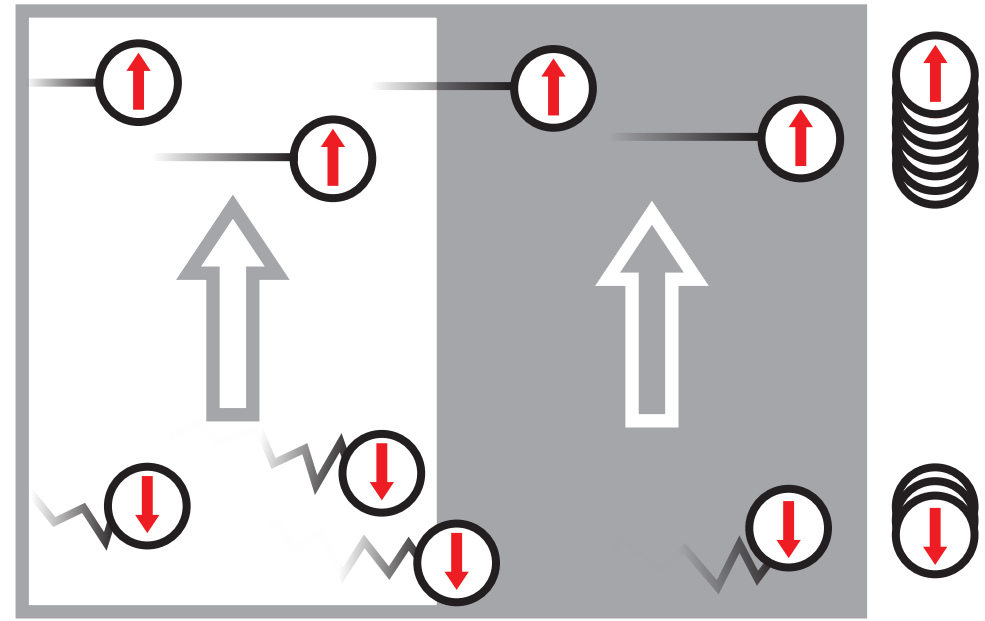
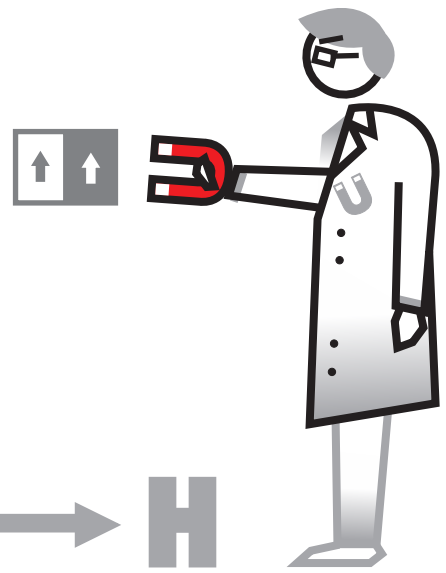
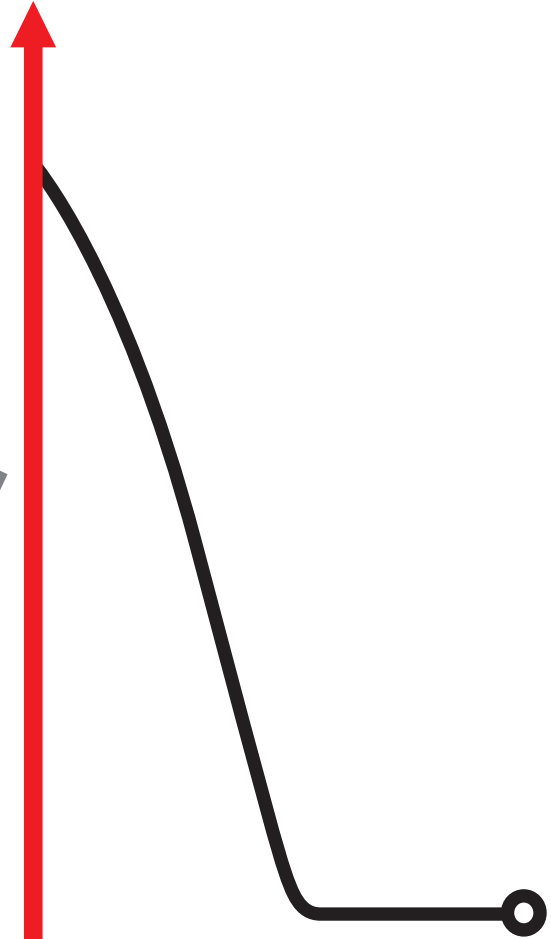


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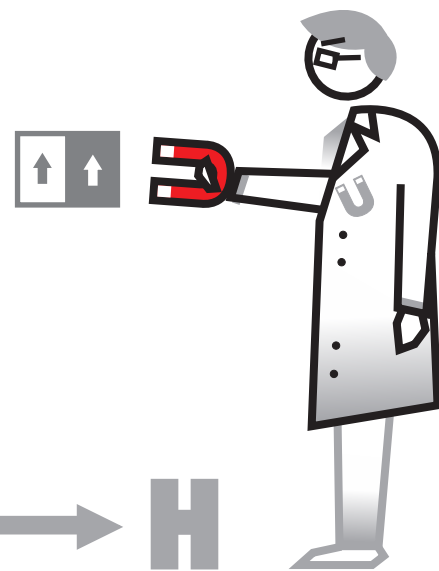
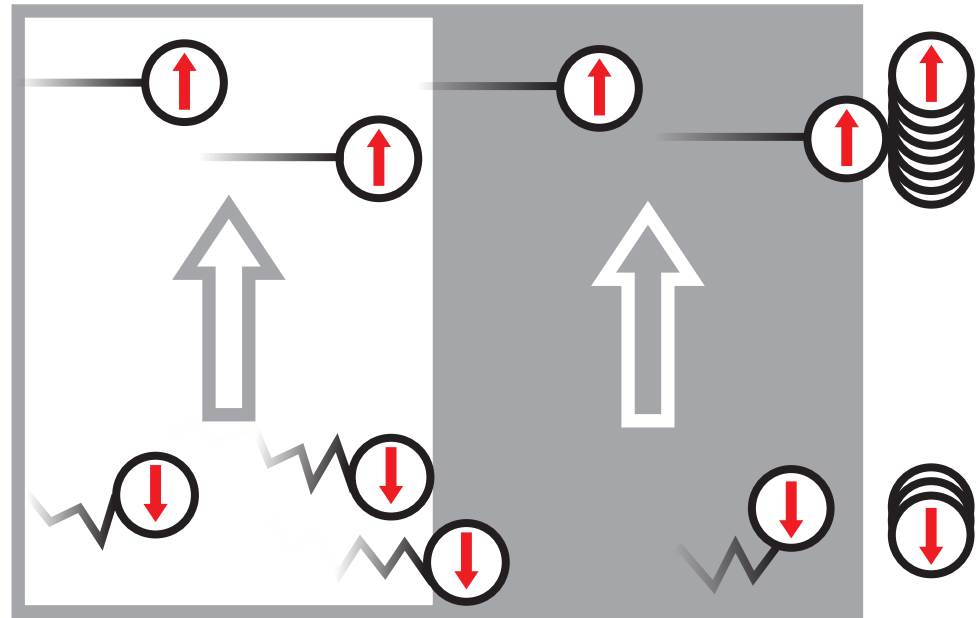
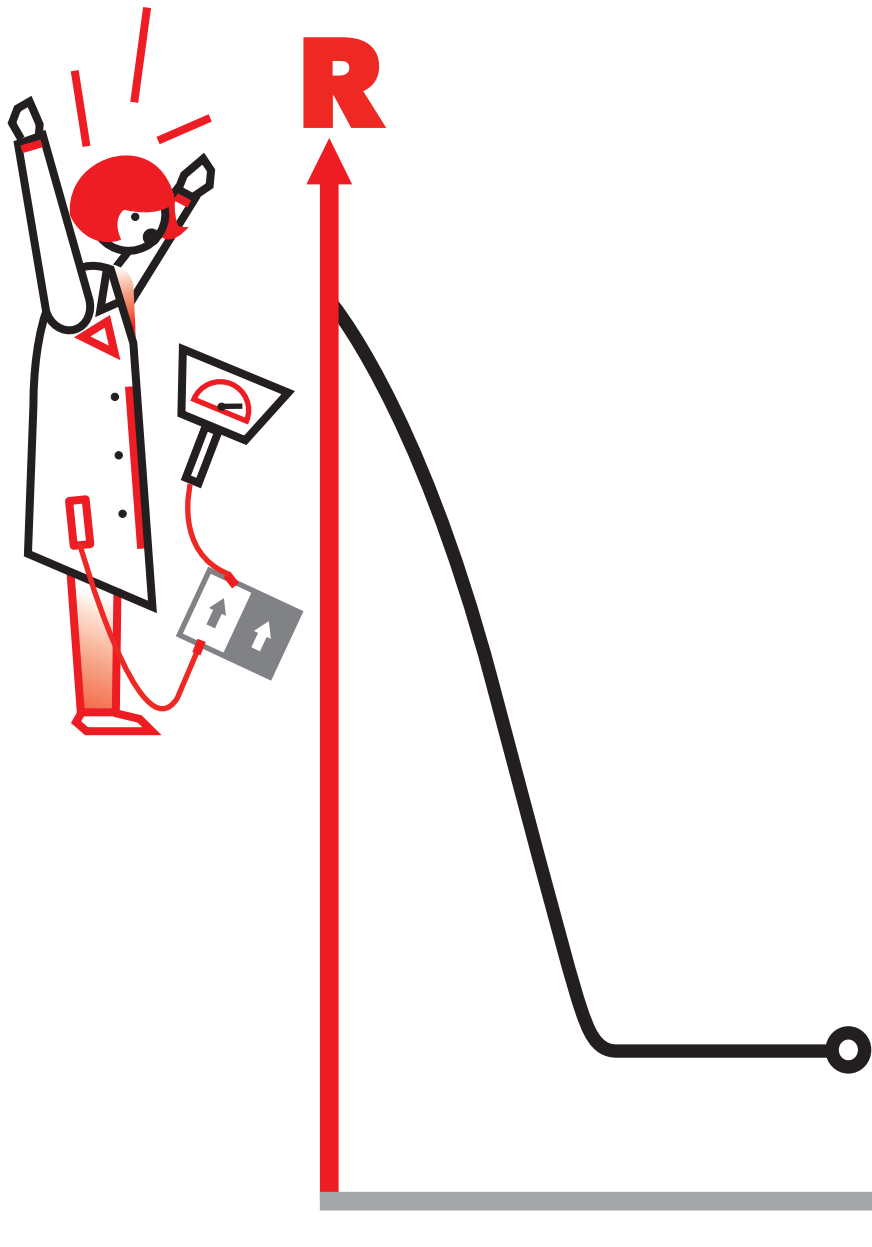




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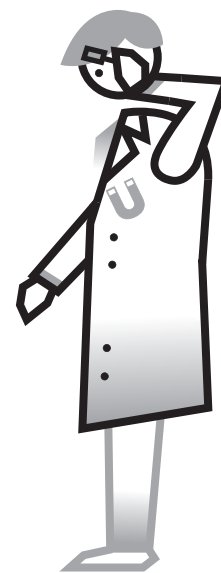
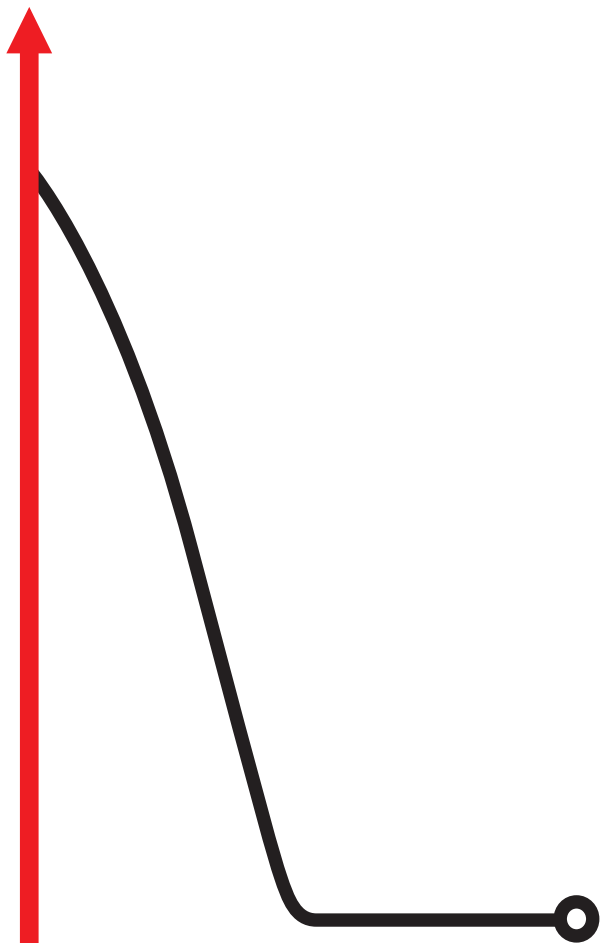




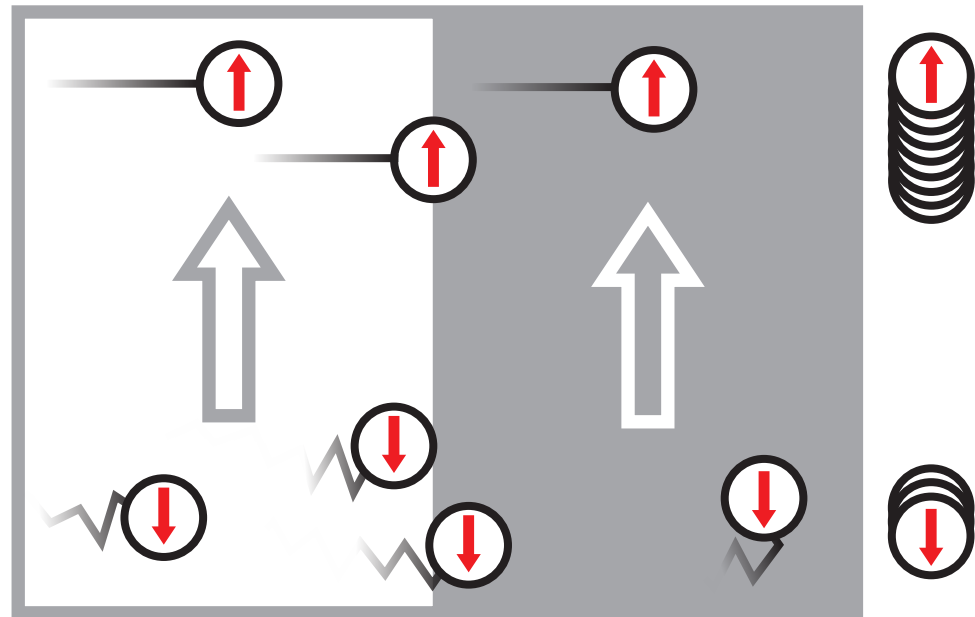




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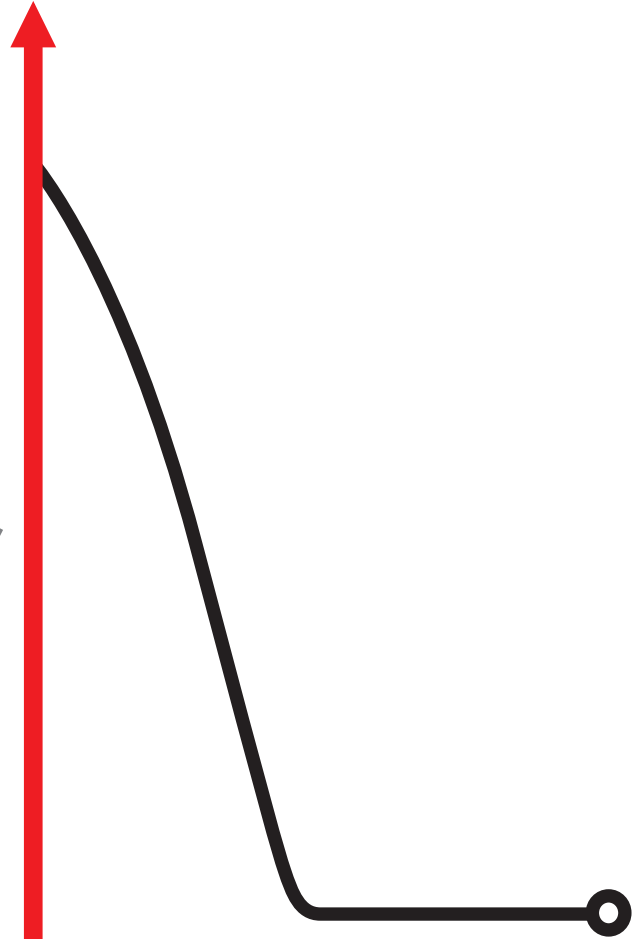


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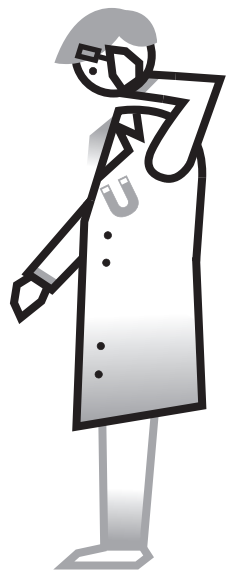
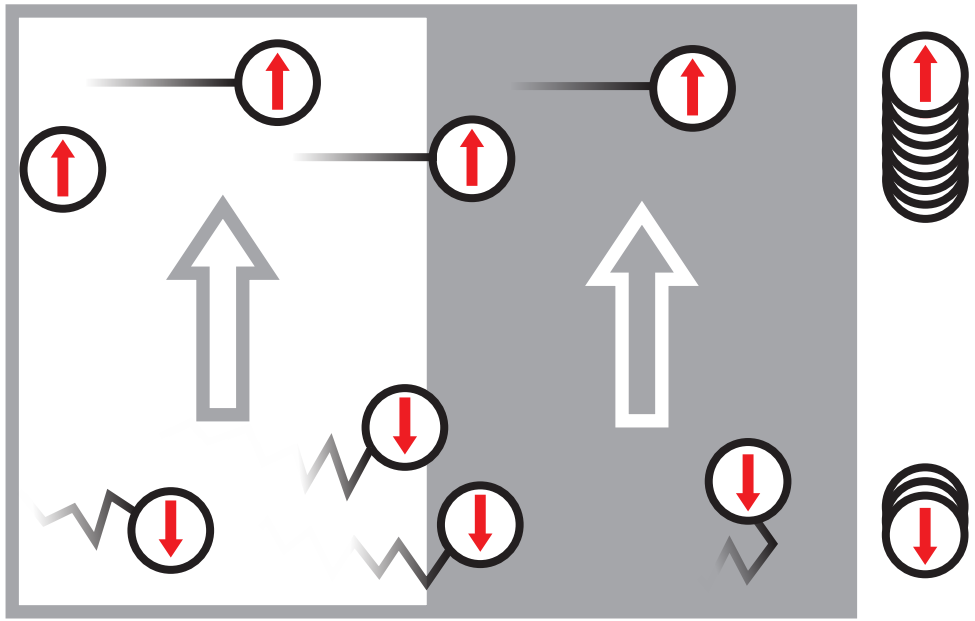




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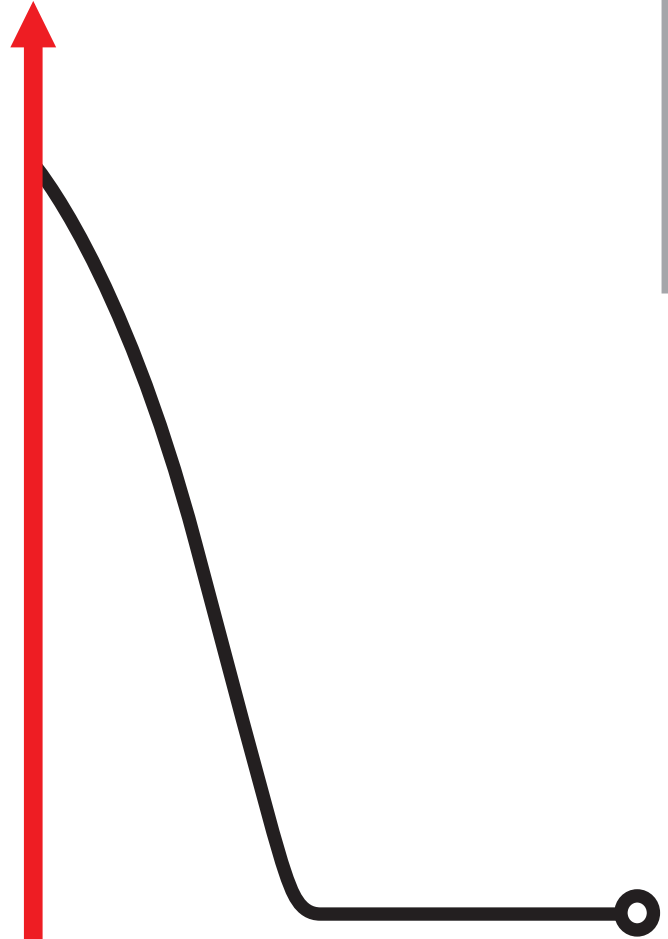


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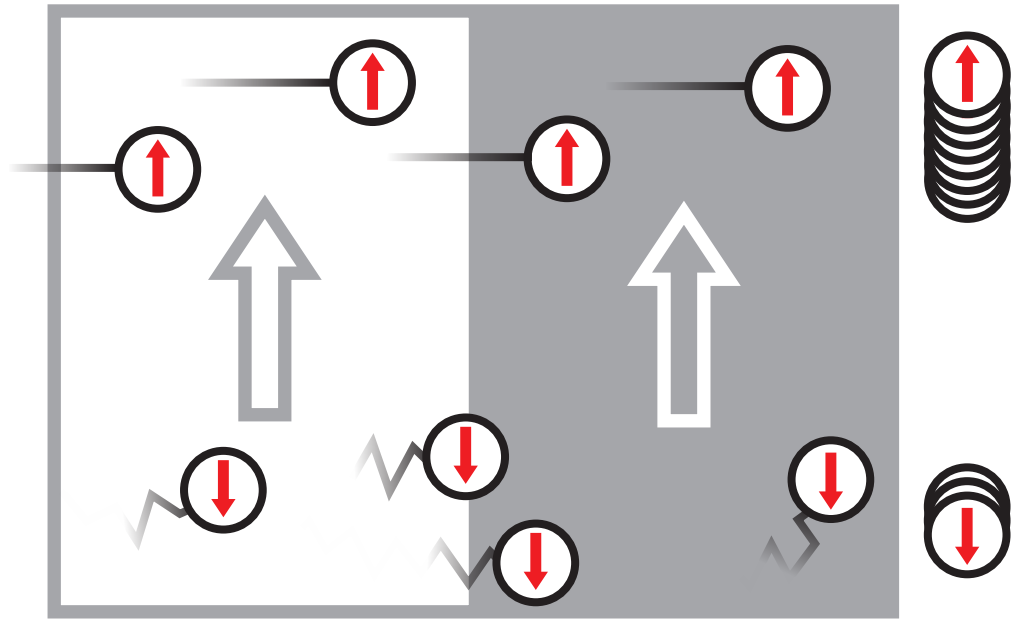
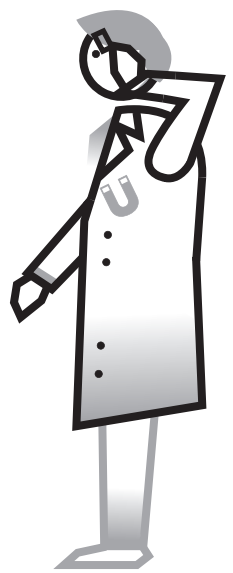




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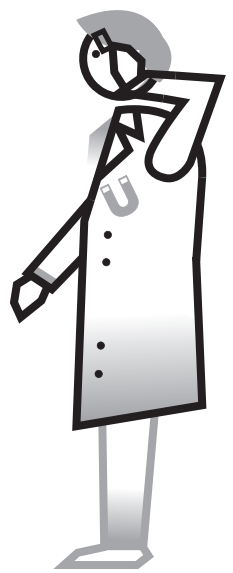
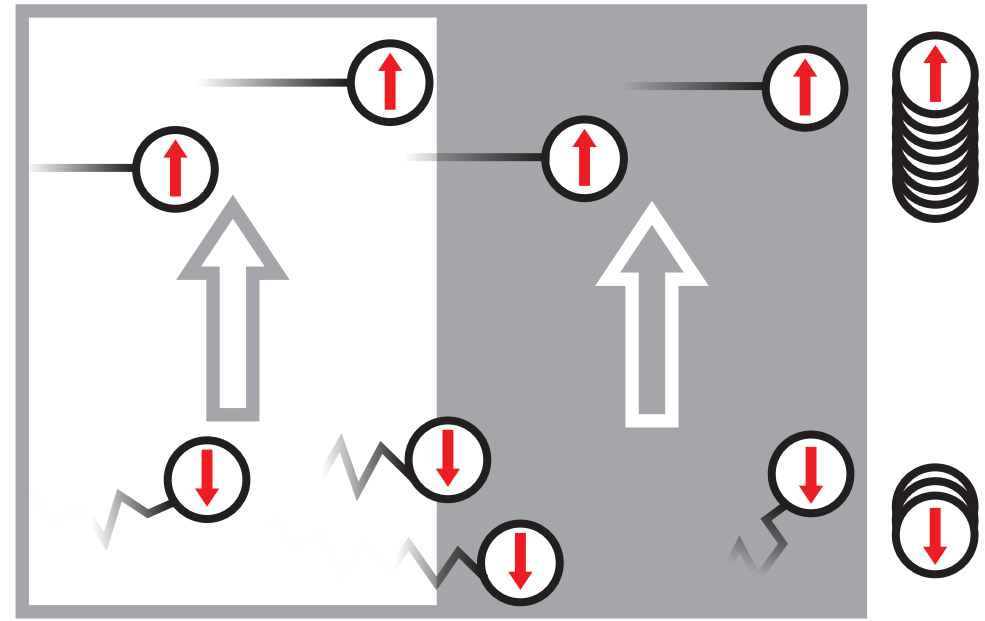
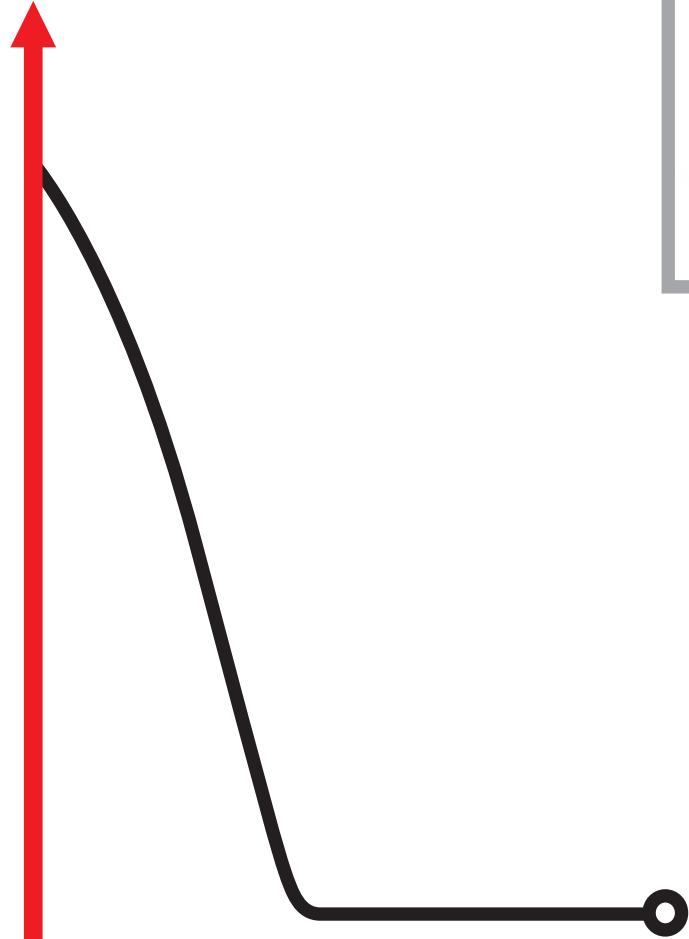


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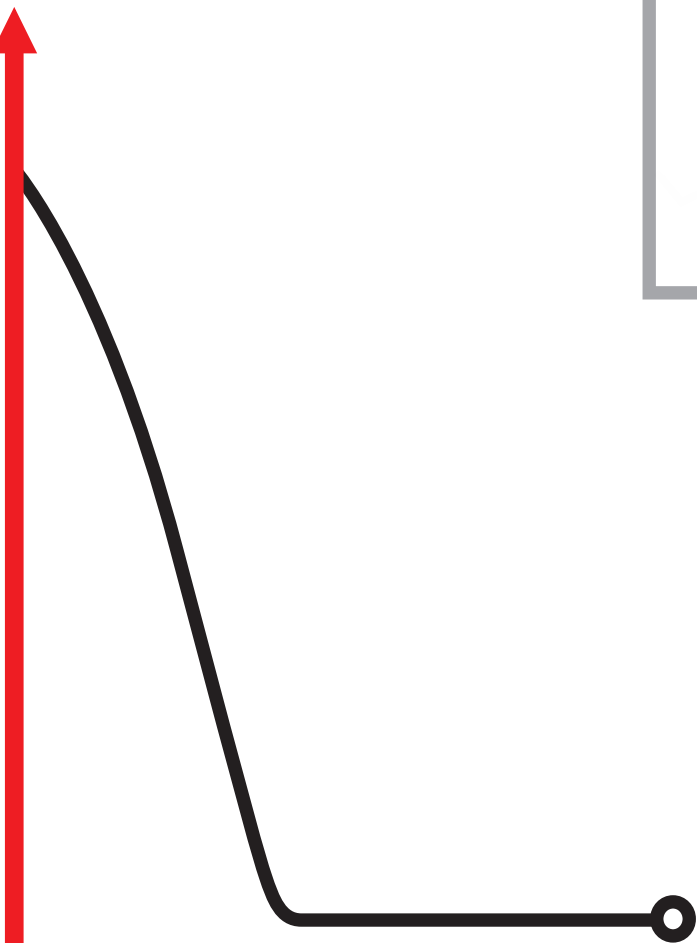
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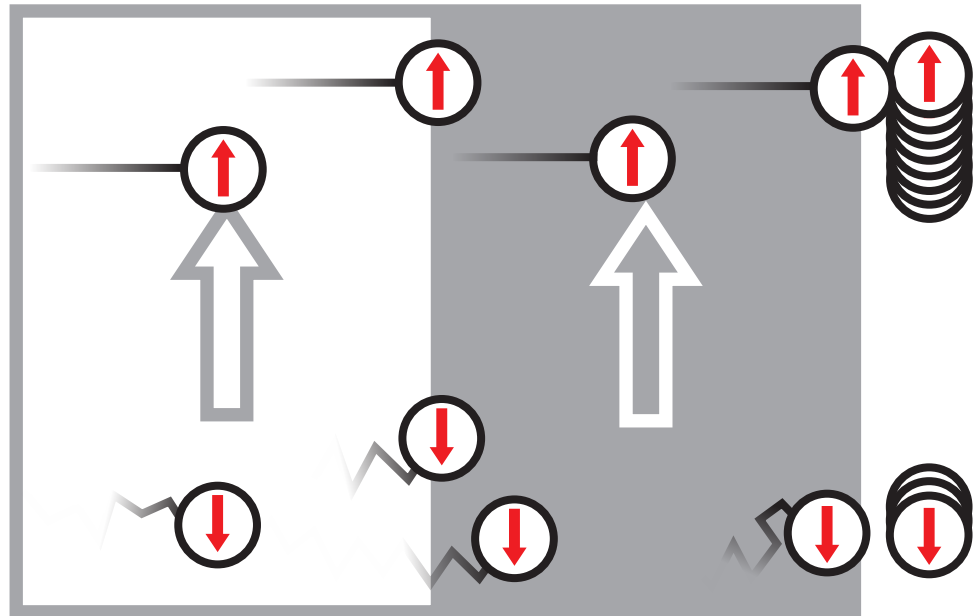
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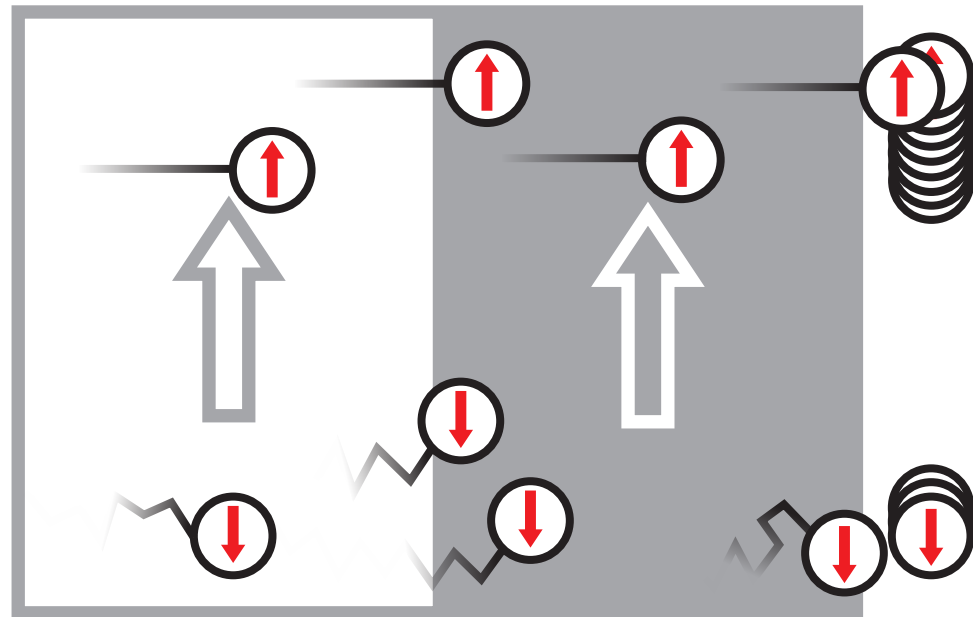


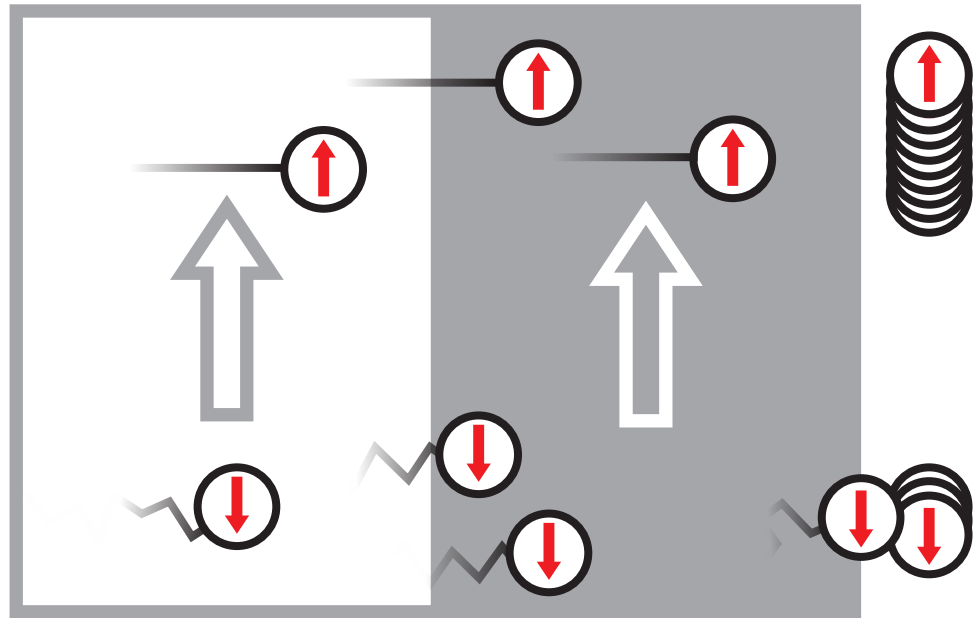
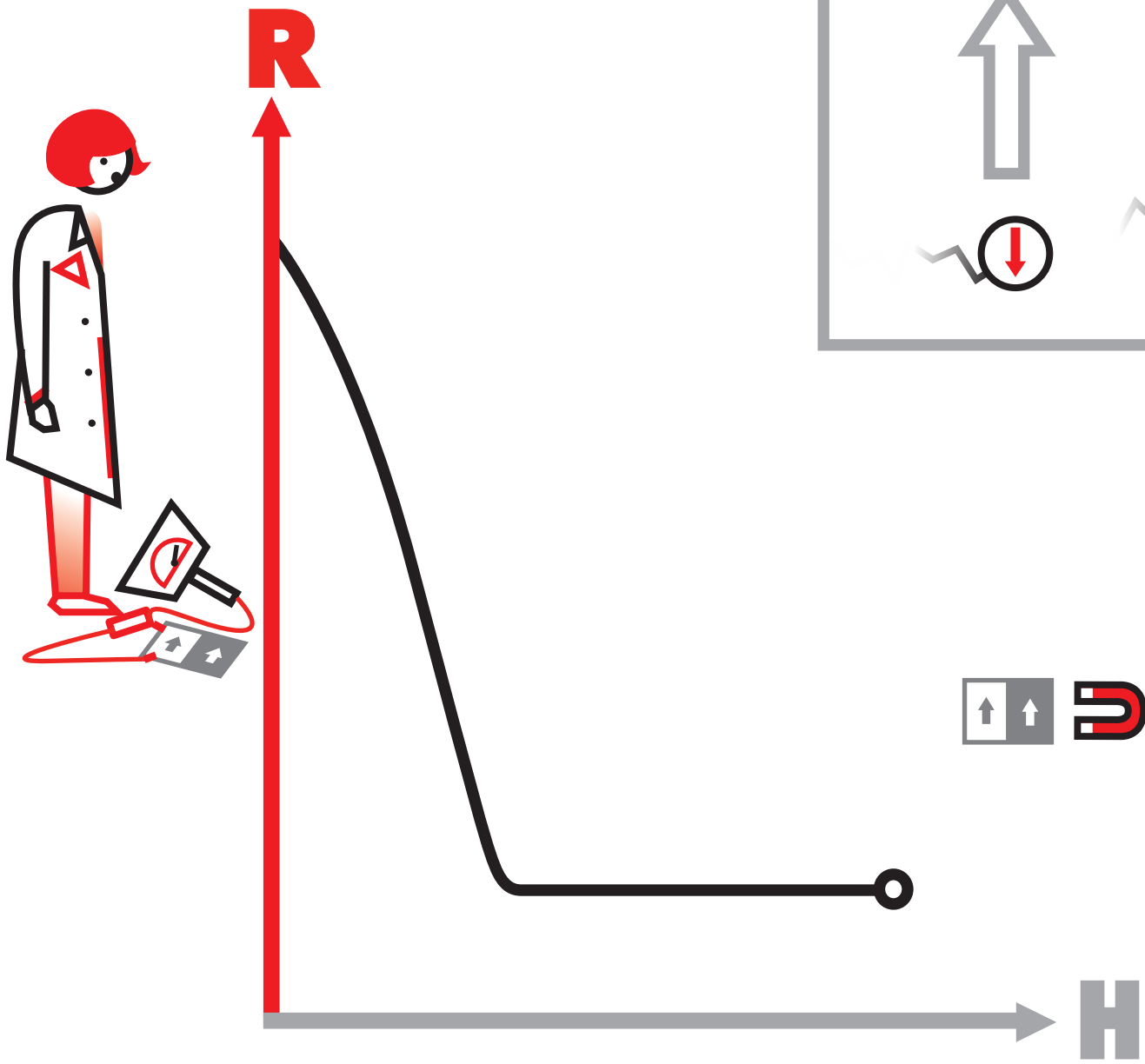


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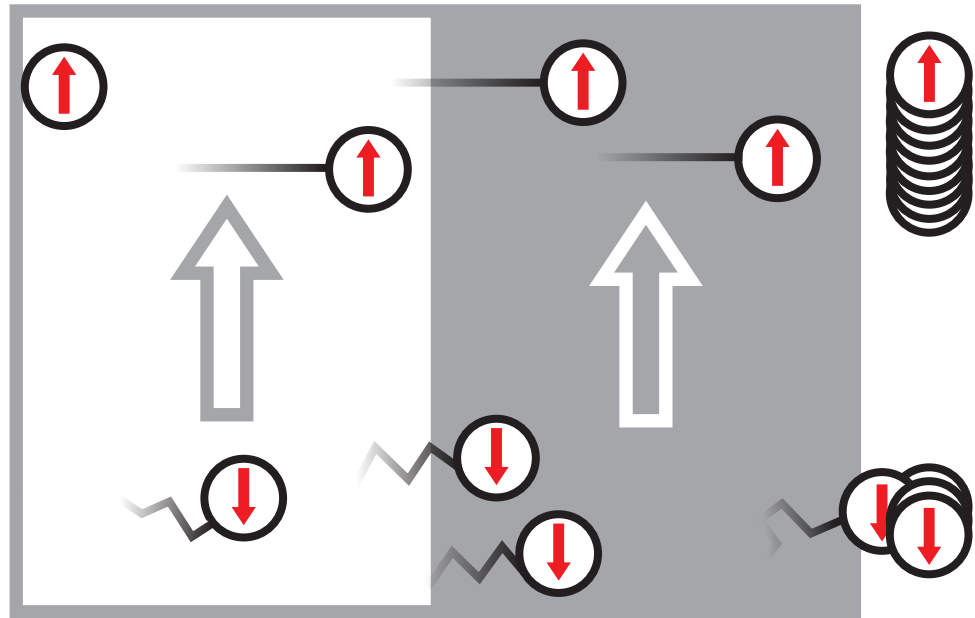
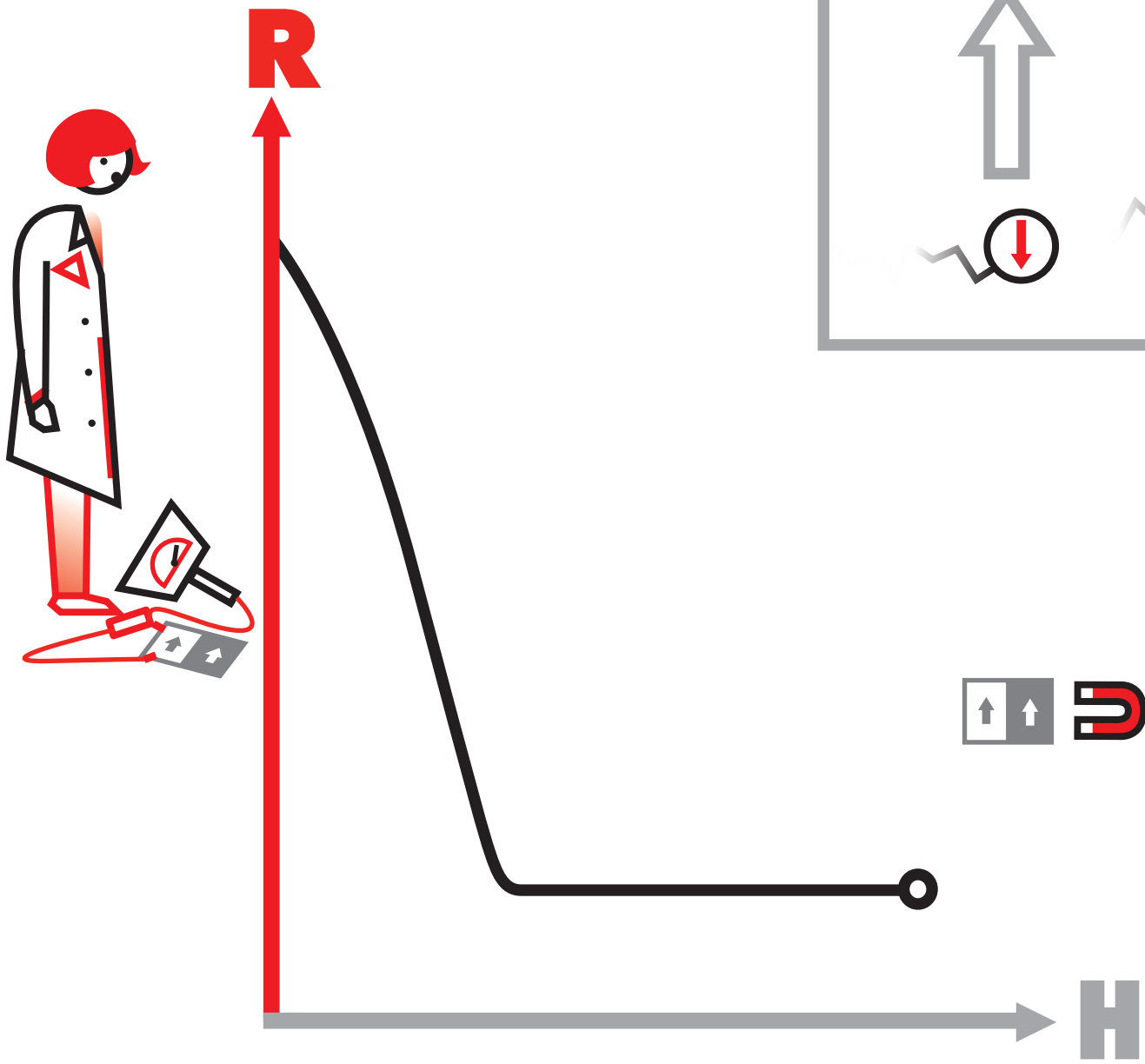


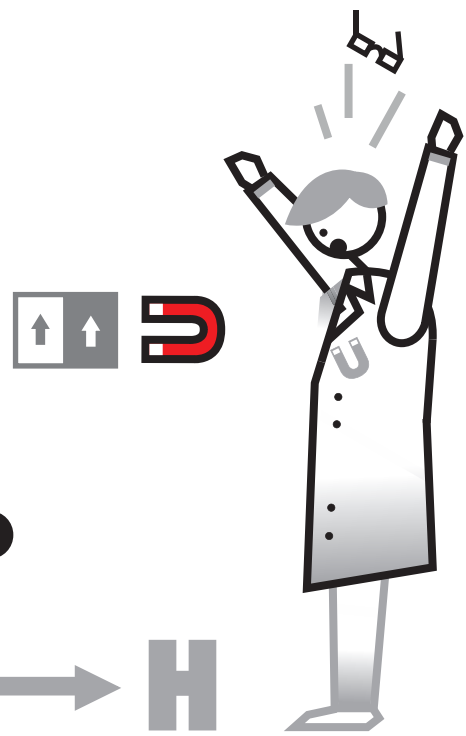
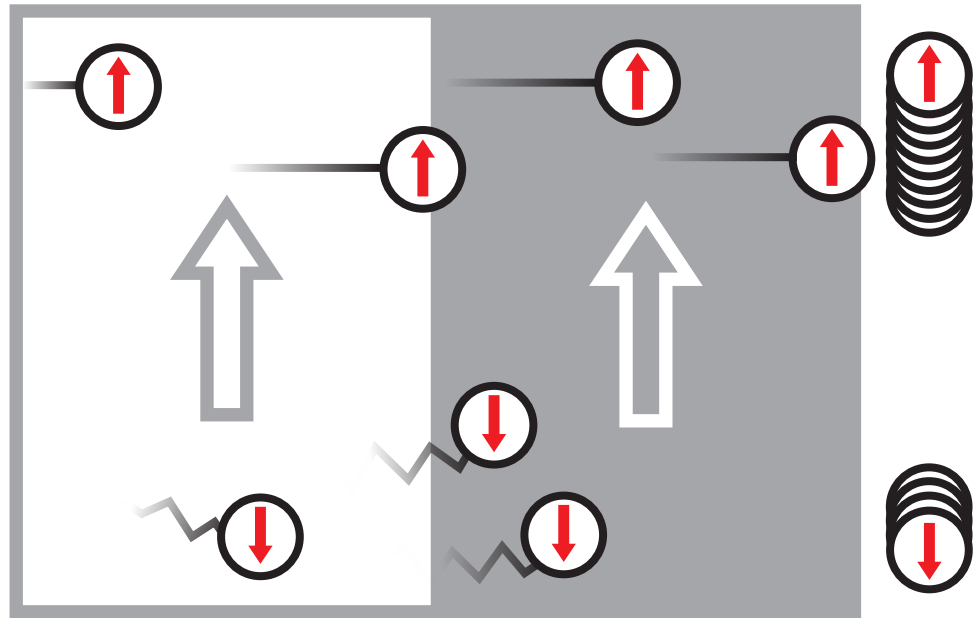
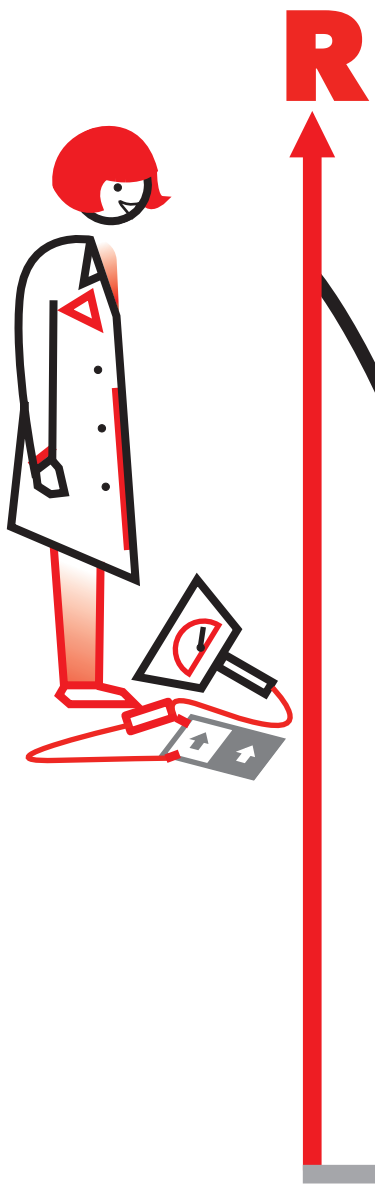
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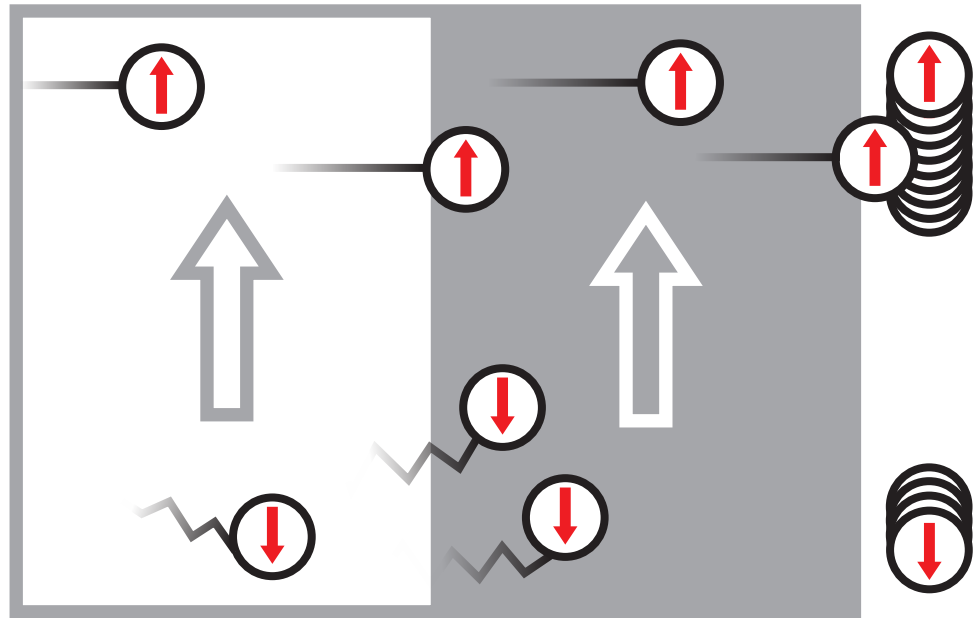
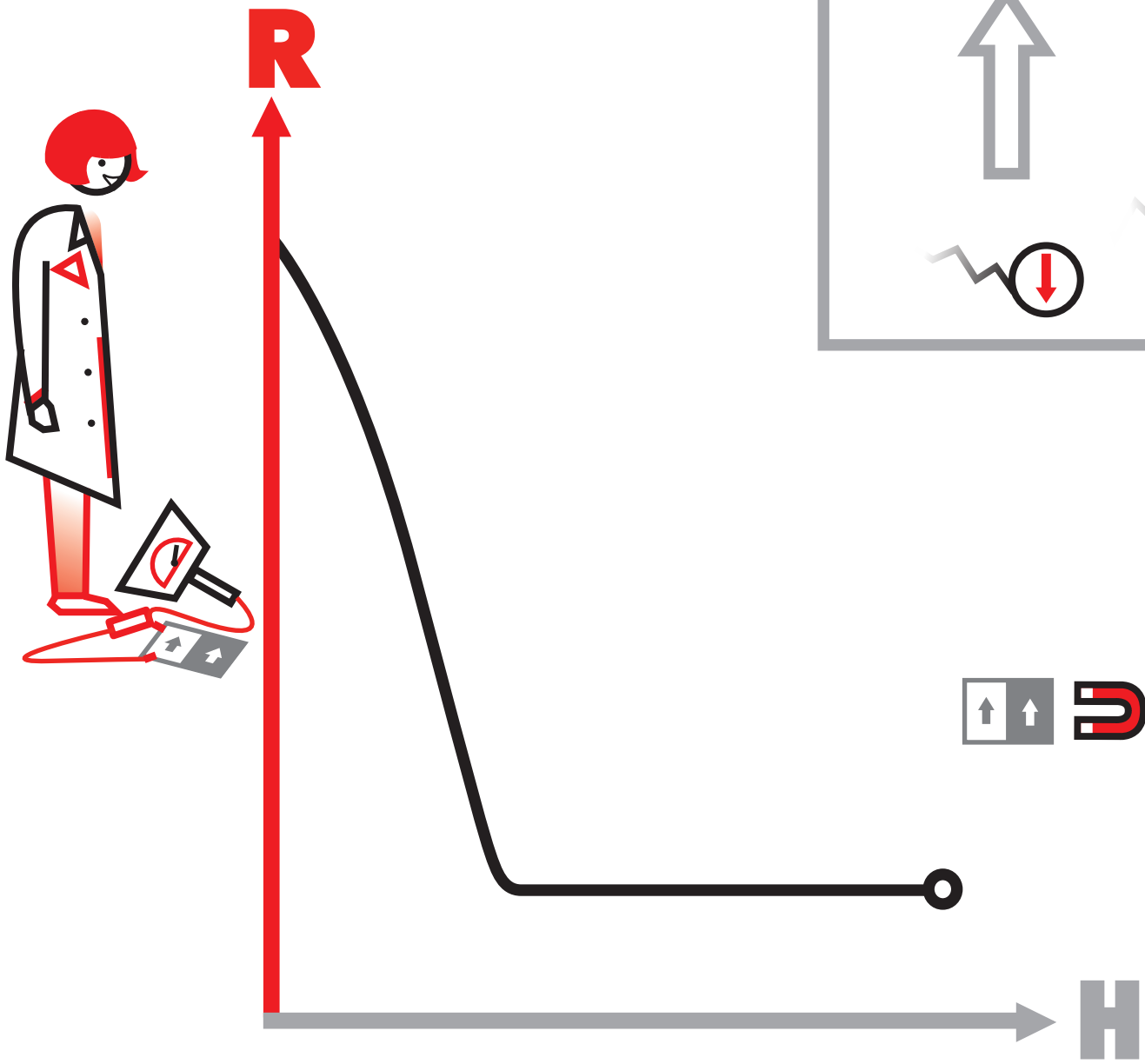


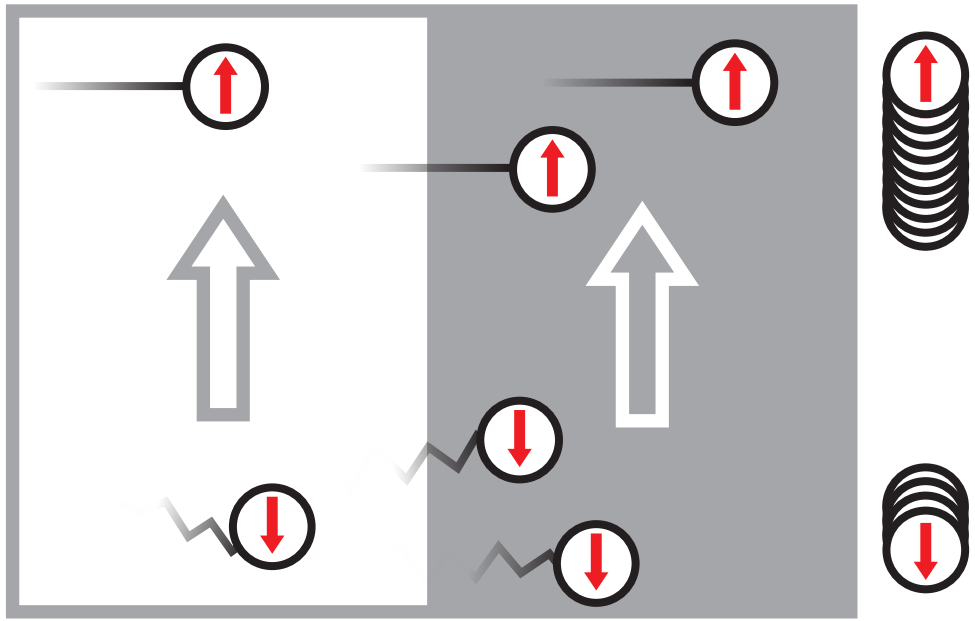
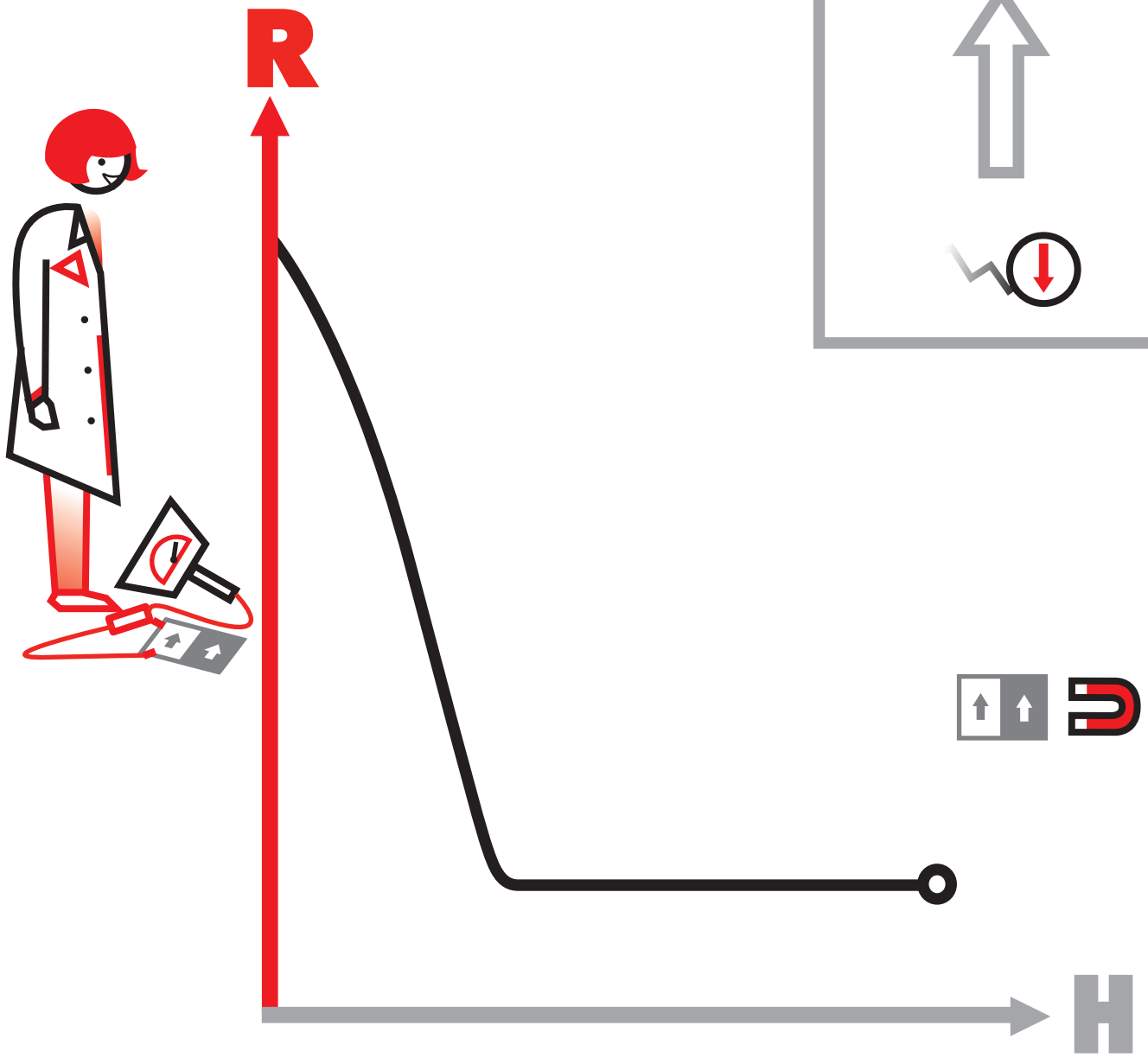


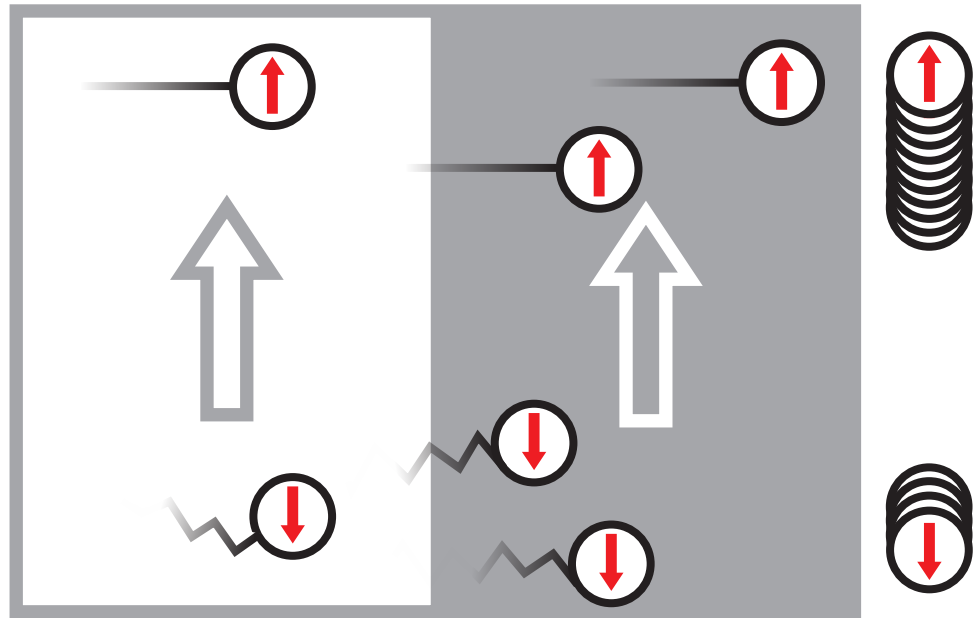
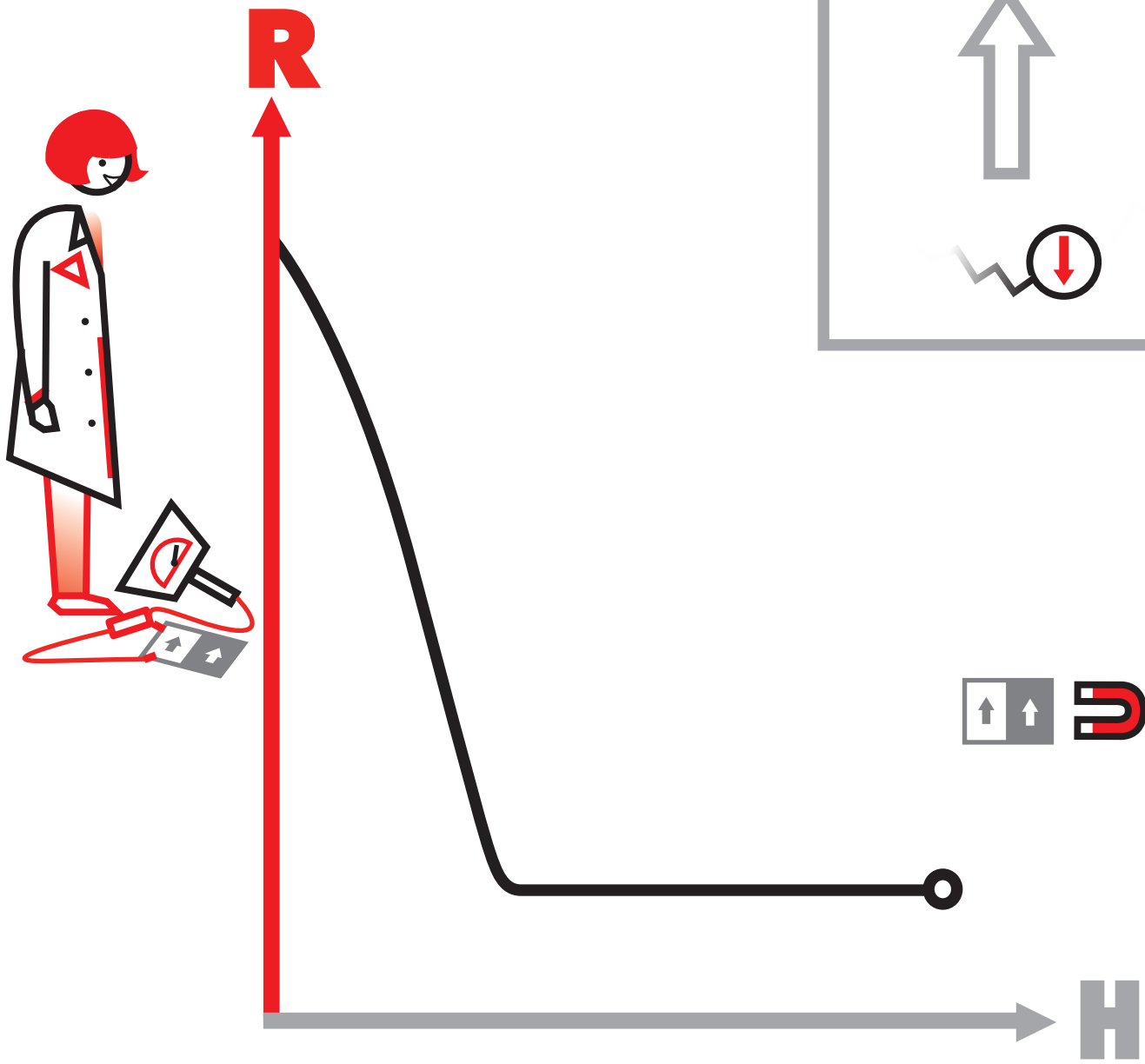


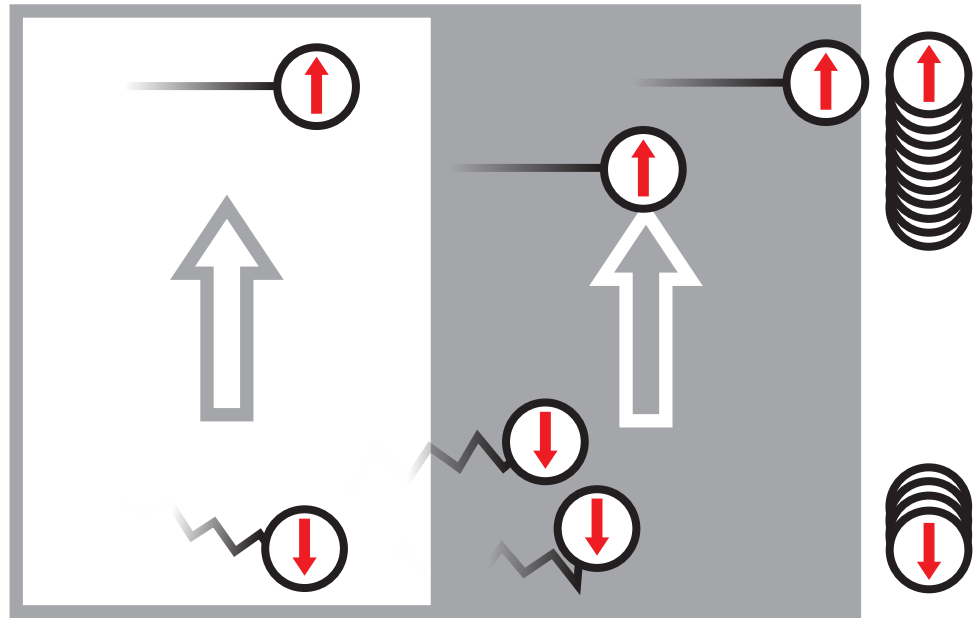
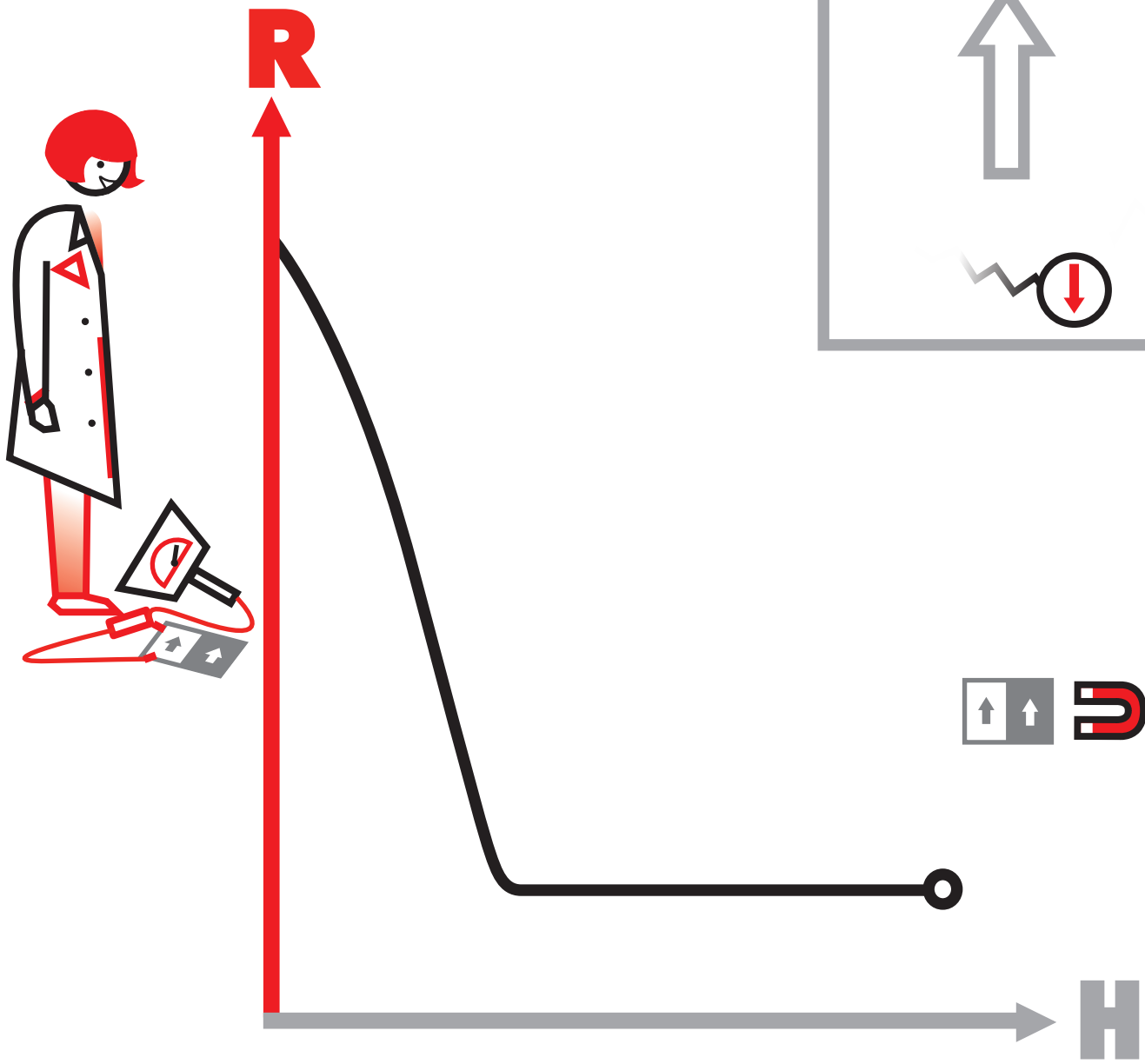


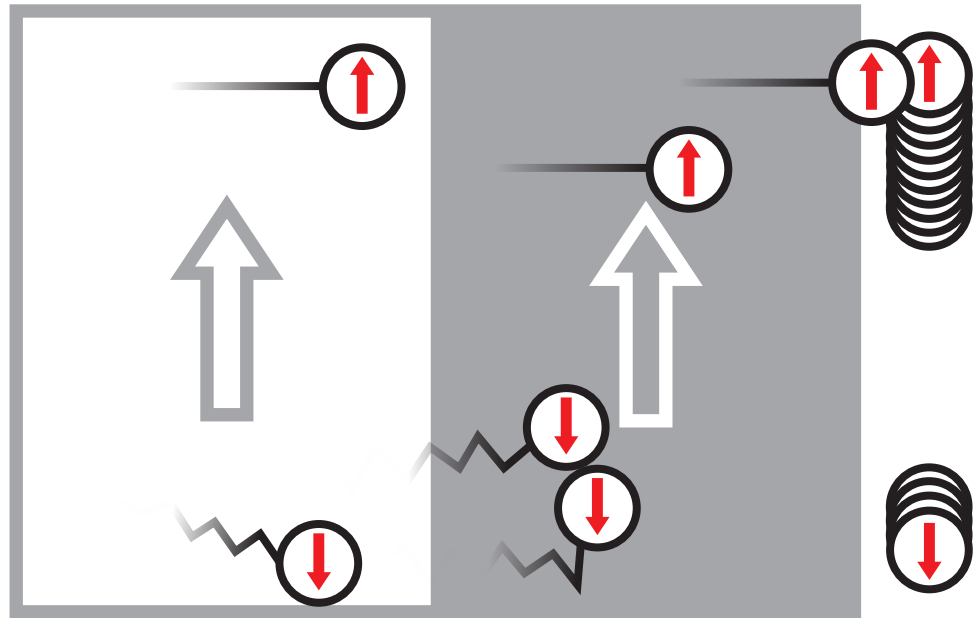
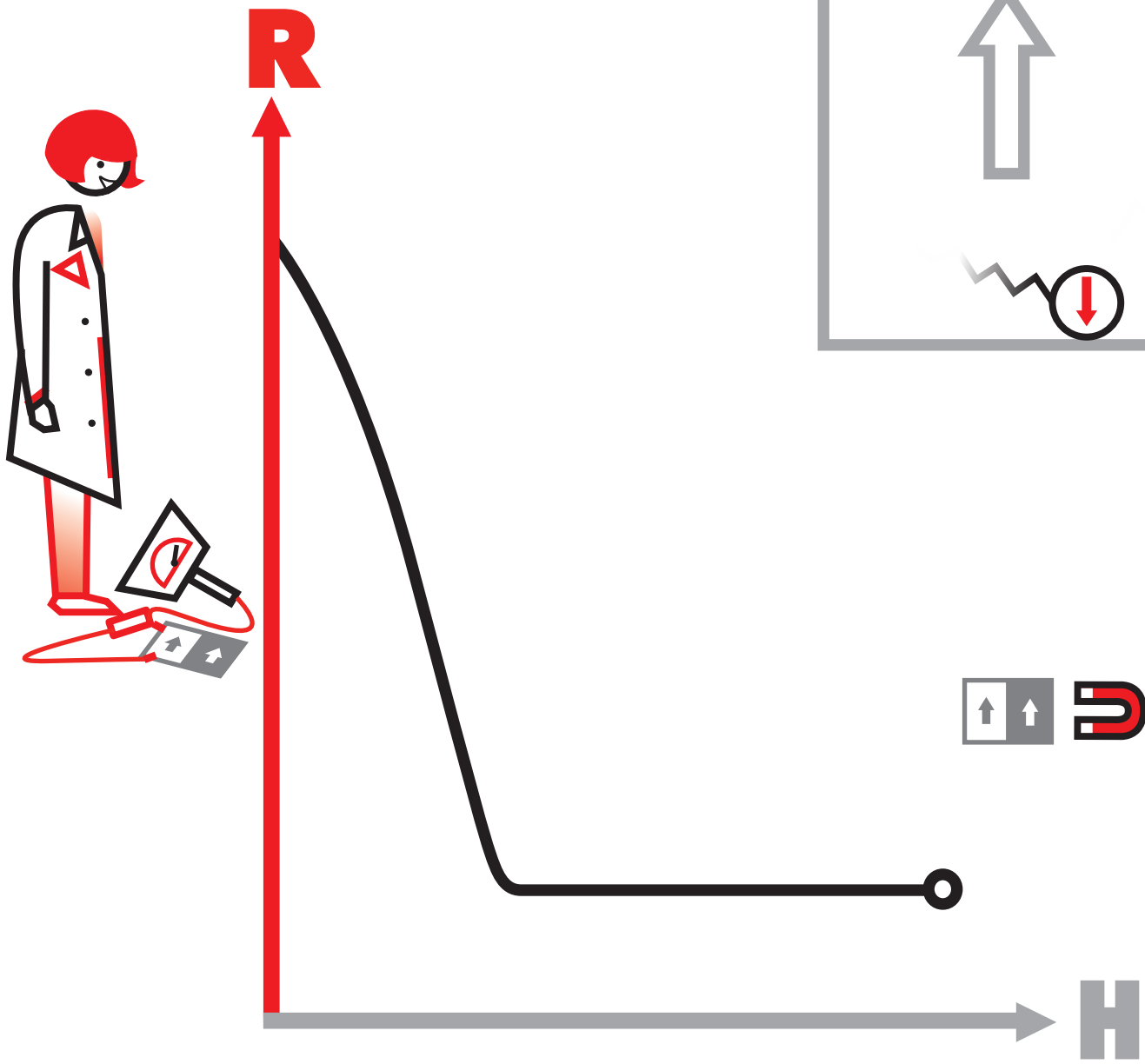


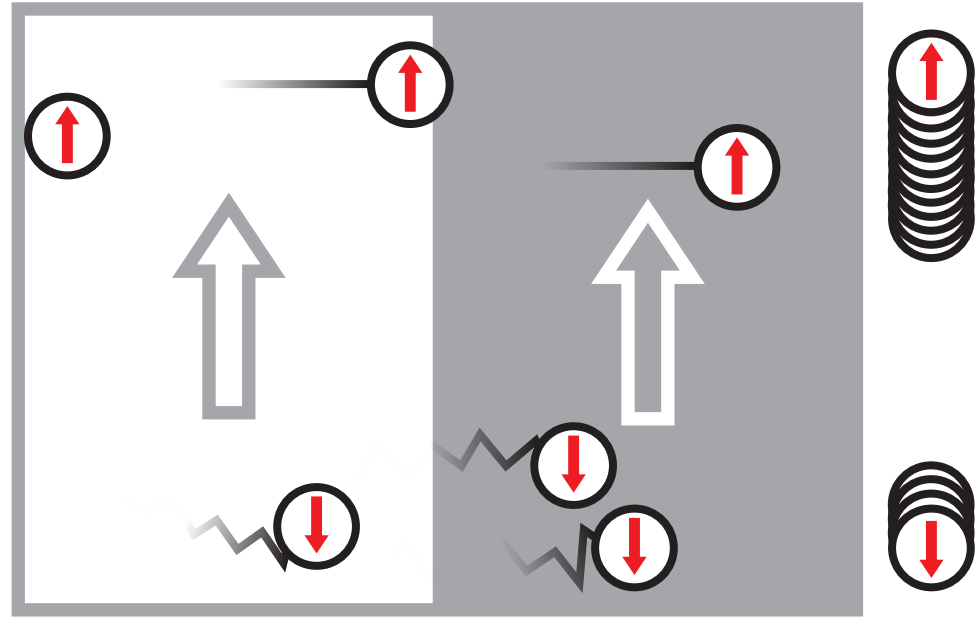
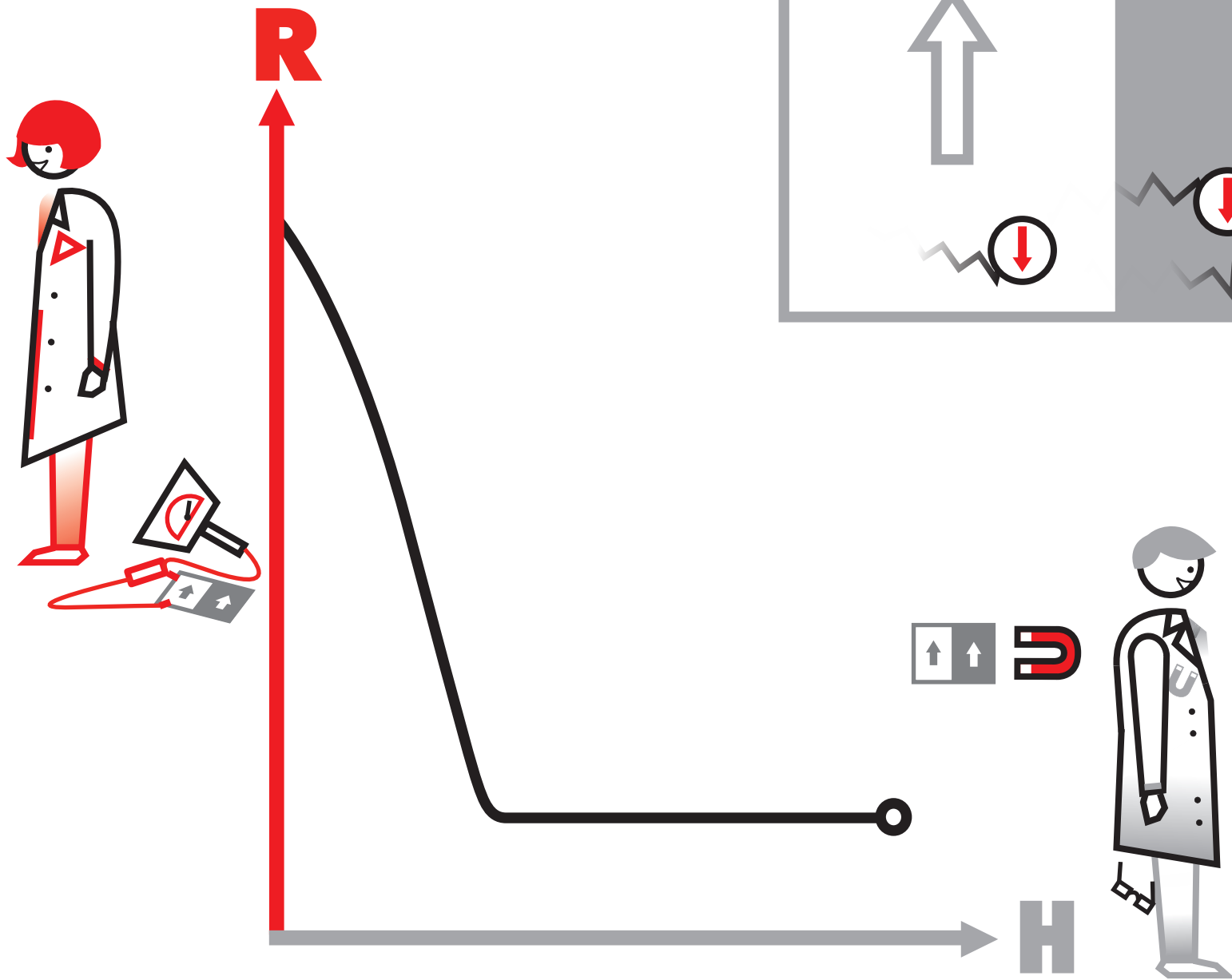




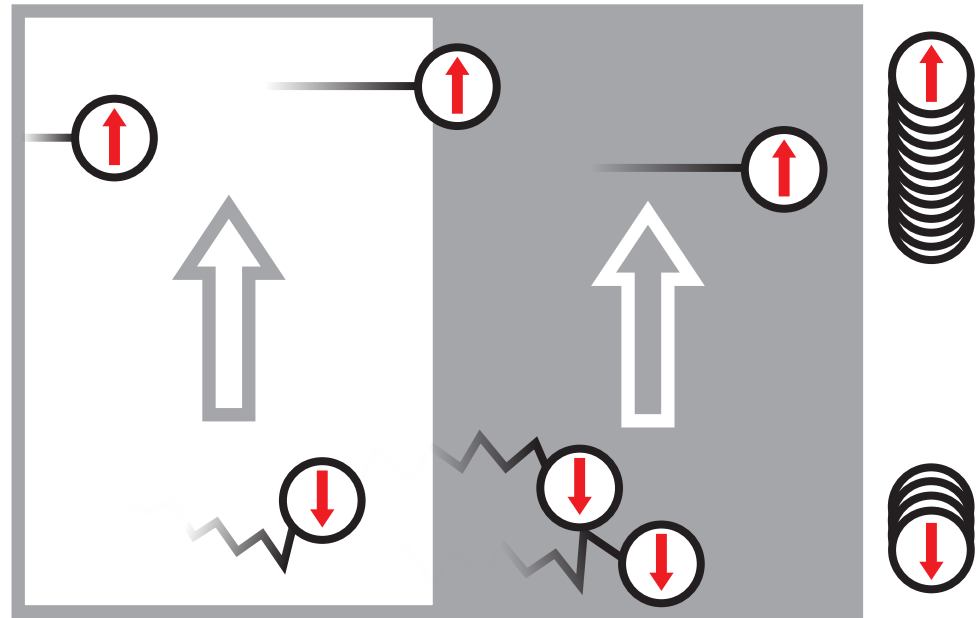
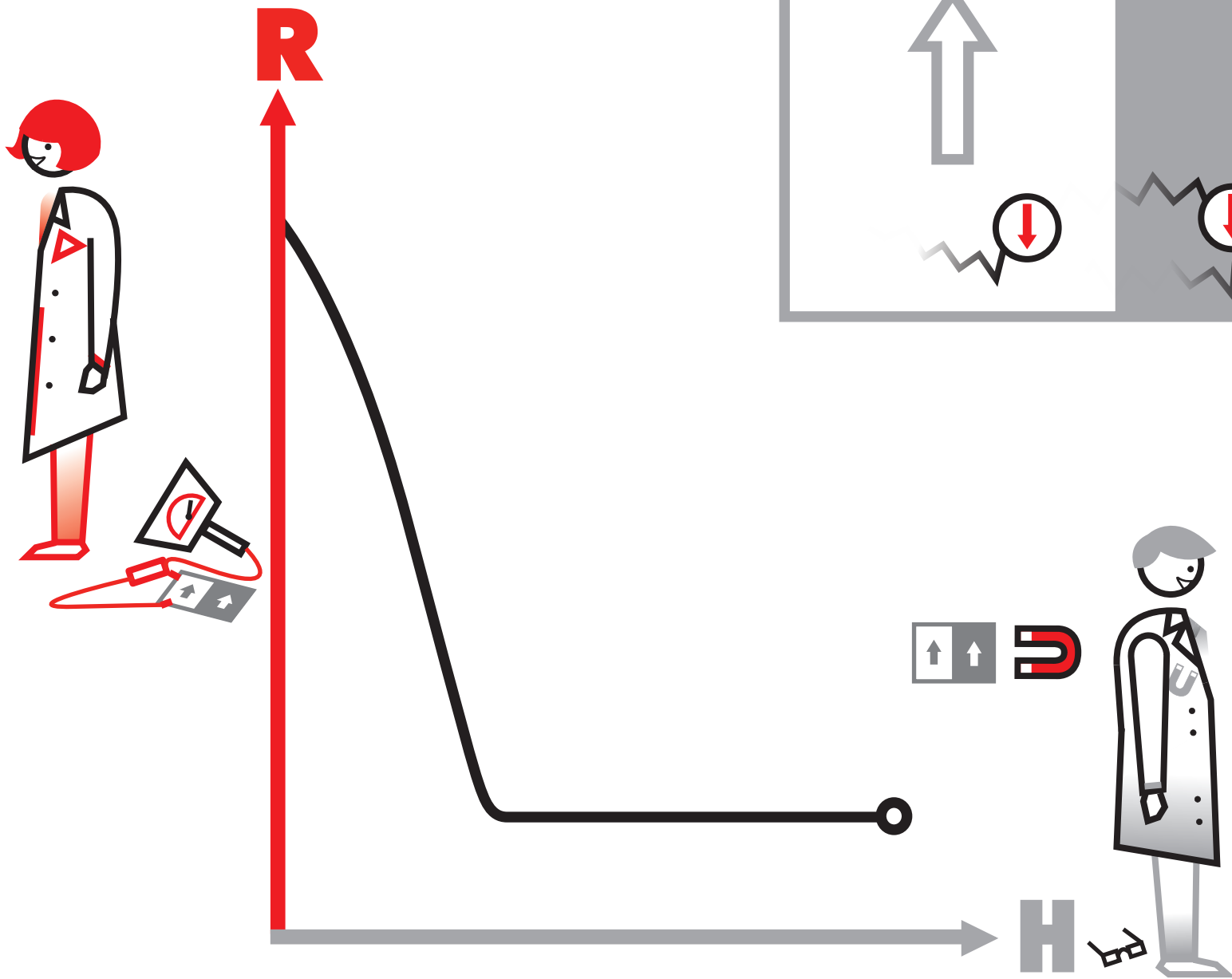


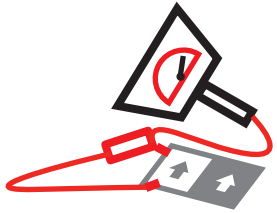
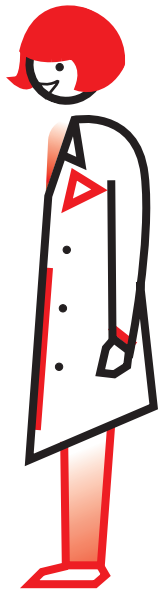




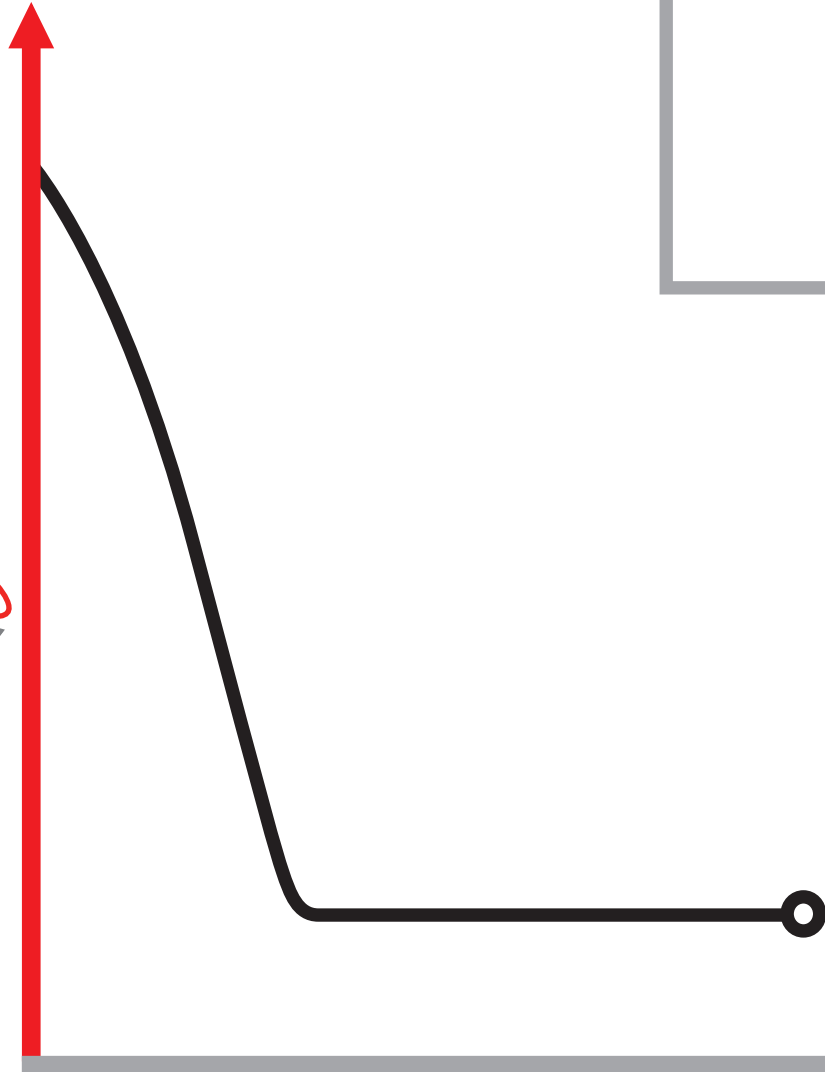




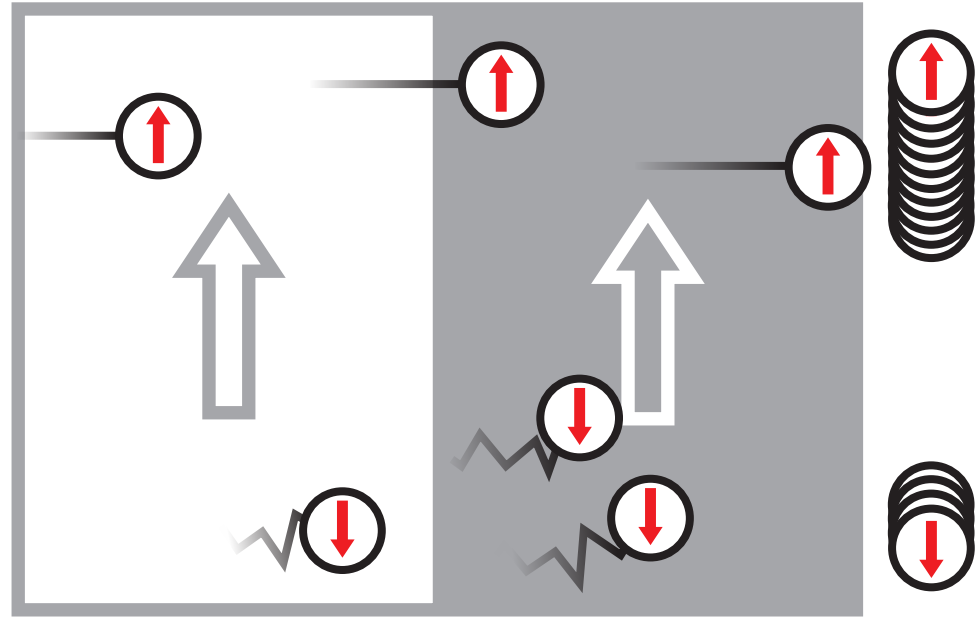
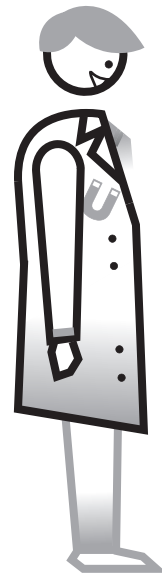


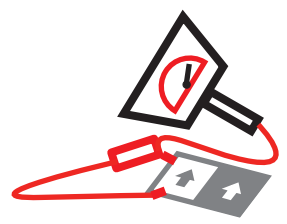
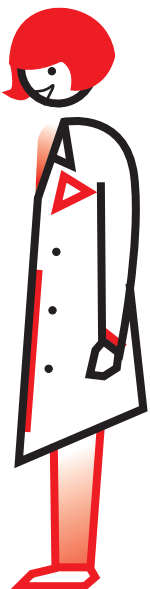


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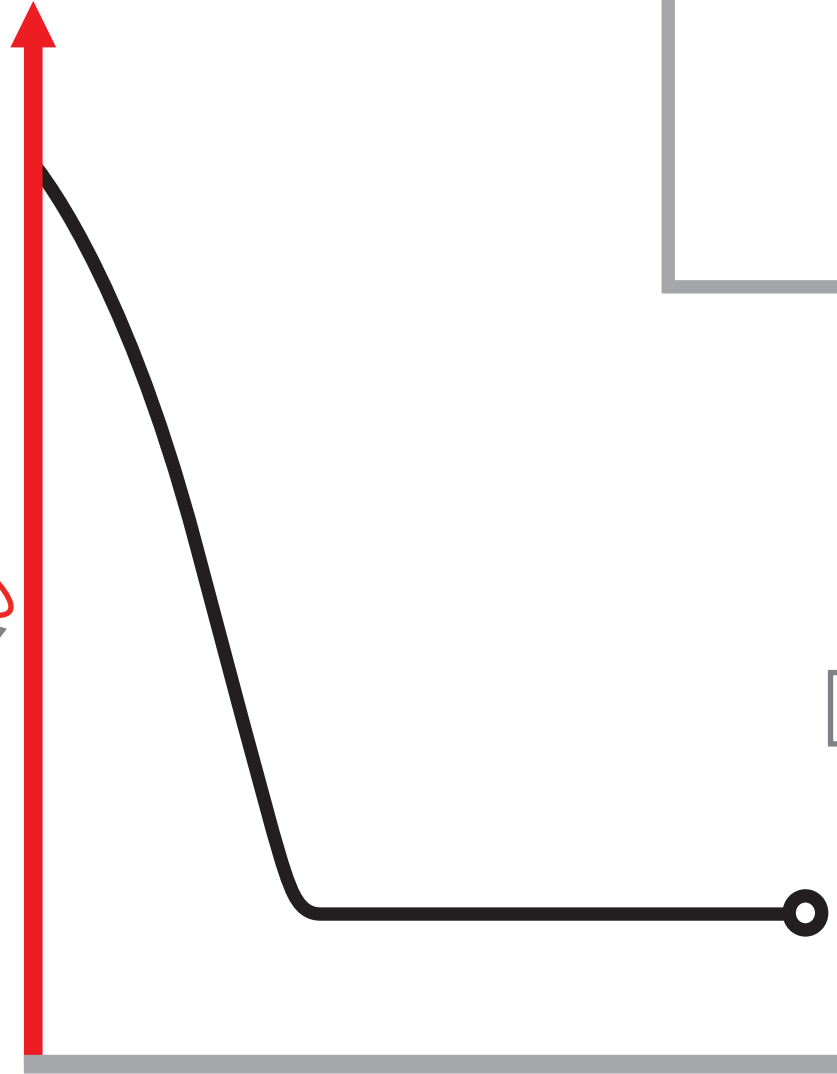


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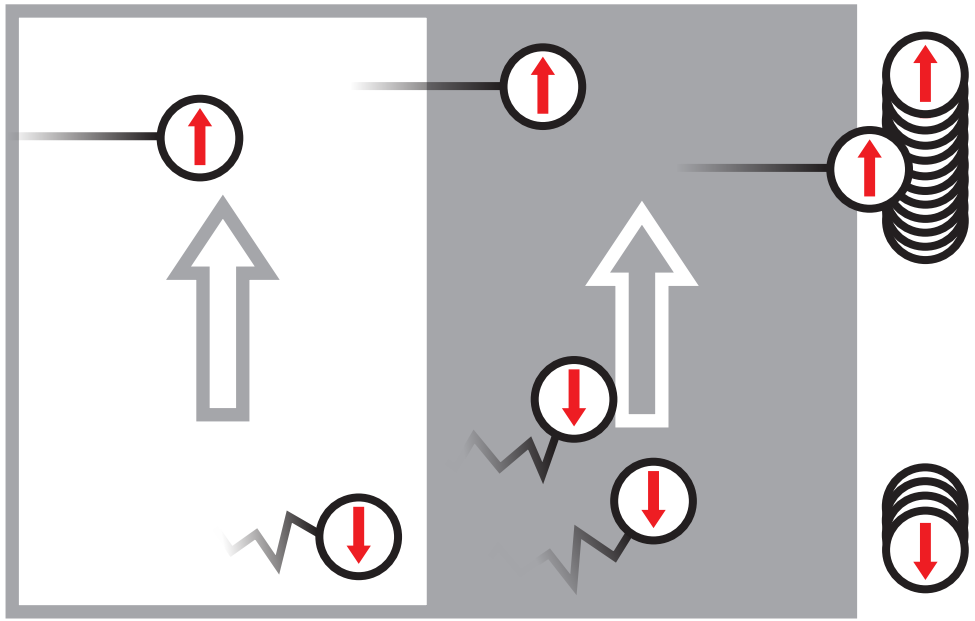
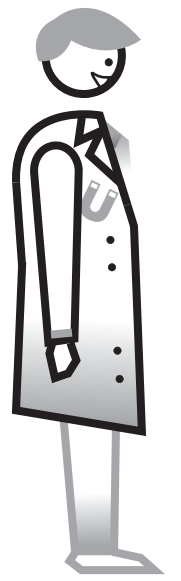


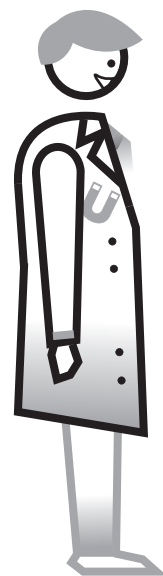
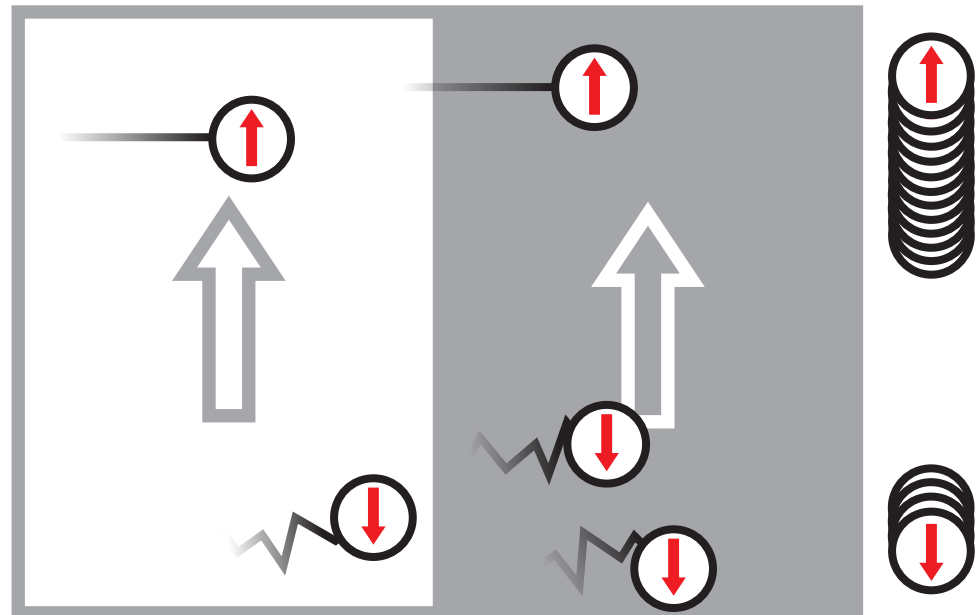
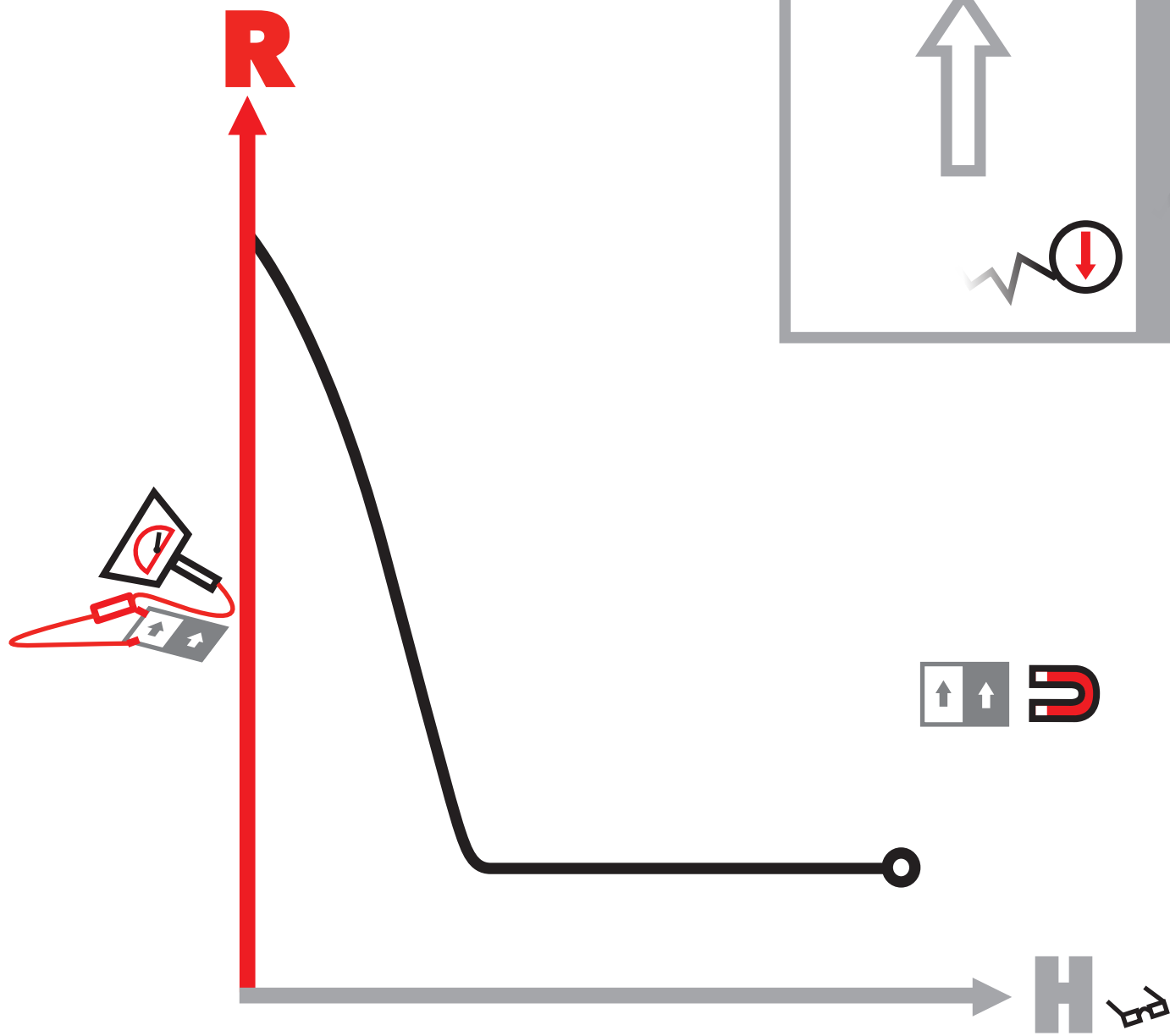


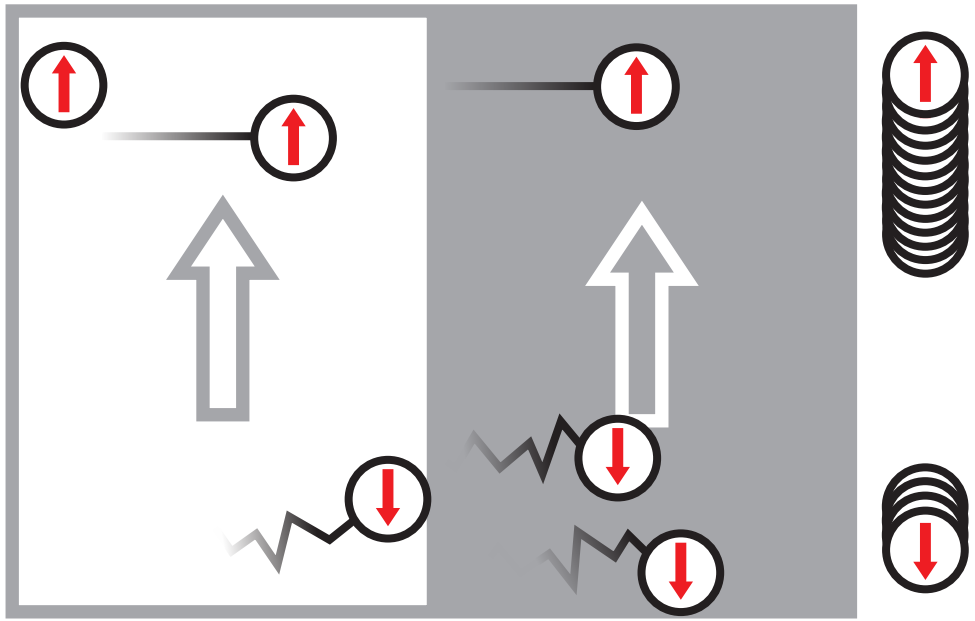
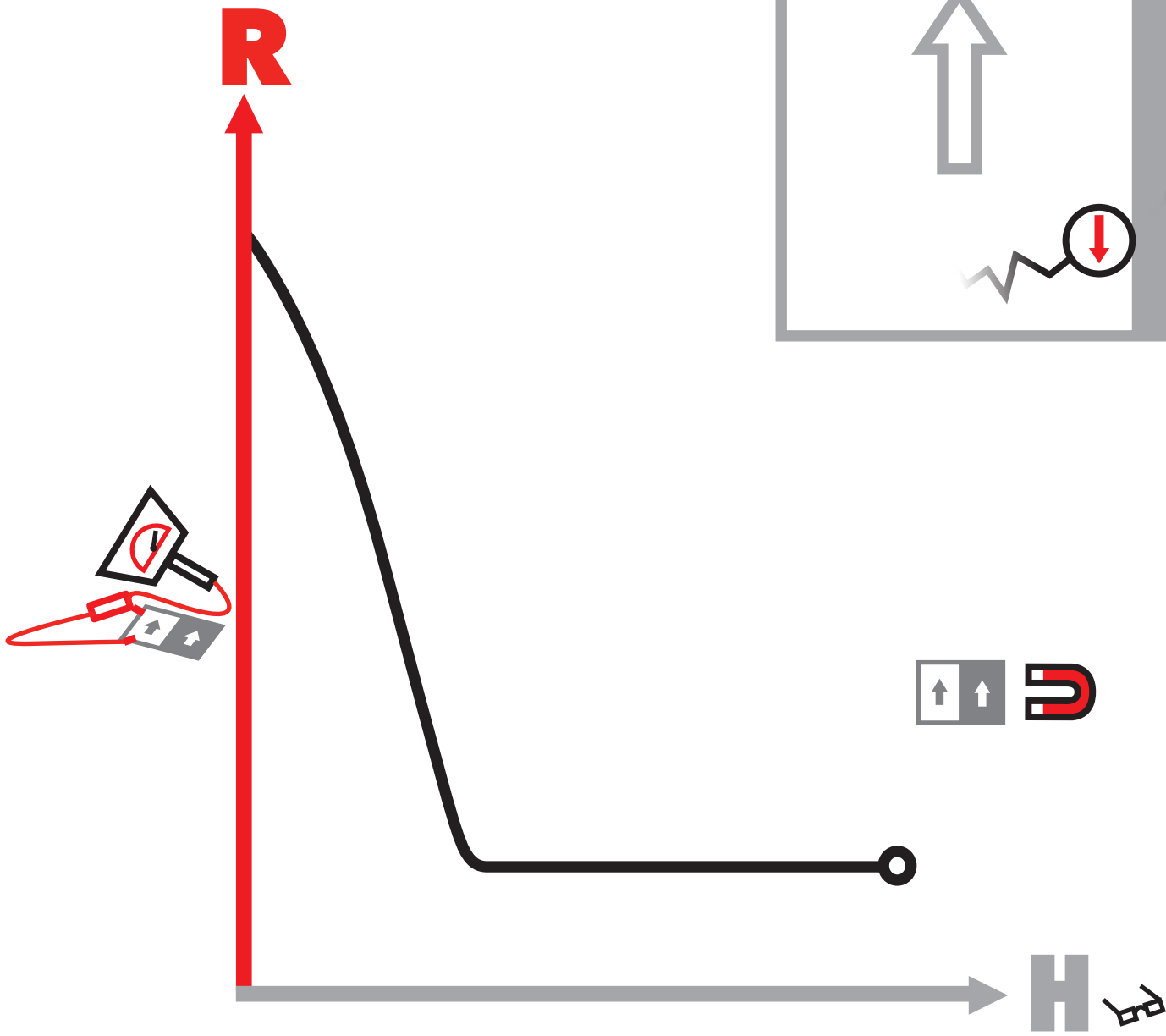
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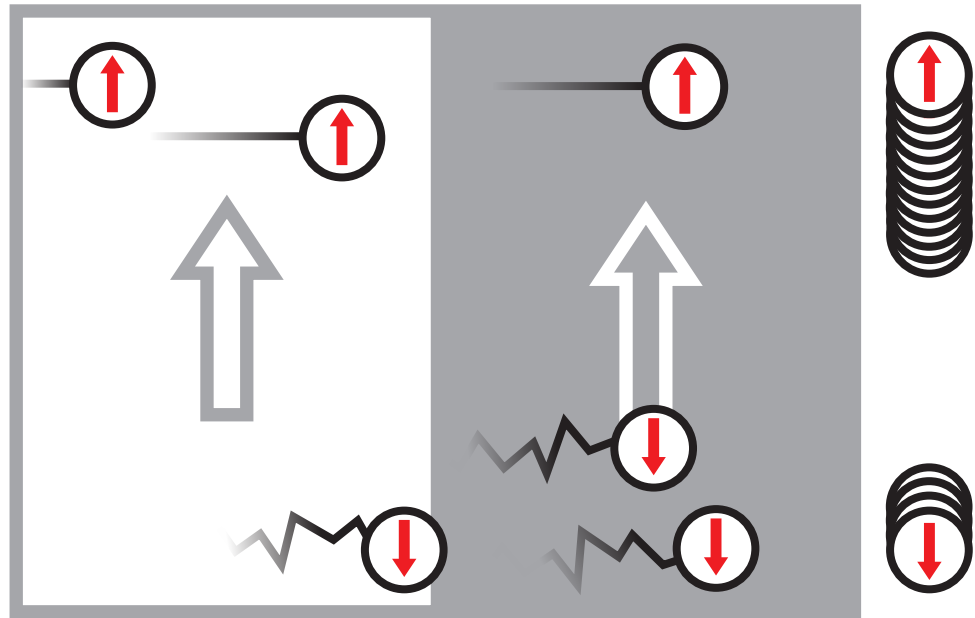
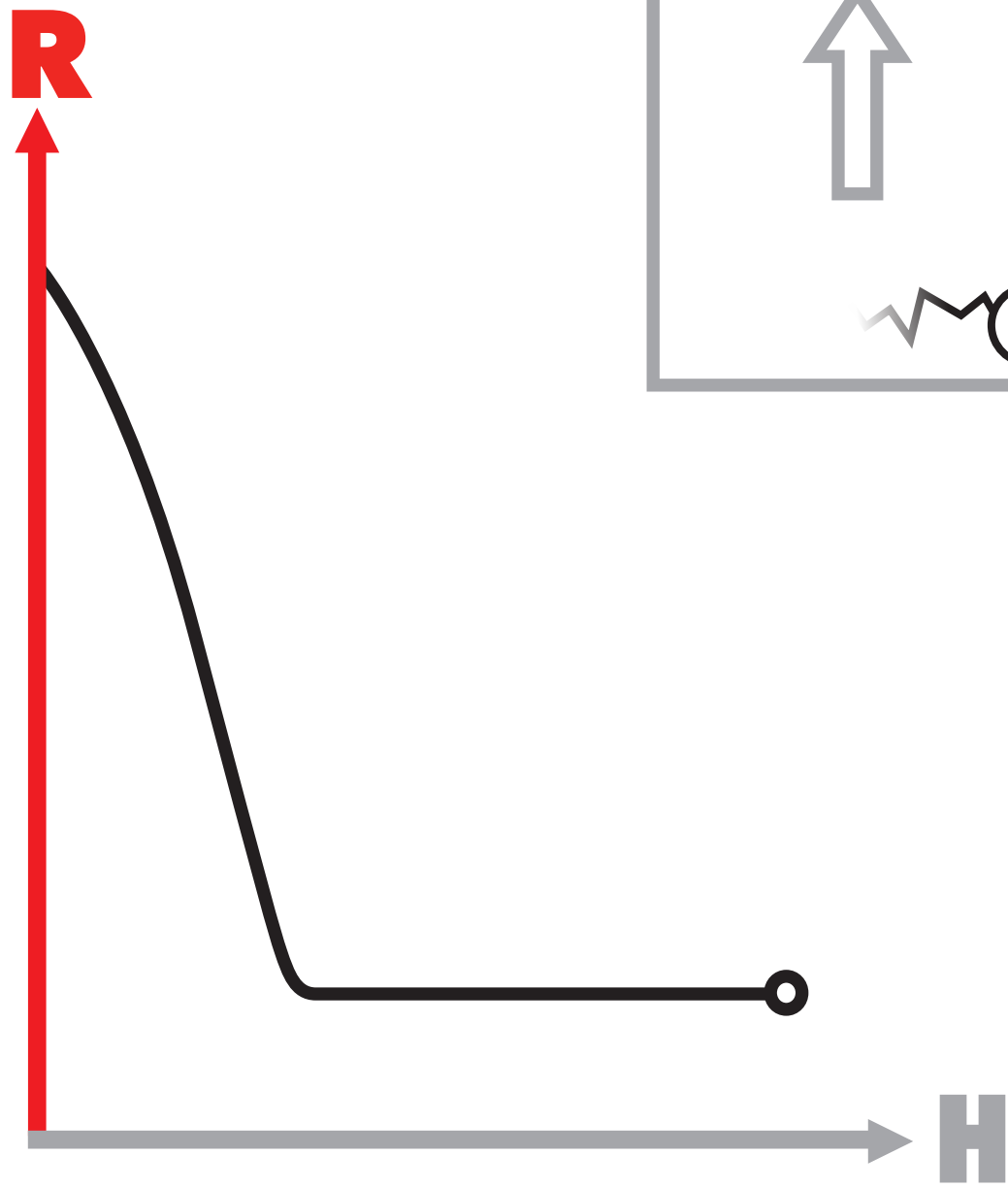


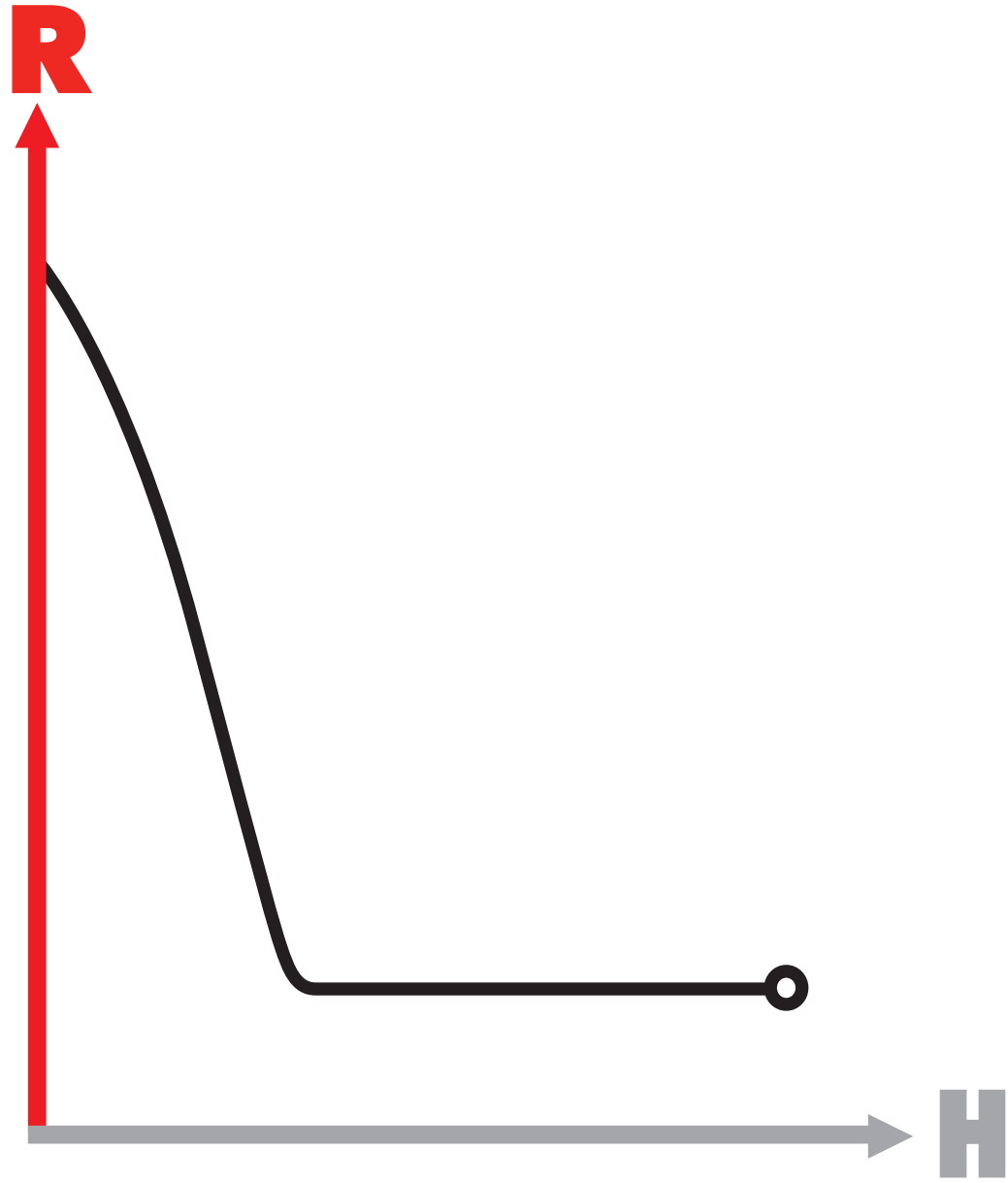
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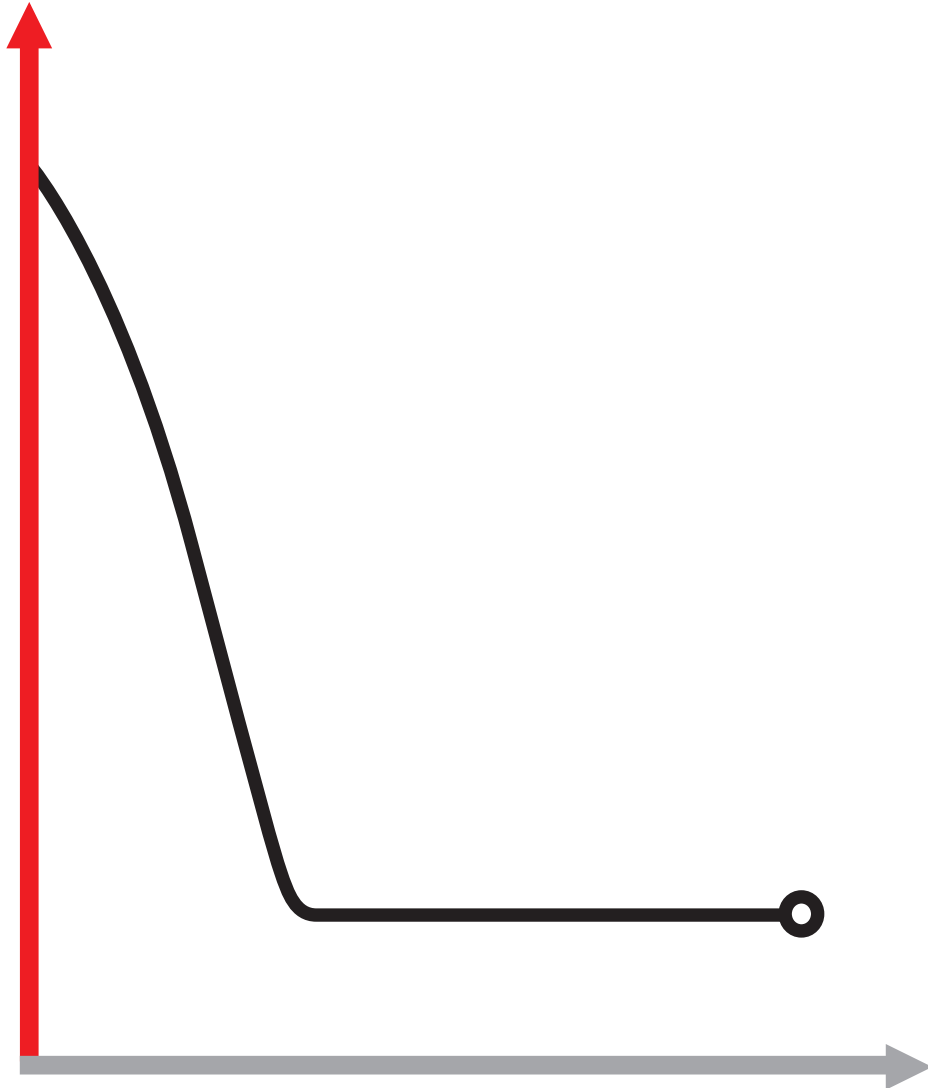




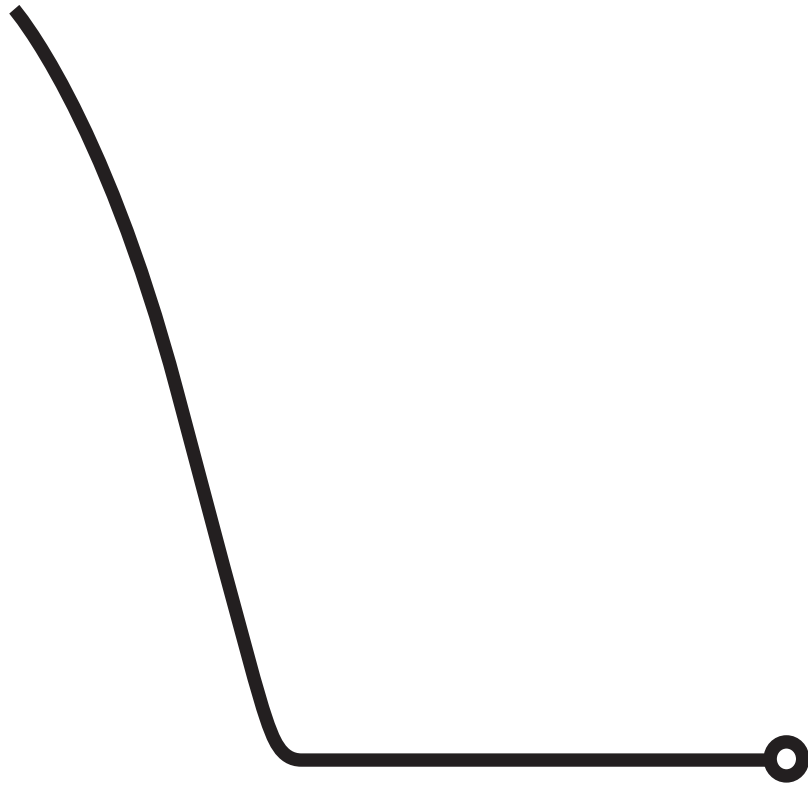


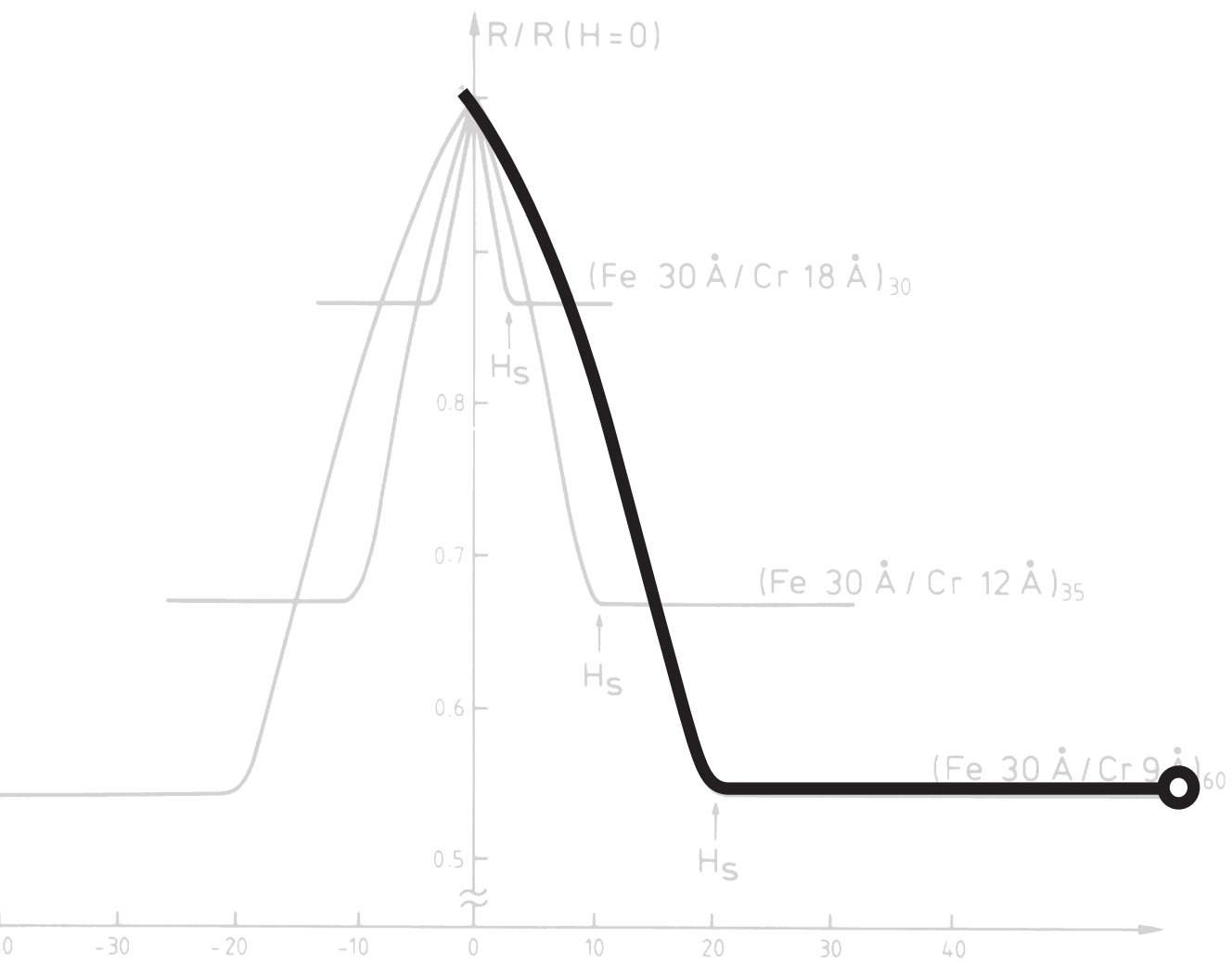


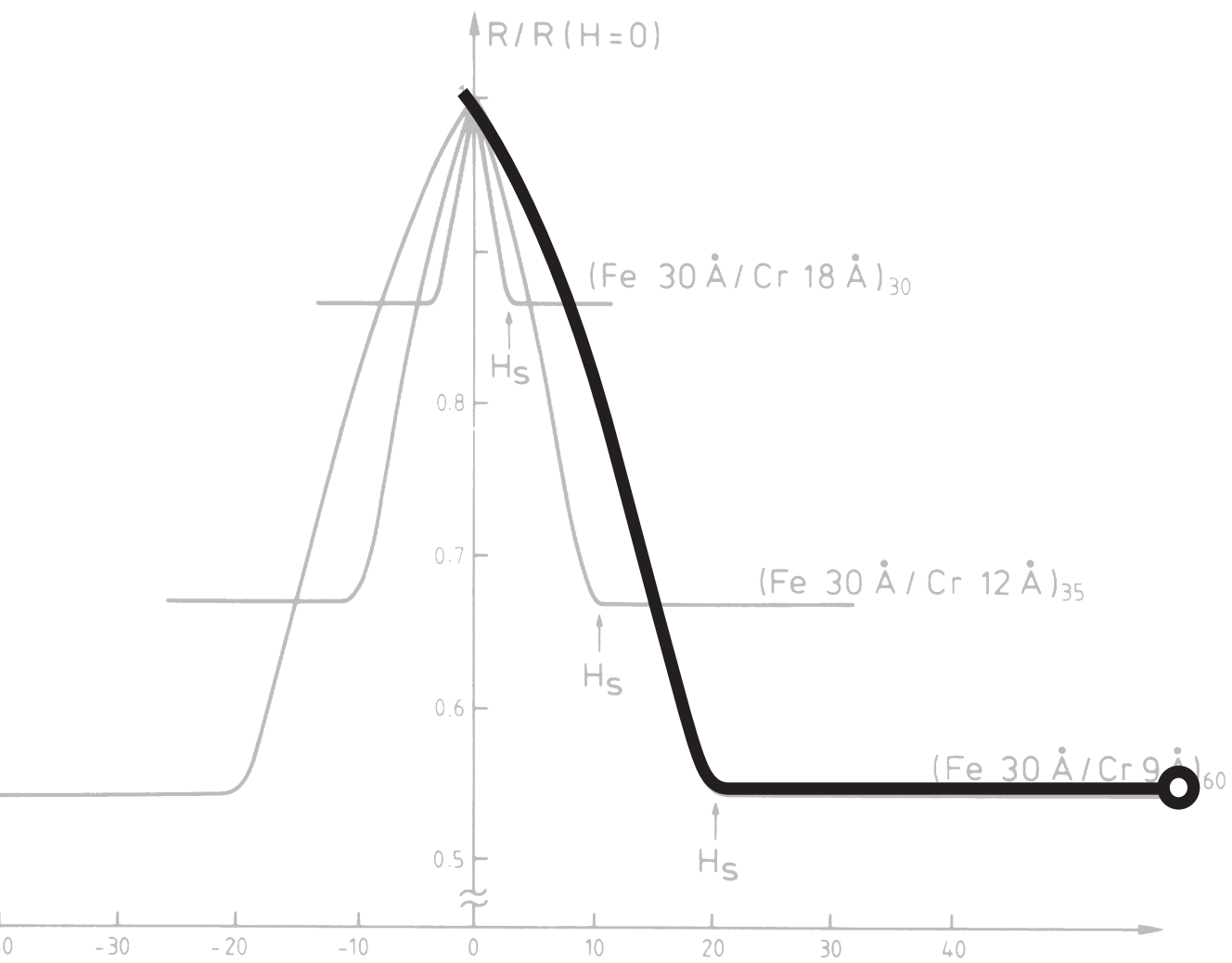


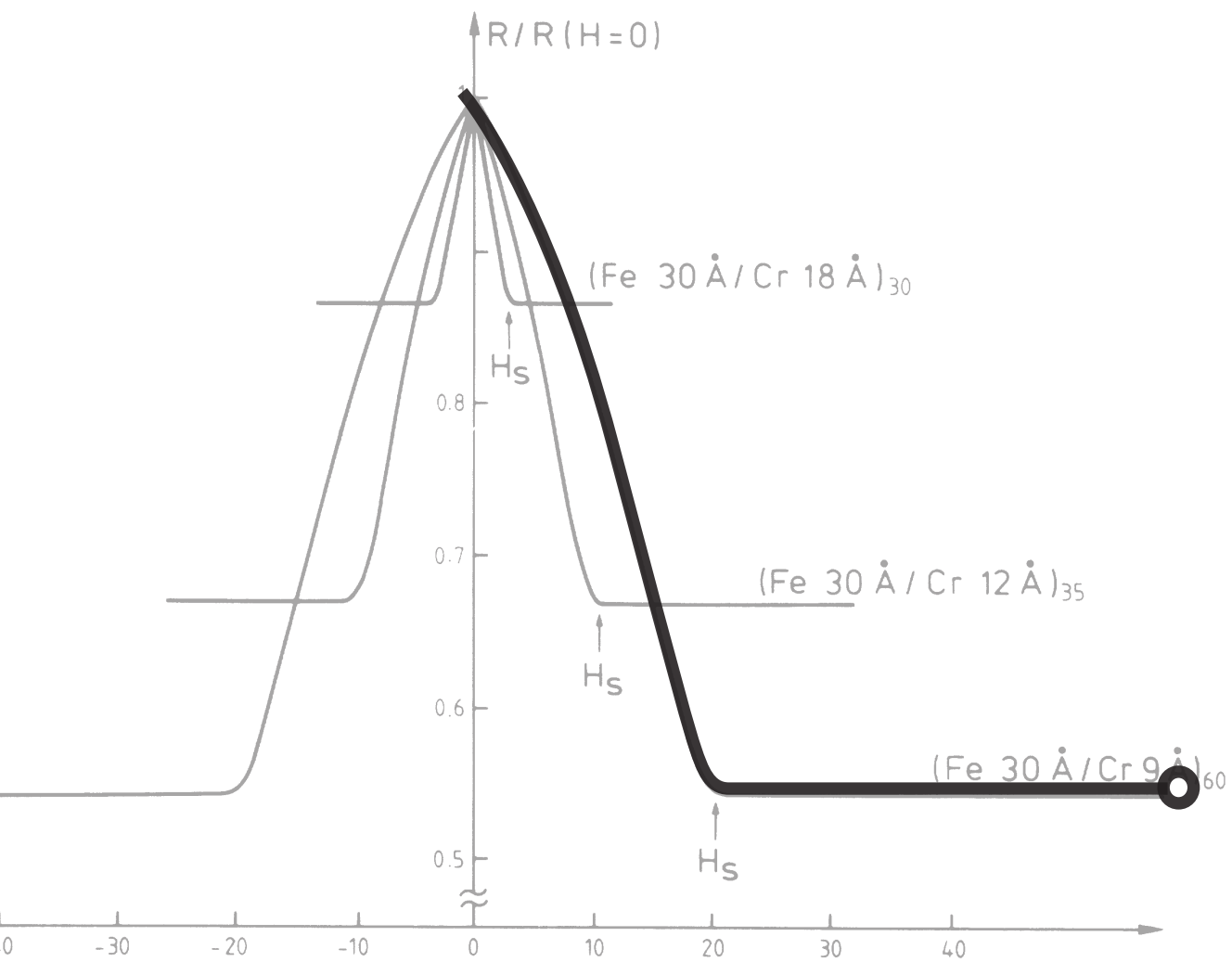


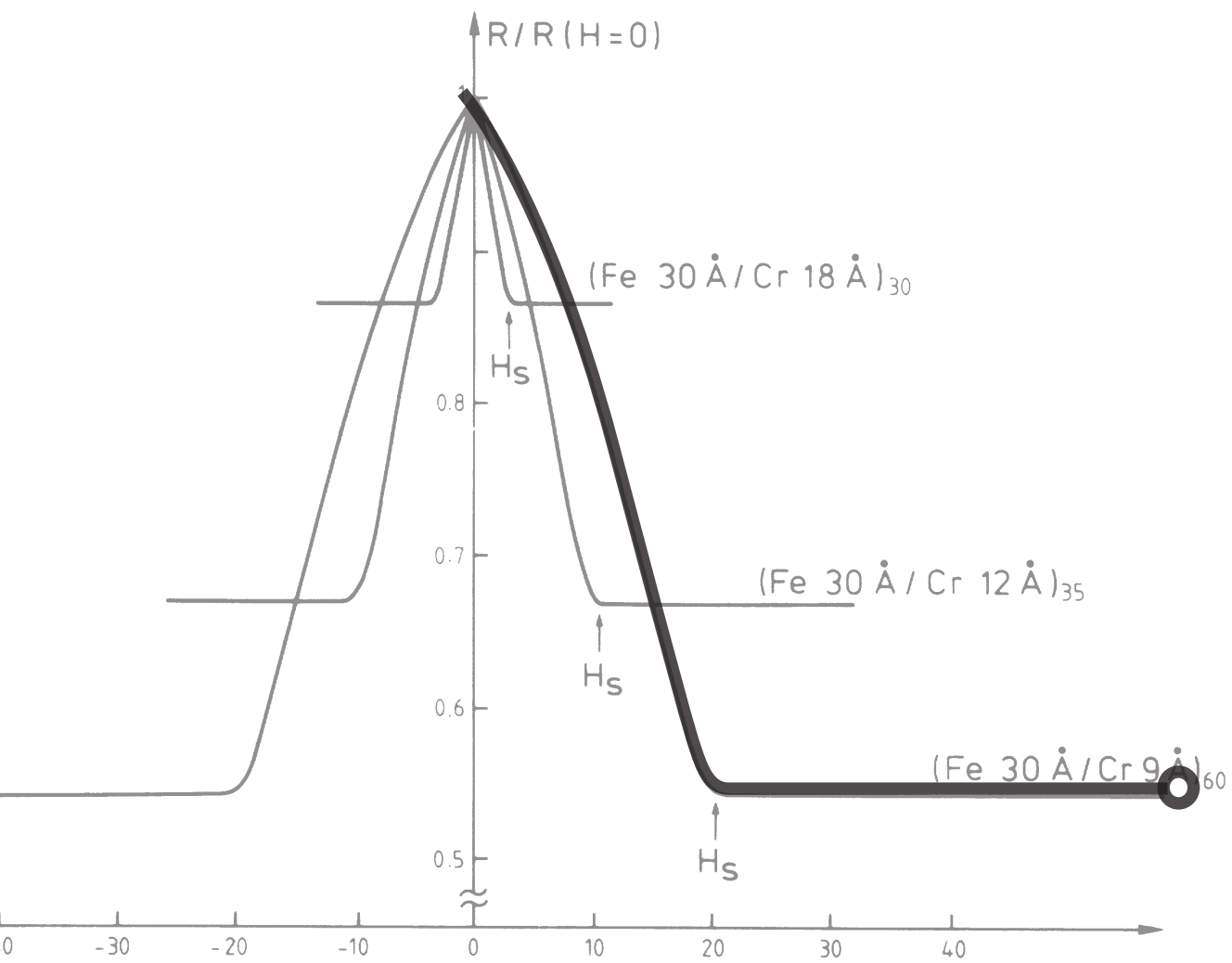


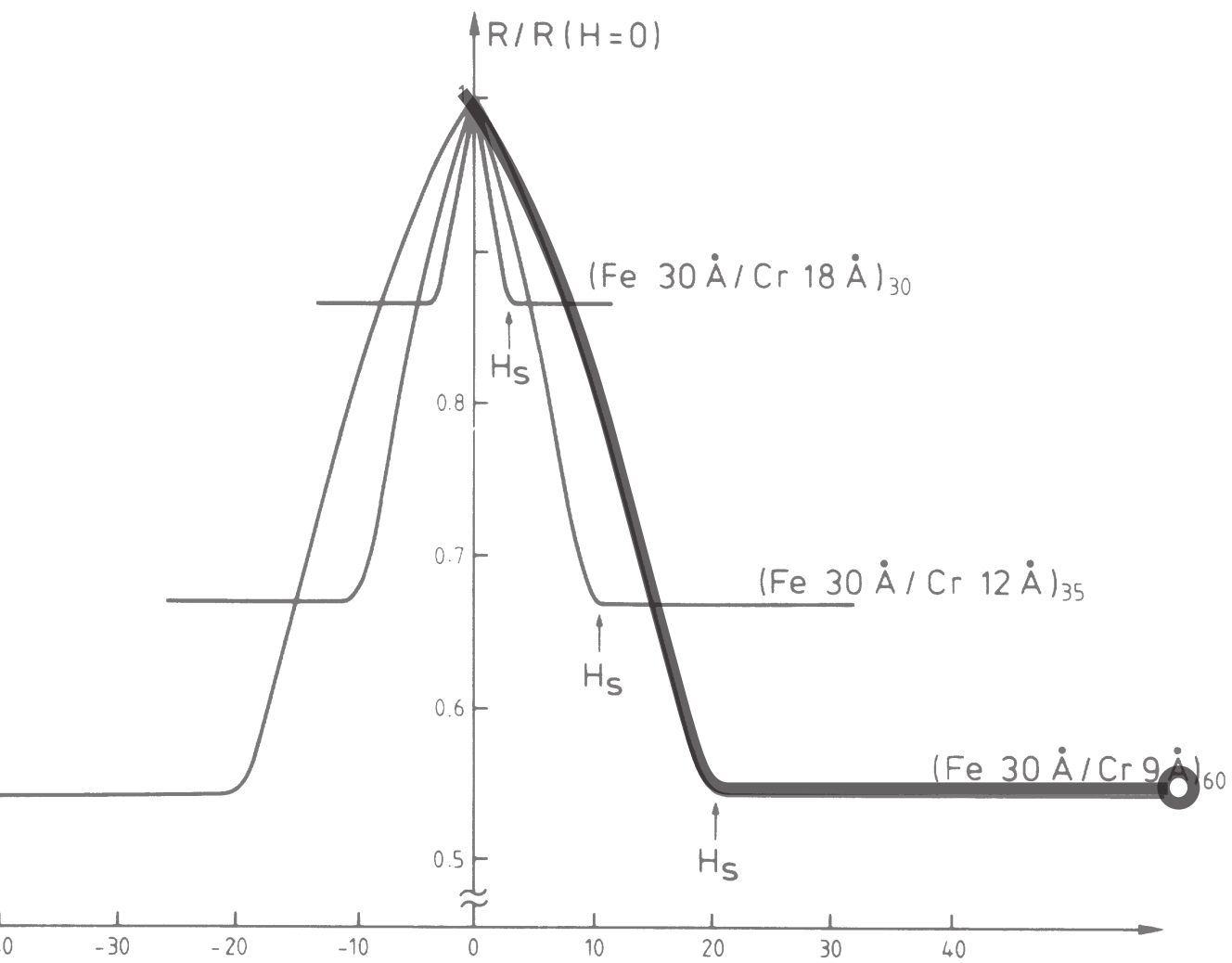


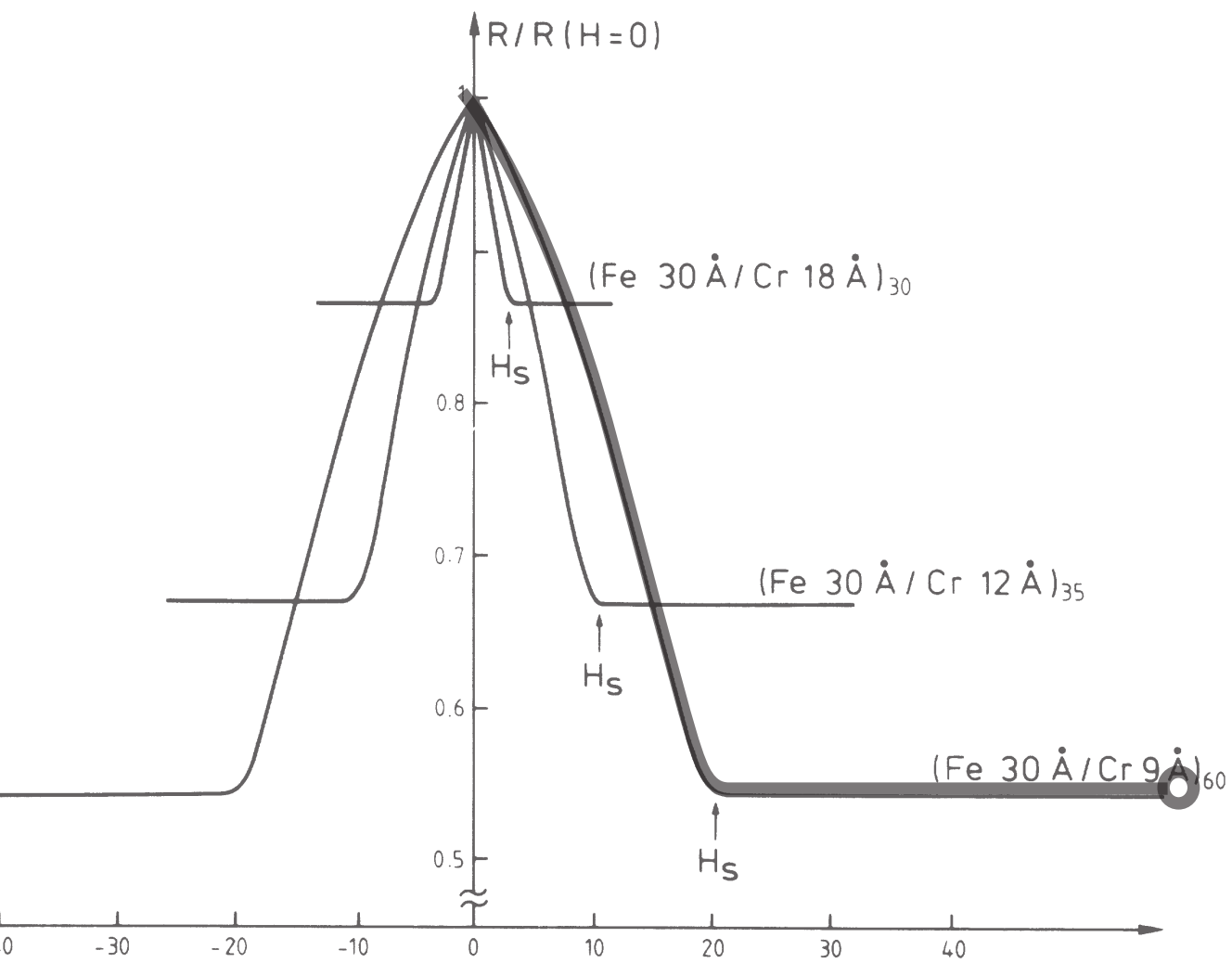


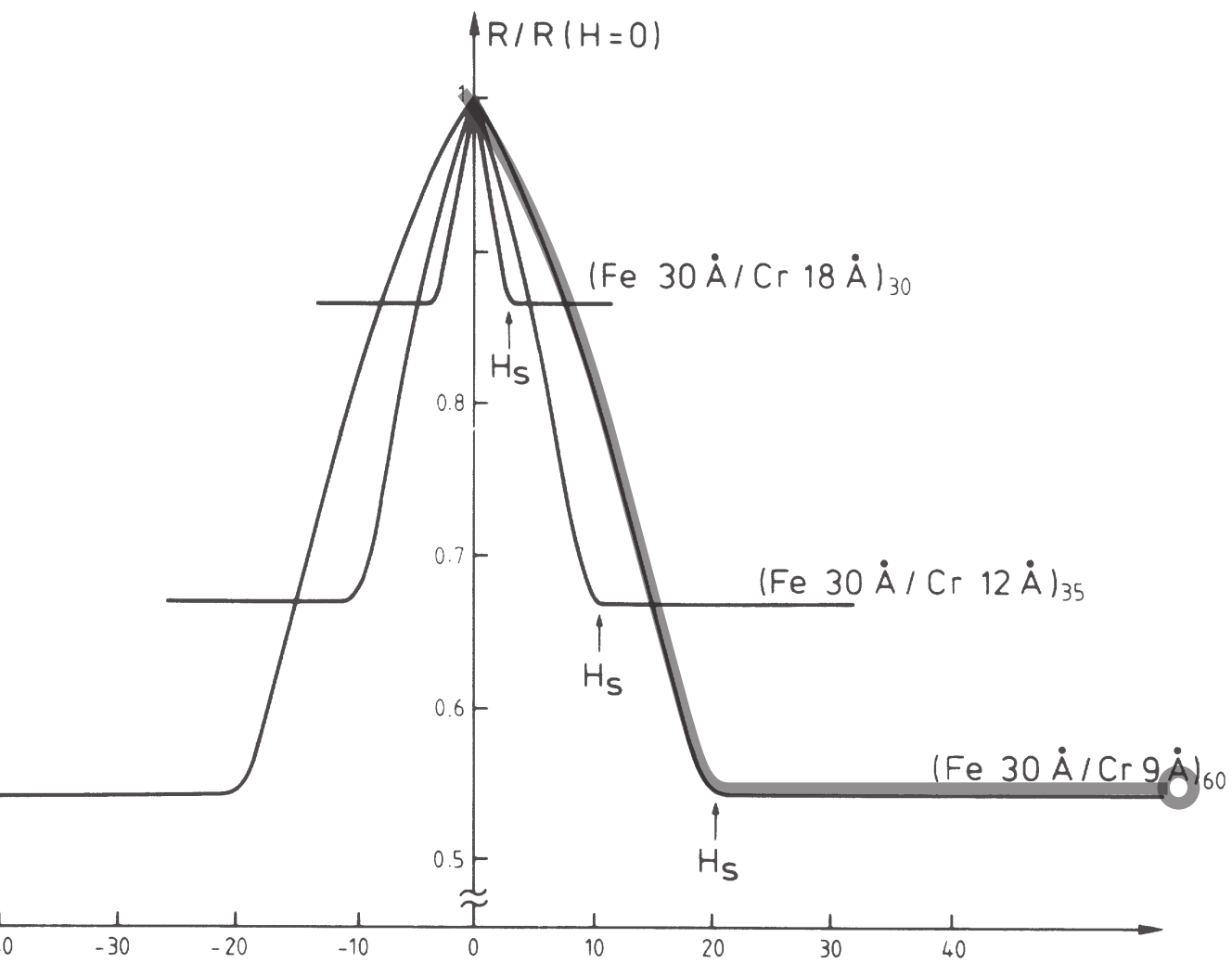




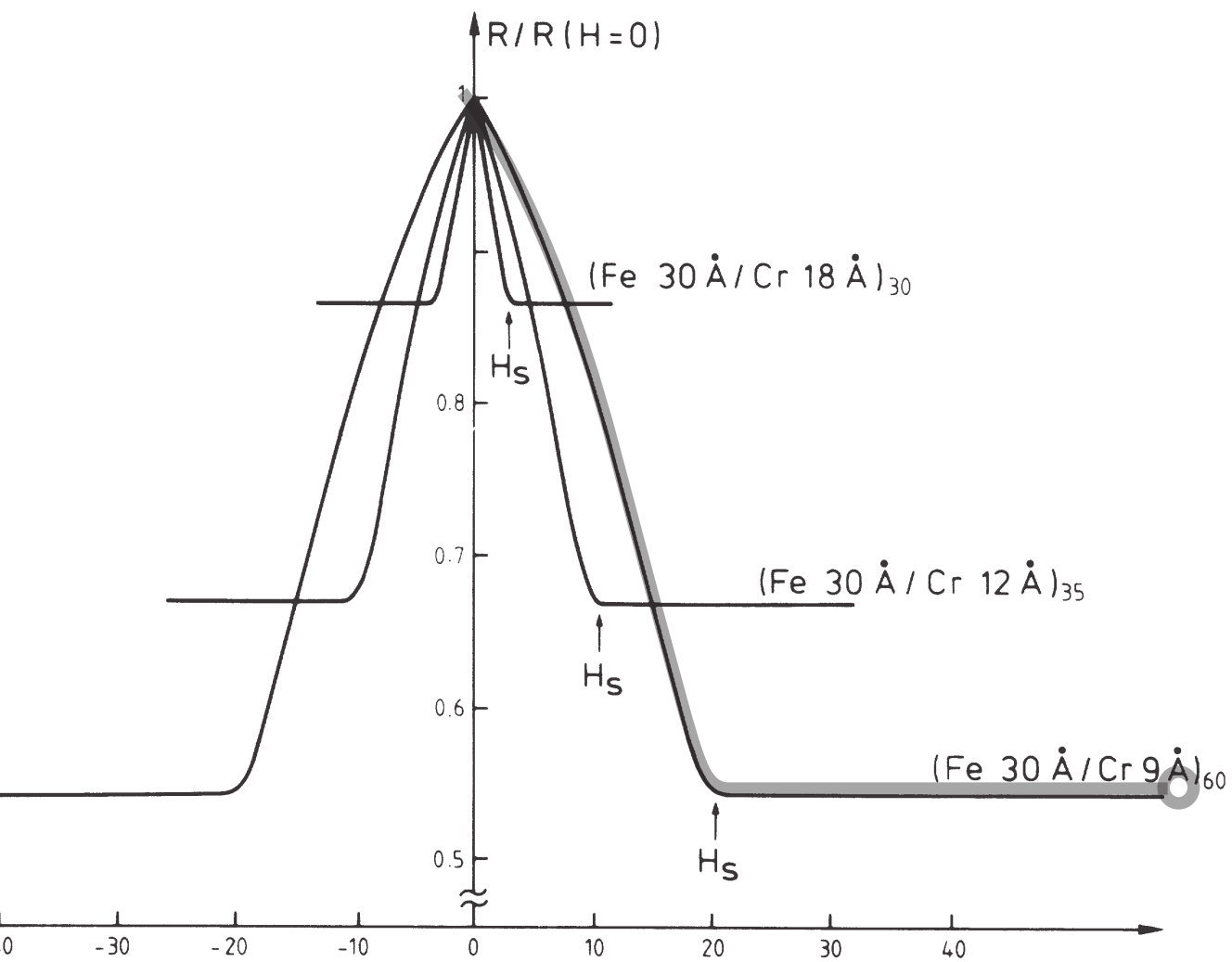


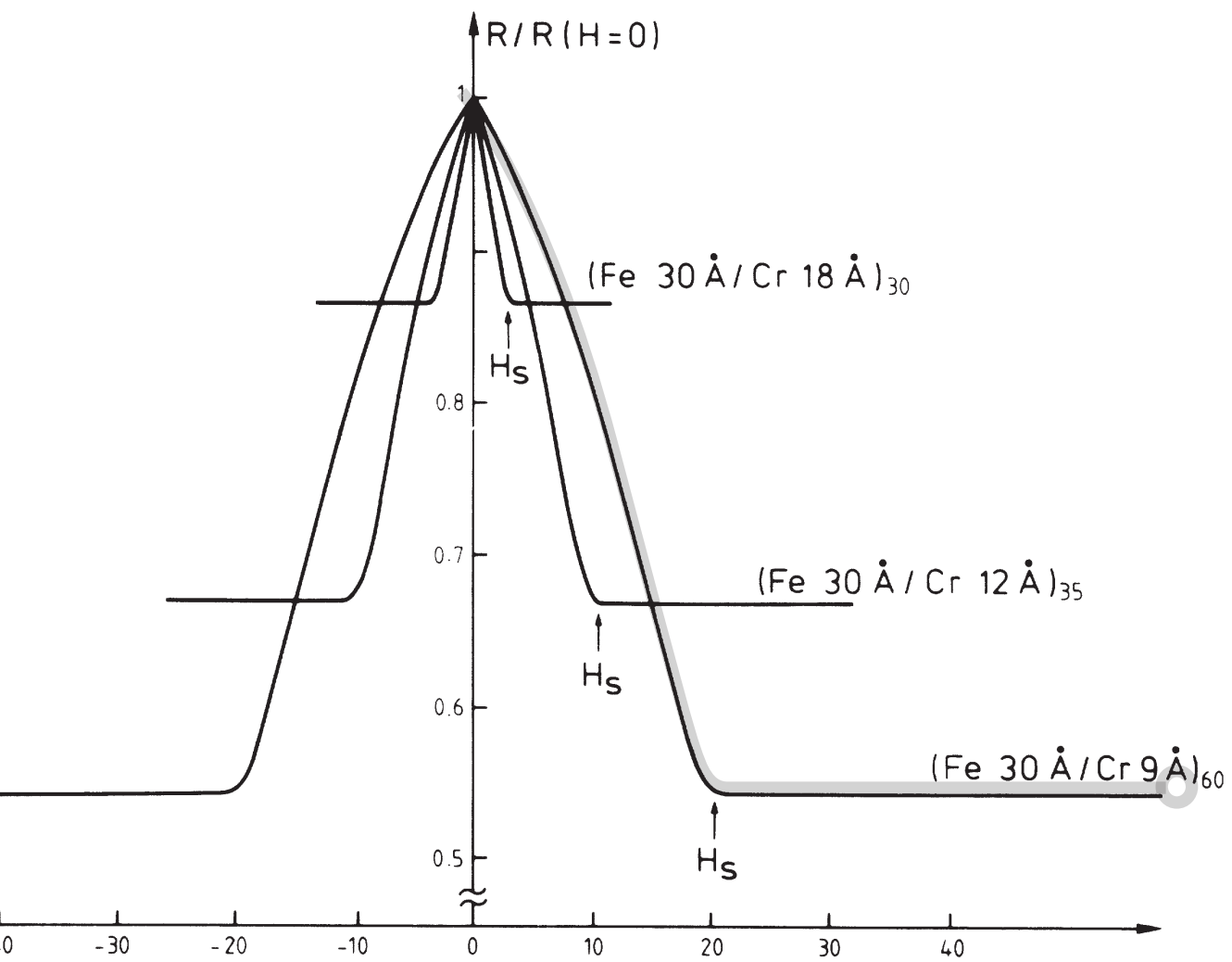


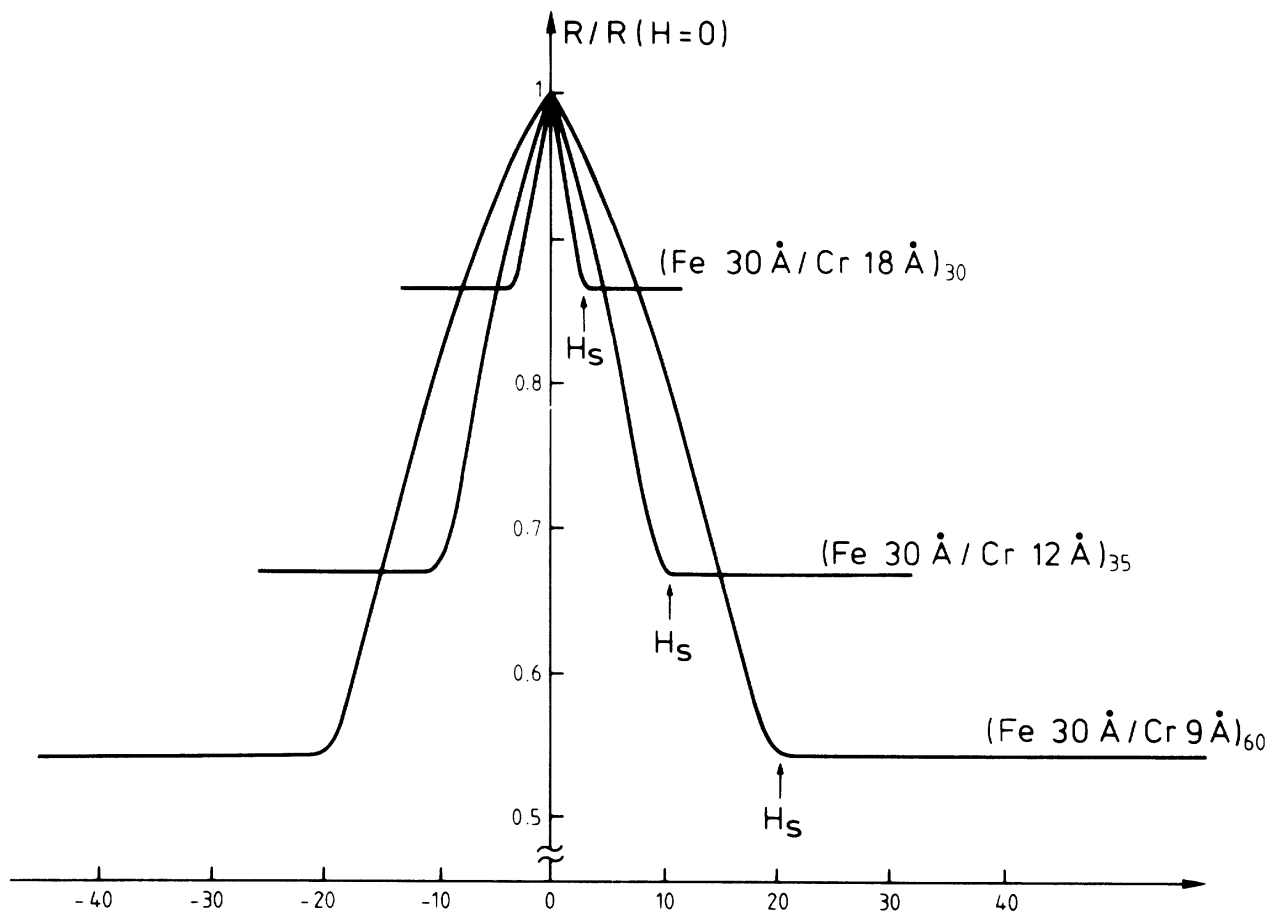


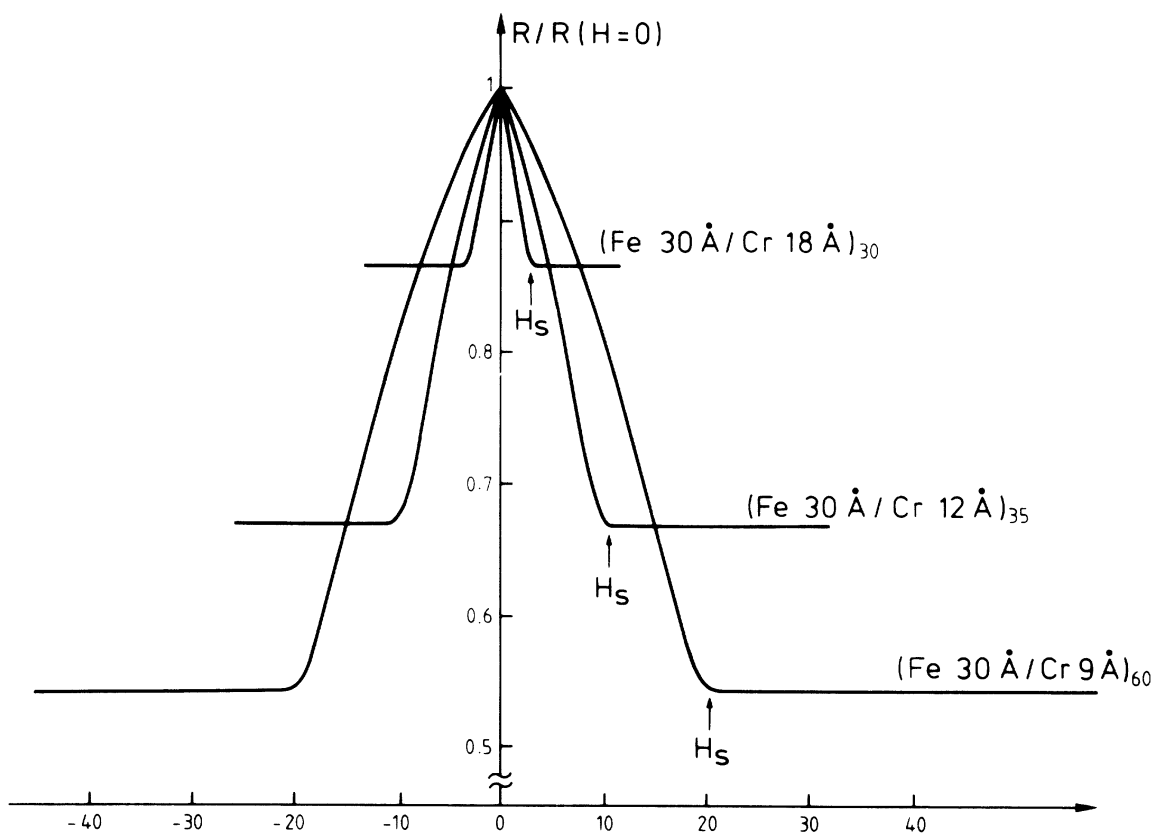


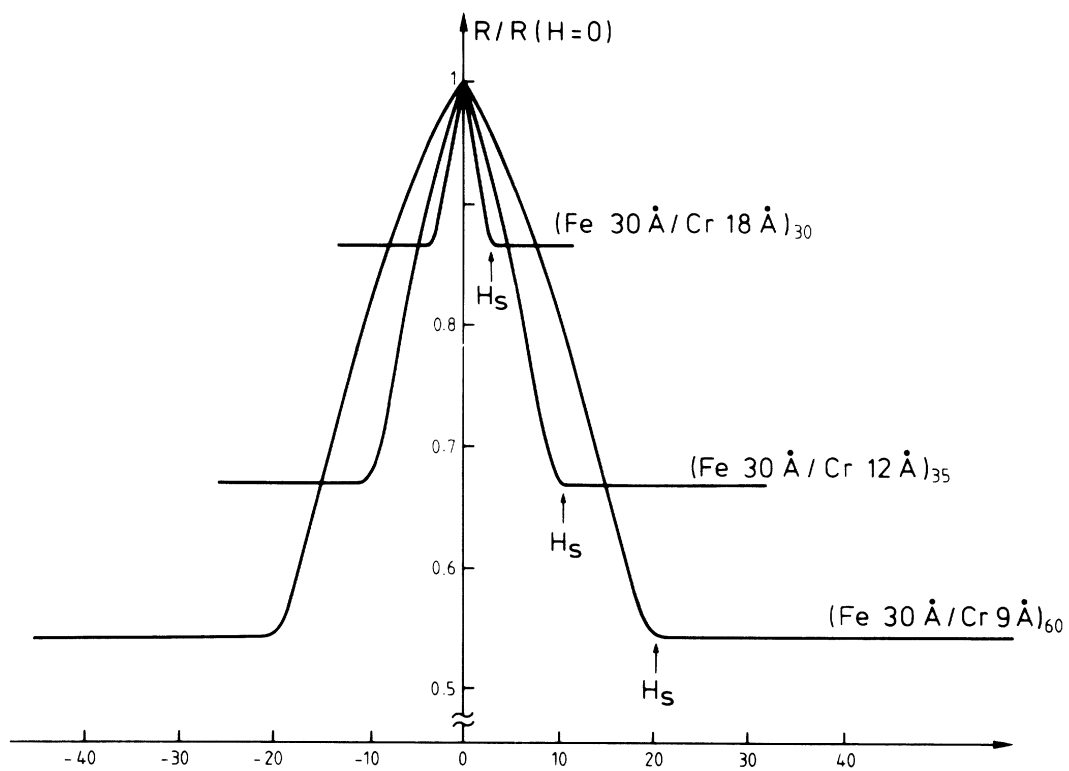


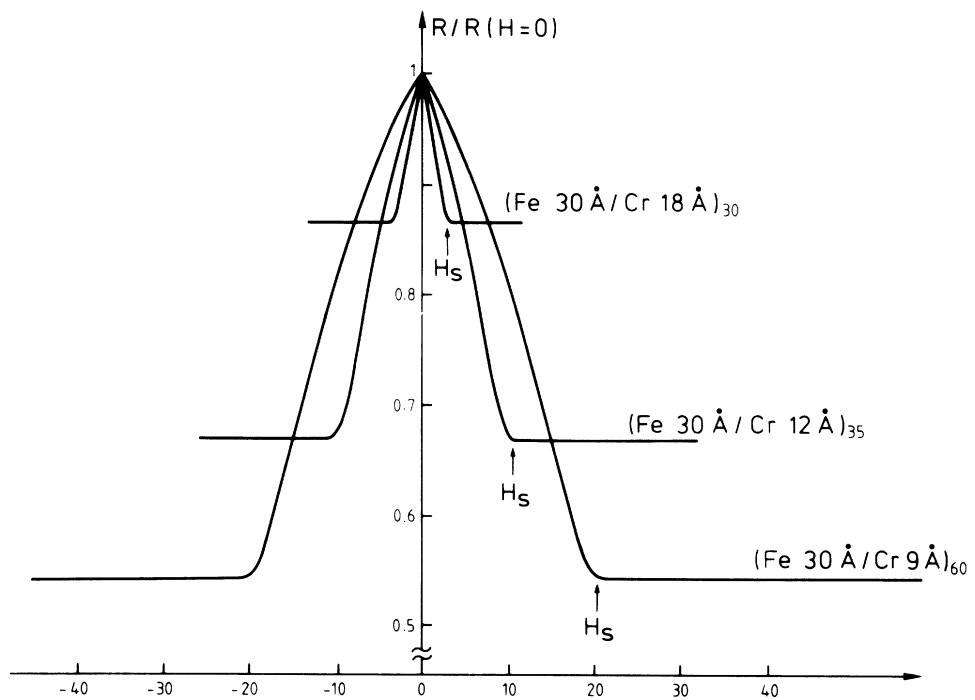


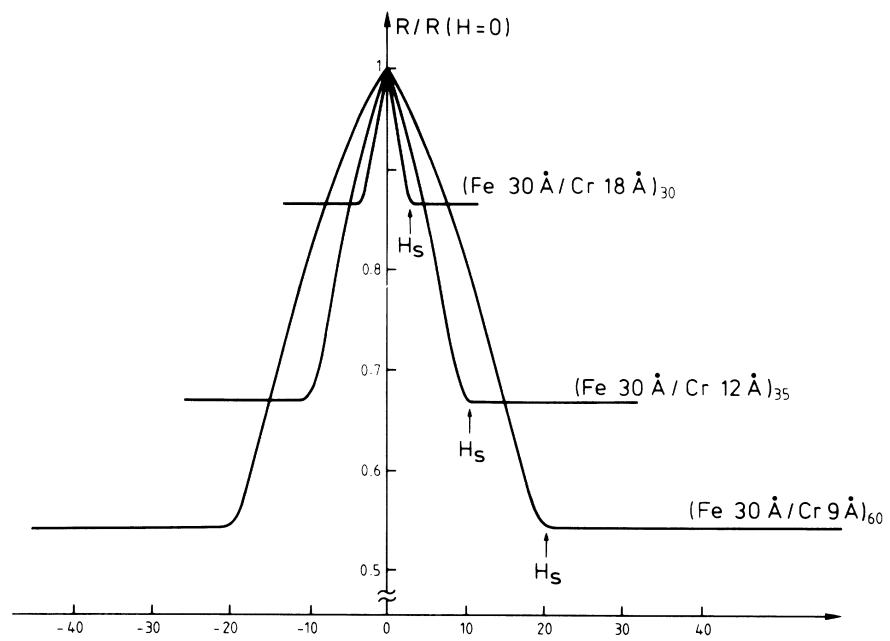


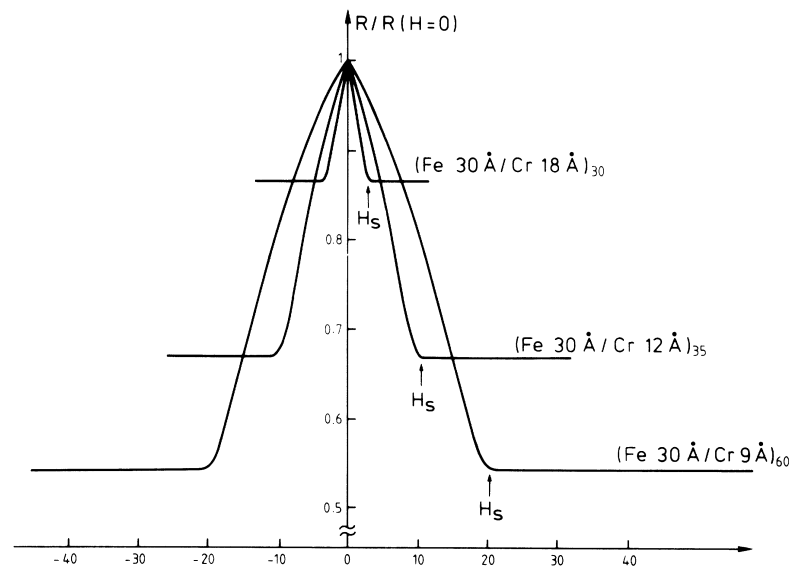














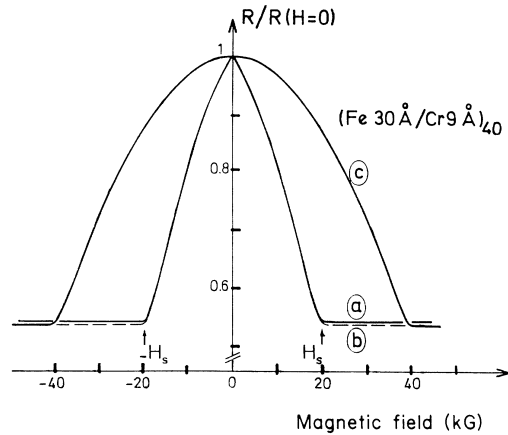


FIG. 2. Magnetoconductance of a  $[(\text{Fe } 30 \text{ \AA})/(\text{Cr } 9 \text{ \AA})]_{40}$  superlattice of 4.2 K. The current is along  $[110]$  and the field is in the layer plane along the current direction (curve *a*), in the layer plane perpendicular to the current (curve *b*), or perpendicular to the layer plane (curve *c*). The resistivity at zero field is  $54 \mu\Omega \text{ cm}$ . There is a small difference between the curves in increasing and decreasing field (hysteresis) that we have not represented in the figure. The superlattice is covered by a  $100\text{-\AA}$  Ag protection layer. This means that the magnetoconductance of the superlattice alone should be slightly higher.

of Grünberg *et al.*<sup>4</sup> and by the spin-polarized low-energy electron-diffraction experiments of Carbone and Alvarado.<sup>10</sup> The AF coupling between the Fe layers has been ascribed to indirect exchange interactions through the Cr layers, but a theoretical model of these interactions is still lacking.<sup>4,9</sup>

The magnetoconductance of the Fe/Cr superlattices has been studied by a classical ac technique on small rectangular samples. Examples of magnetoconductance curves at 4.2 K are shown in Figs. 2 and 3. The resistance decreases during the magnetization process and becomes practically constant when the magnetization is saturated. The curves *a* and *b* in Fig. 2 are obtained for applied fields in the plane of layers in the longitudinal and transverse directions, respectively. The field  $H_S$  is the field needed to overcome the AF couplings and to saturate the magnetization (compare with Fig. 1). In contrast, fields applied perpendicularly to the layers (curve *c*) have to overcome not only the AF coupling but also the magnetic anisotropy, so that the magnetoconductance is saturated at a field higher than  $H_S$ .

The most remarkable result exhibited in Figs. 2 and 3 is the huge value of the magnetoconductance. For  $t_{\text{Cr}}=9 \text{ \AA}$  and  $T=4.2 \text{ K}$ , see Fig. 2, there is almost a factor of 2 between the resistivities at zero field and in the saturated state, respectively (in absolute value, the resistivity change is about  $23 \mu\Omega \text{ cm}$ ). By comparison of the results for three different samples in Fig. 3, it can be seen

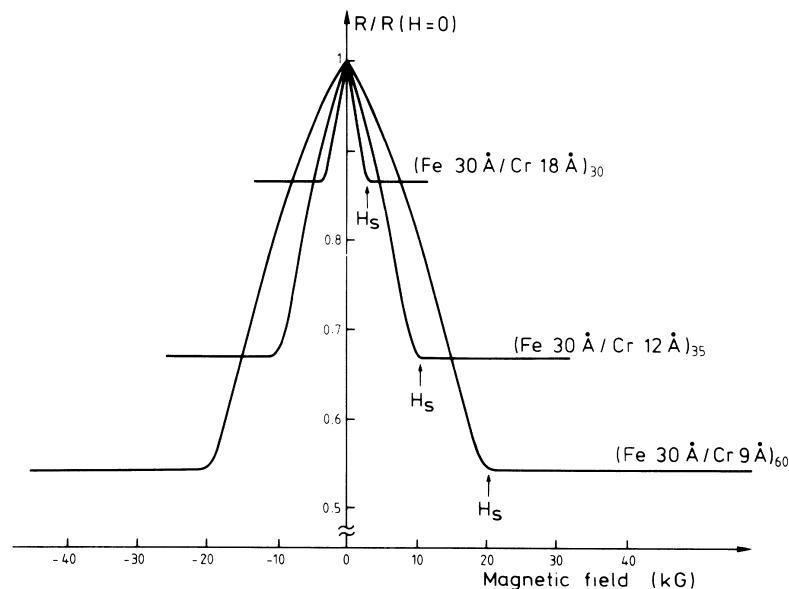
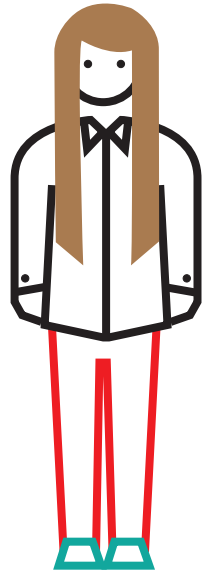


FIG. 3. Magnetoconductance of three Fe/Cr superlattices at 4.2 K. The current and the applied field are along the same  $[110]$  axis in the plane of the layers.

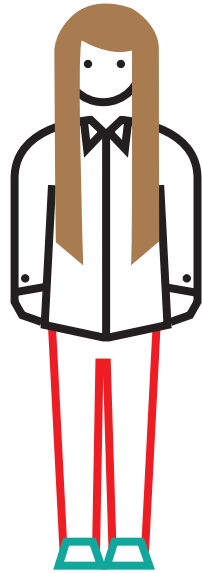


**GRAPHISM**

# GRAPHISM

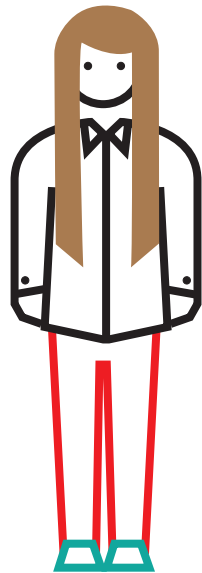


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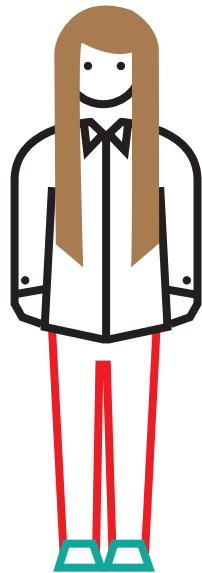
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**GRAPHISM**



**Chloé  
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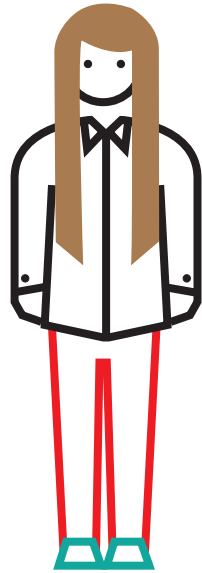
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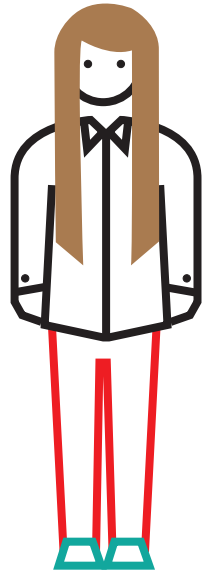


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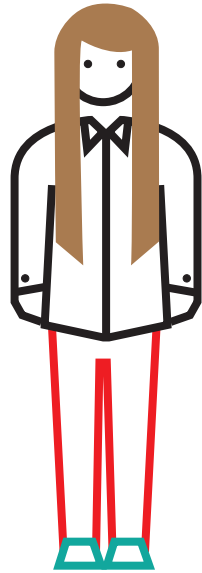
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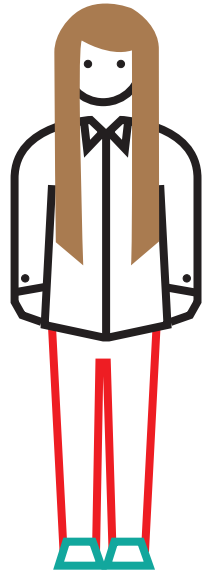
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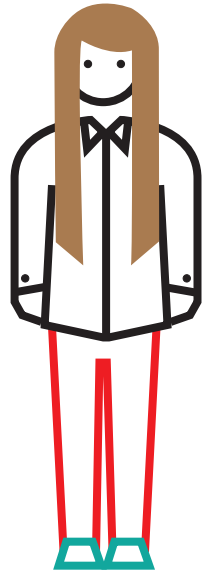
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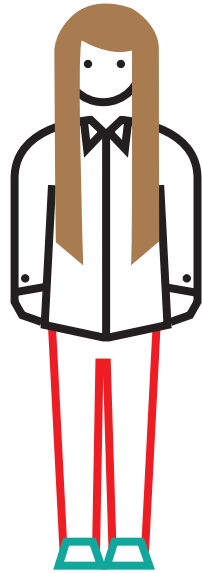
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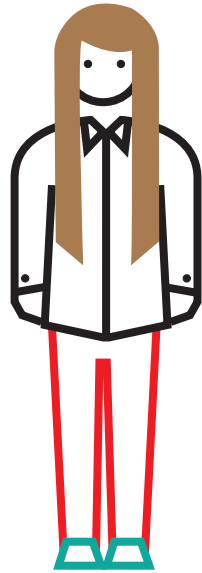
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**Chloé  
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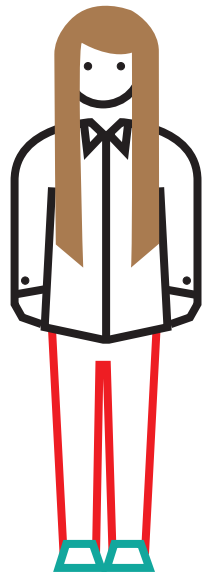


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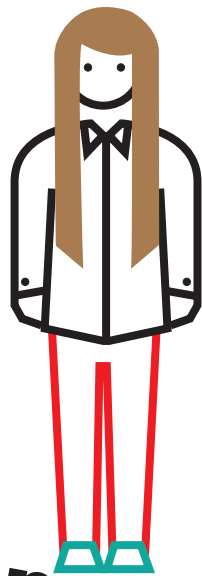
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**DSAA DIS  
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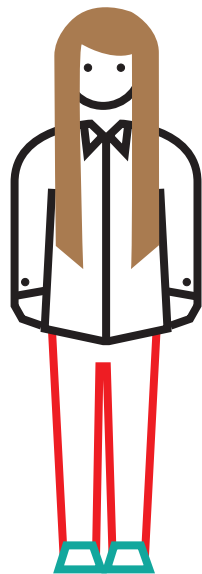
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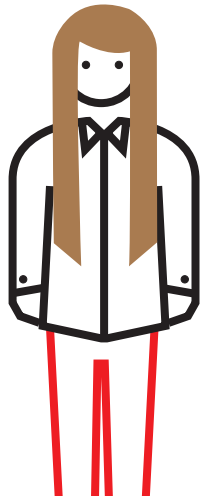
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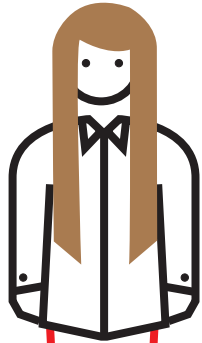
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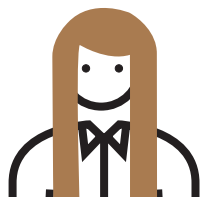
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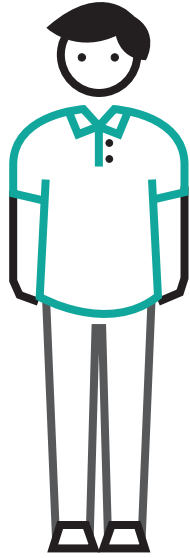
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**PHYSICS**

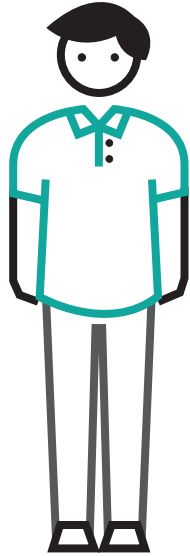
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# PHYSICS REIMAGINED



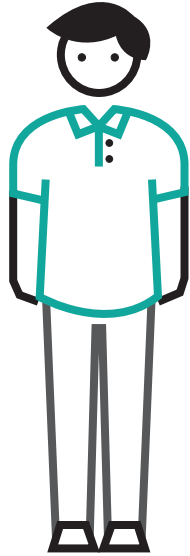


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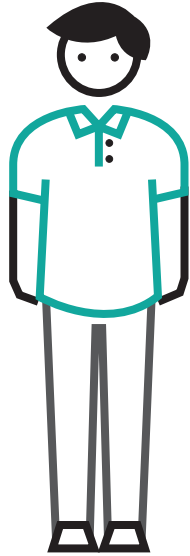
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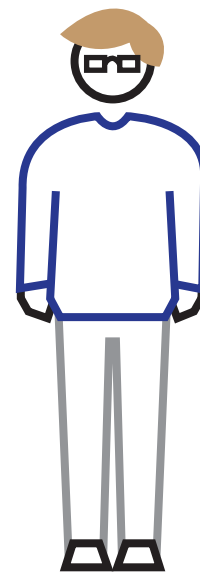
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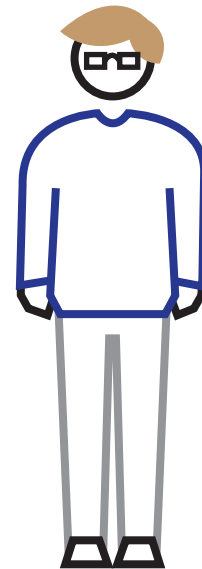
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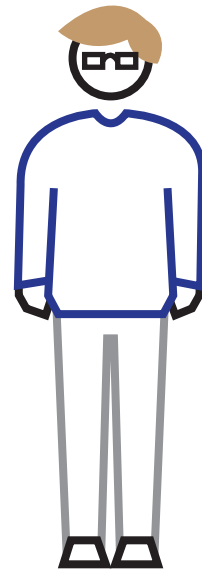
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Frederic



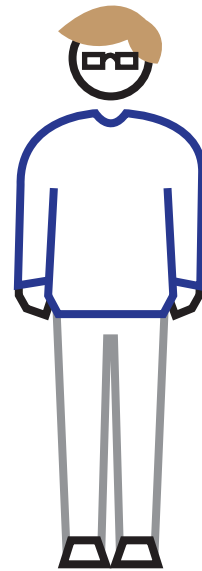
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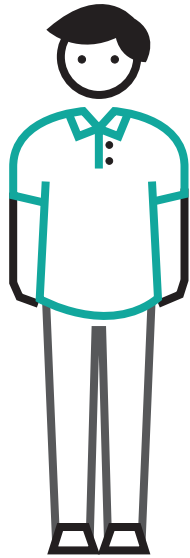


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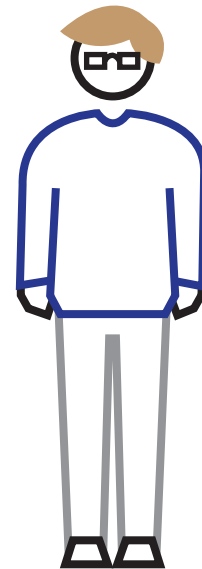
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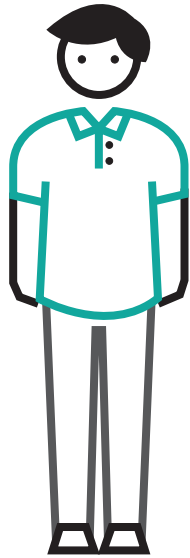
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**Frederic BOUQUET**



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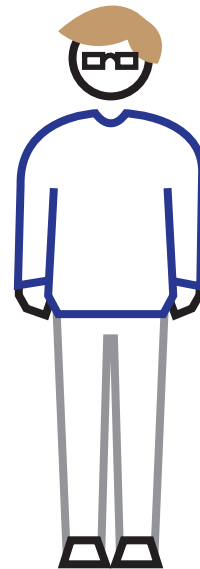


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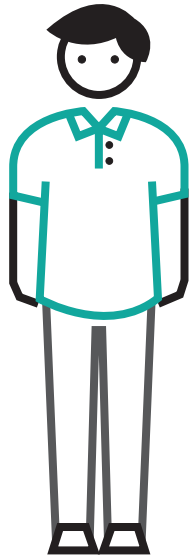


**Julien BOBROFF**  
**Frederic BOUQUET**

**La Physique**

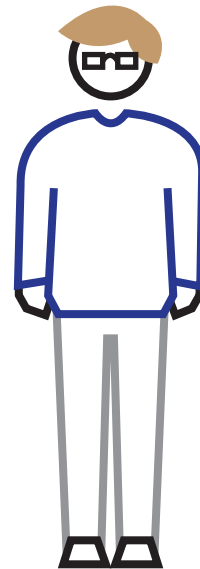


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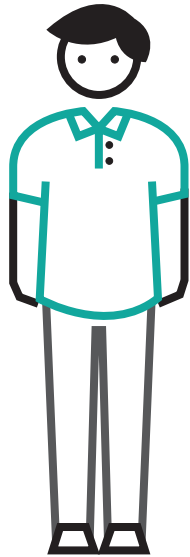


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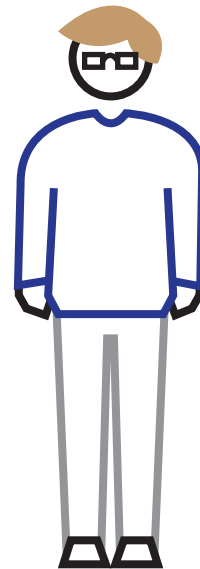
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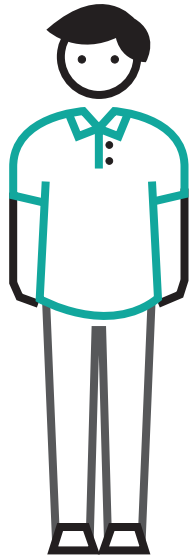


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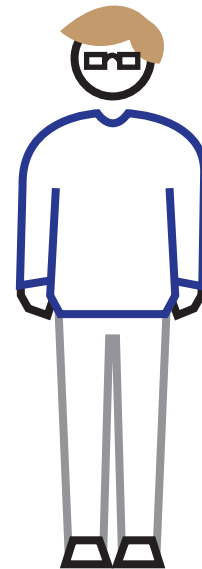


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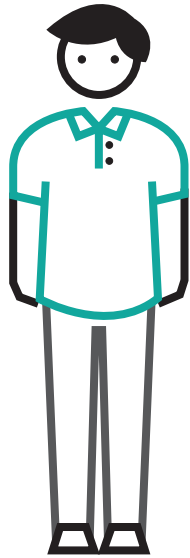


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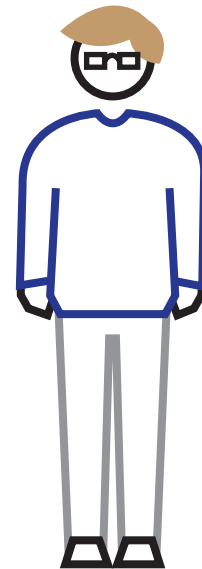
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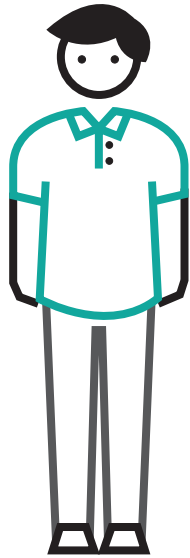


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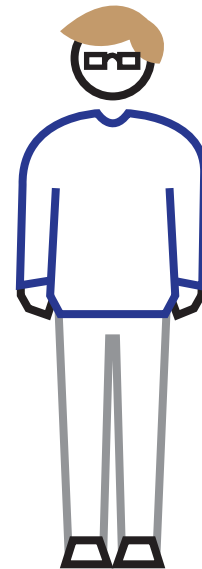
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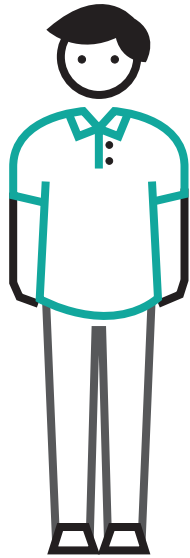


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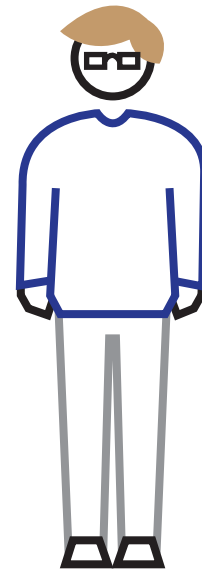


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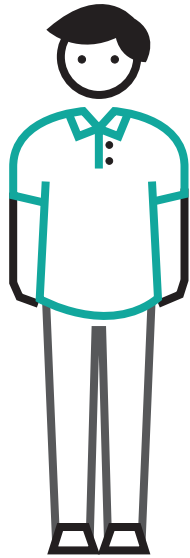


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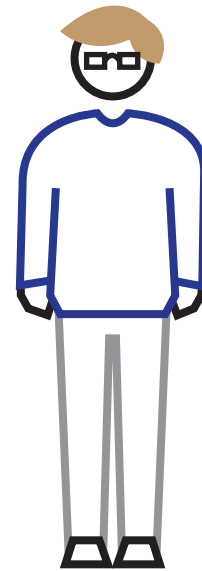


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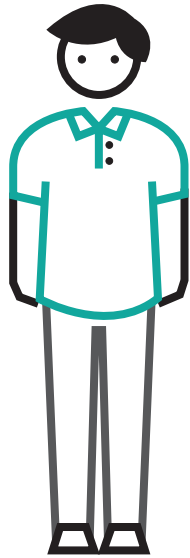
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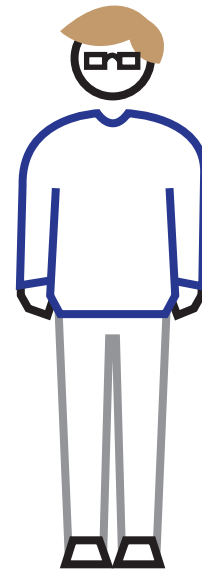
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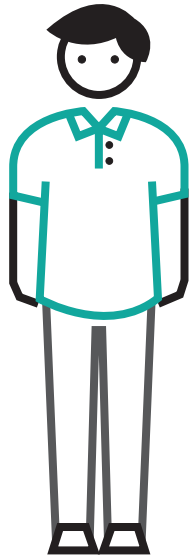


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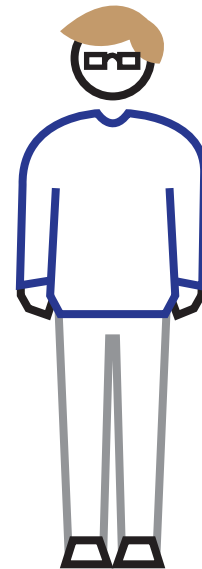


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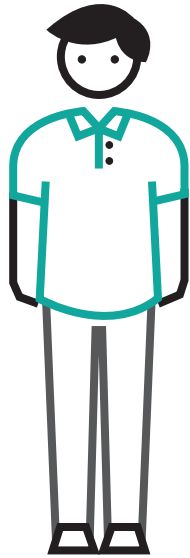


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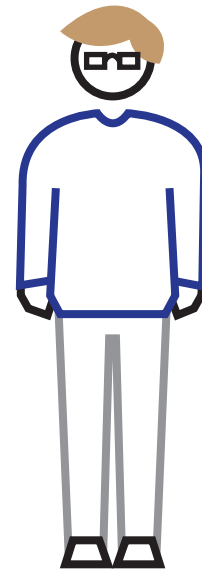


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**Université Paris-**

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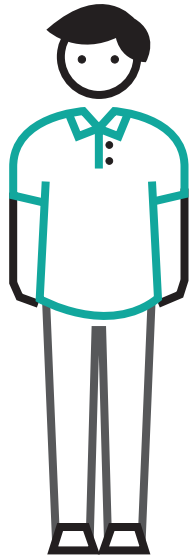


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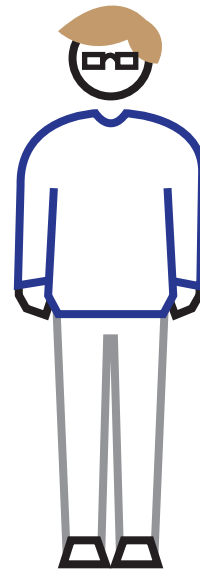


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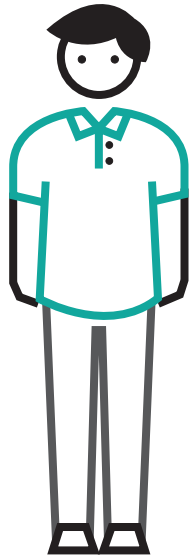


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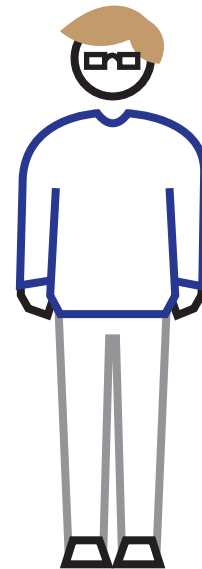


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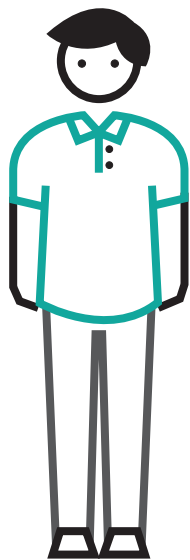


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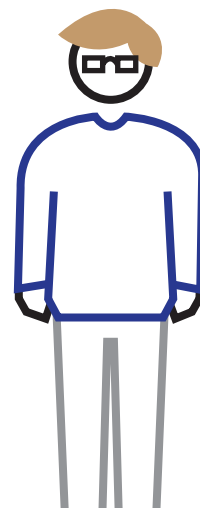
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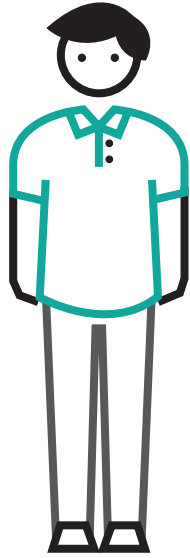


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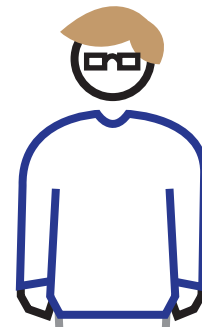


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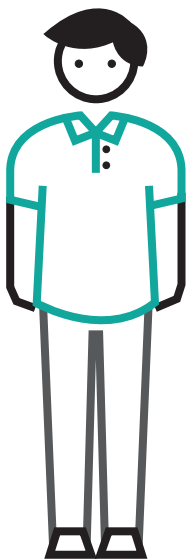


**Julien BOBROFF**  
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**Institut de Physique des Solides**  
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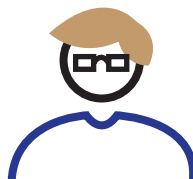


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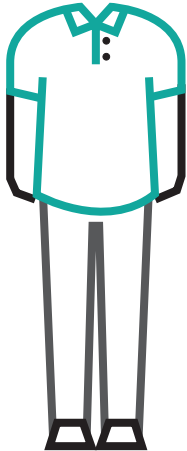
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**de Physique des Solides**  
**Paris-Sud, CNRS**





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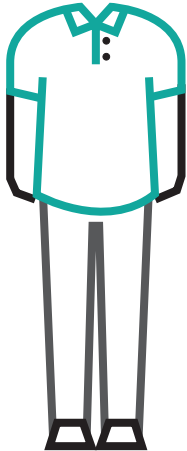
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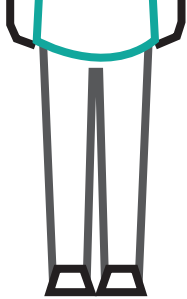


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